

Reconnection-mediated turbulence

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Outline

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- ☞ Context & Motivation: turbulence in space and astrophysical plasmas
- ☞ Magnetic reconnection: basic concepts
- ☞ Phenomenology of Alfvénic turbulence: from weak to strong

2. Further Developments in Theoretical Models

- ☞ From dynamic alignment to reconnection-mediated regime

3. Recent Progress via Numerical Simulations

- ☞ 2D simulation of freely decaying magneto-hydrodynamic (MHD) turbulence
- ☞ 3D collisions of Alfvén-wave packets in reduced MHD

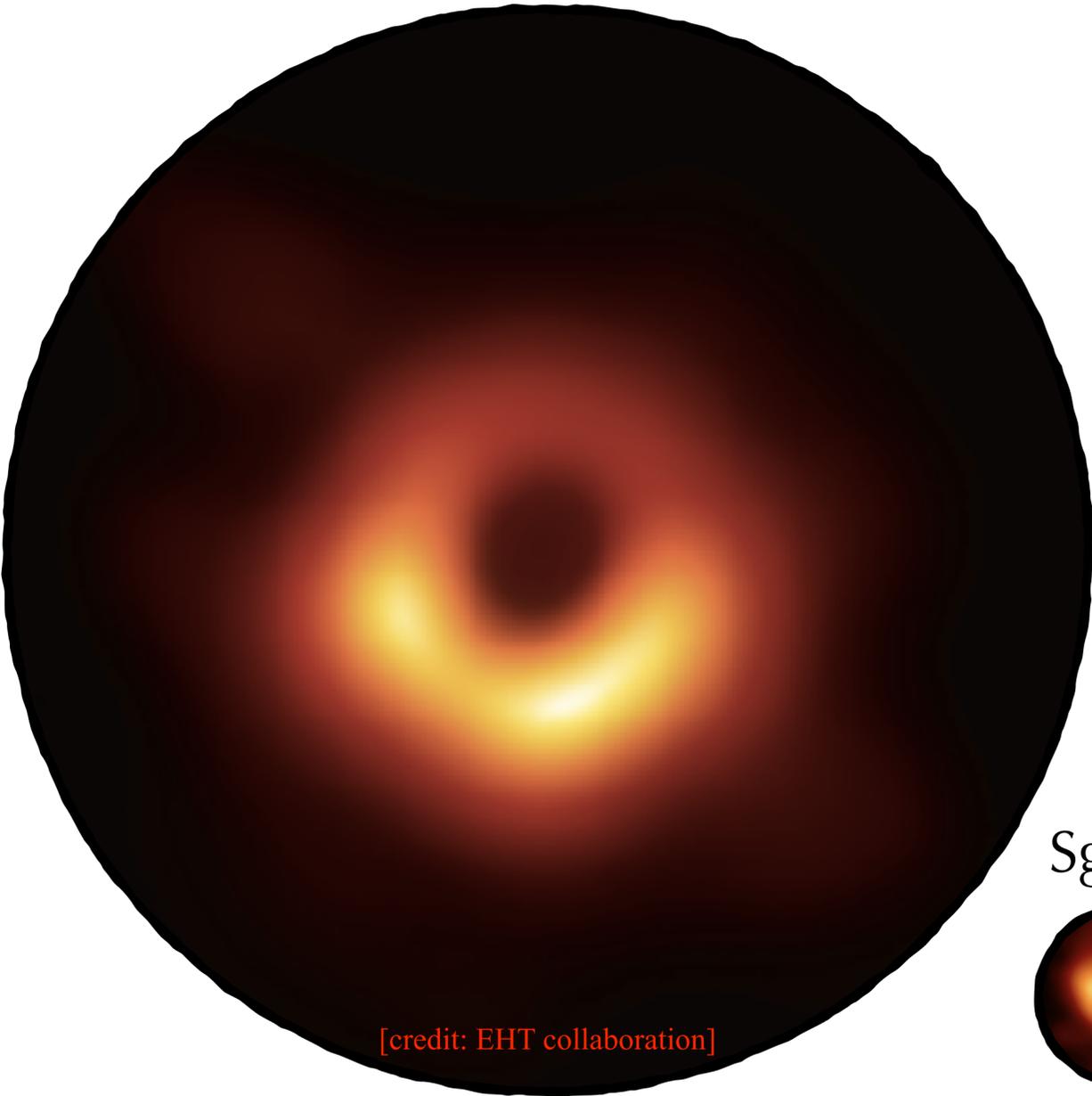
Introduction

Context & Motivation



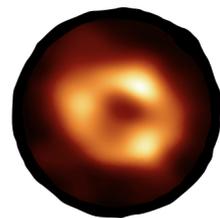
Turbulence in space and astrophysical plasmas

M87*



[credit: EHT collaboration]

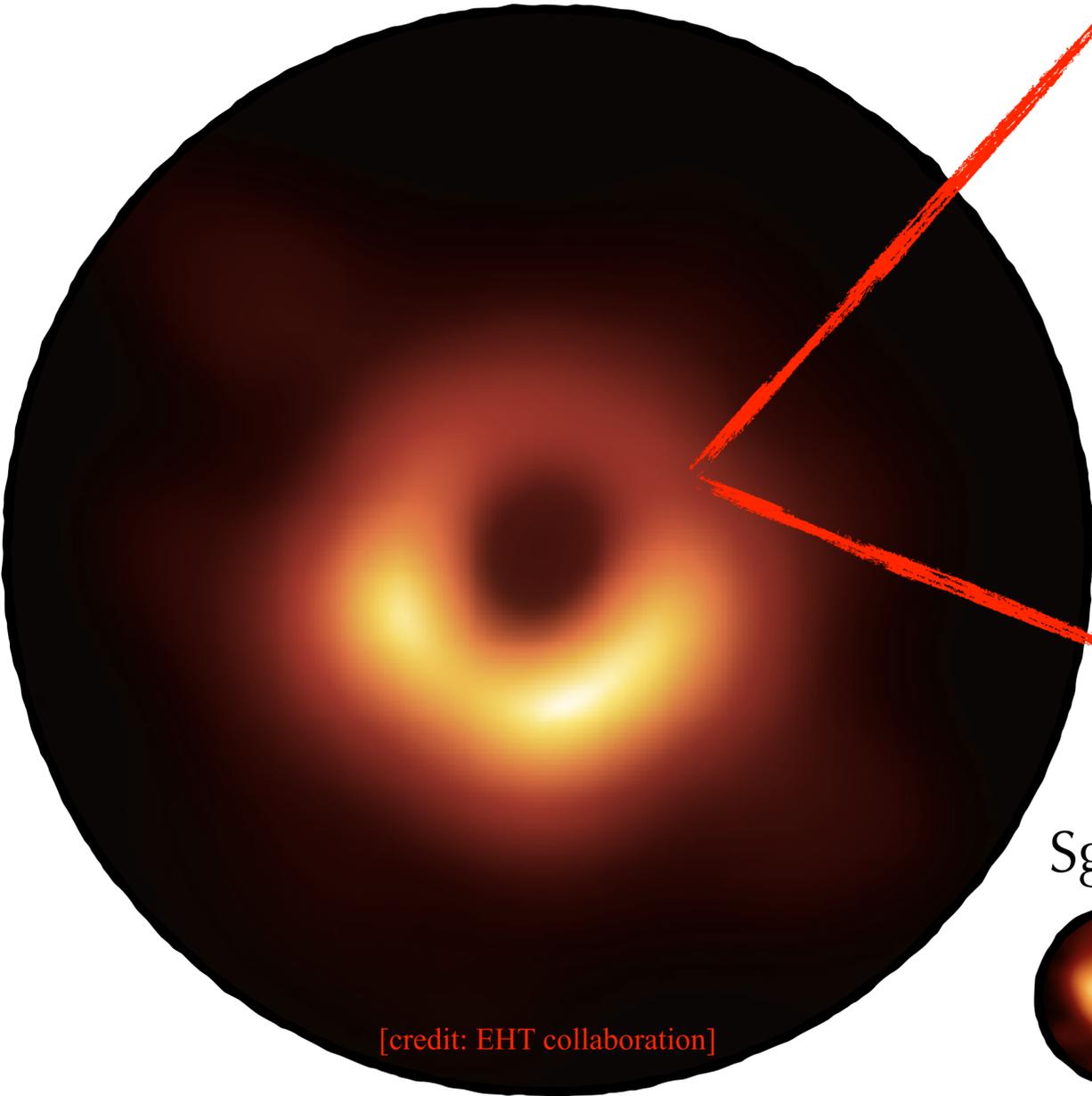
Sgr A*



[credit: EHTC]

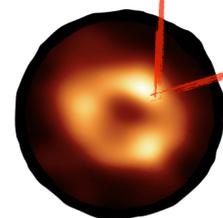
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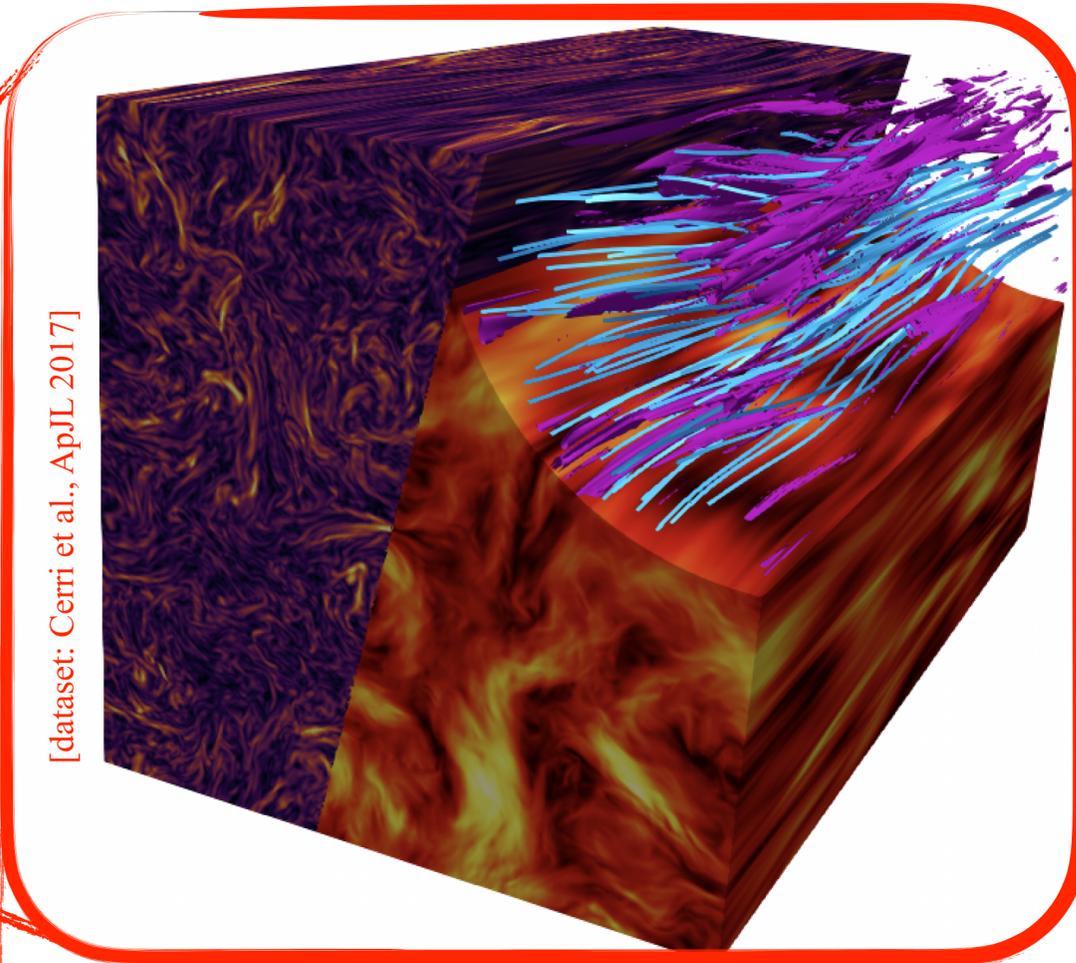


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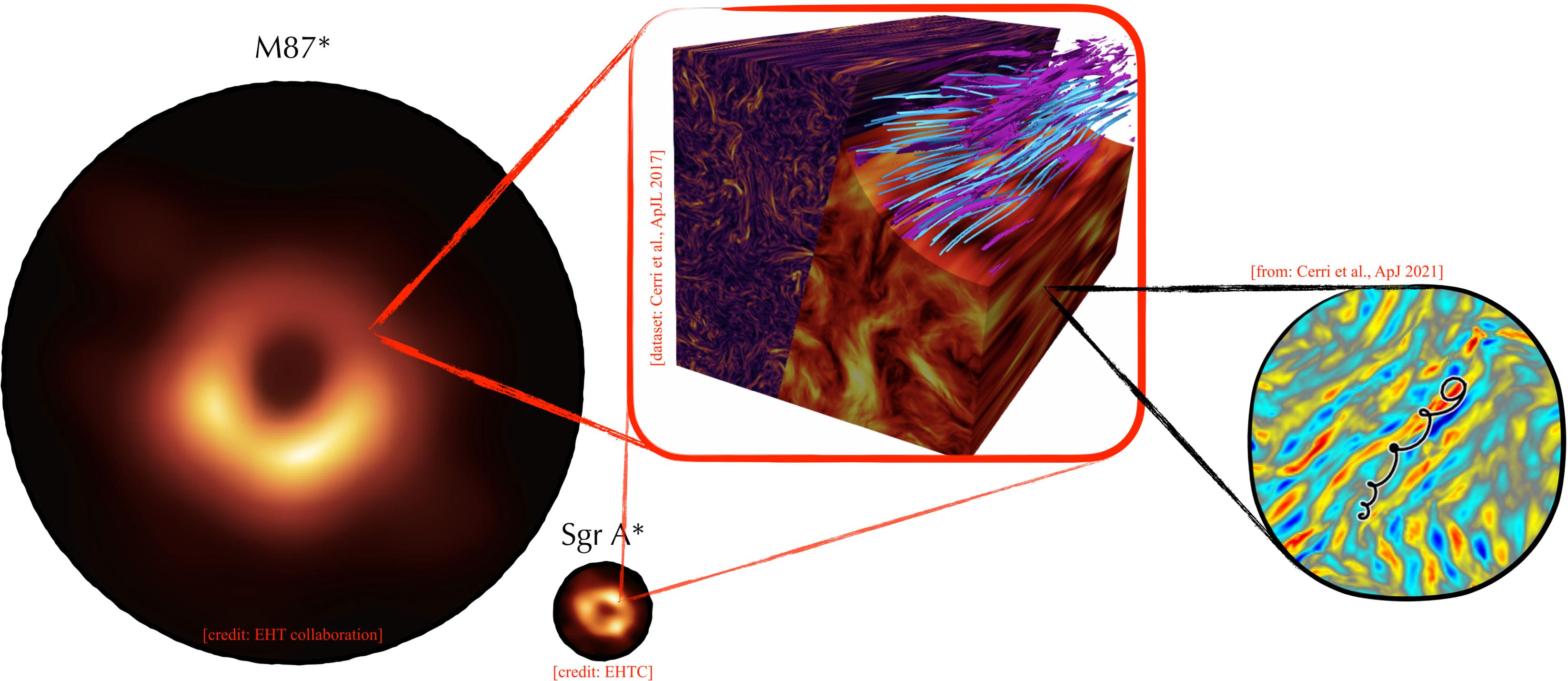


[credit: EHTC]

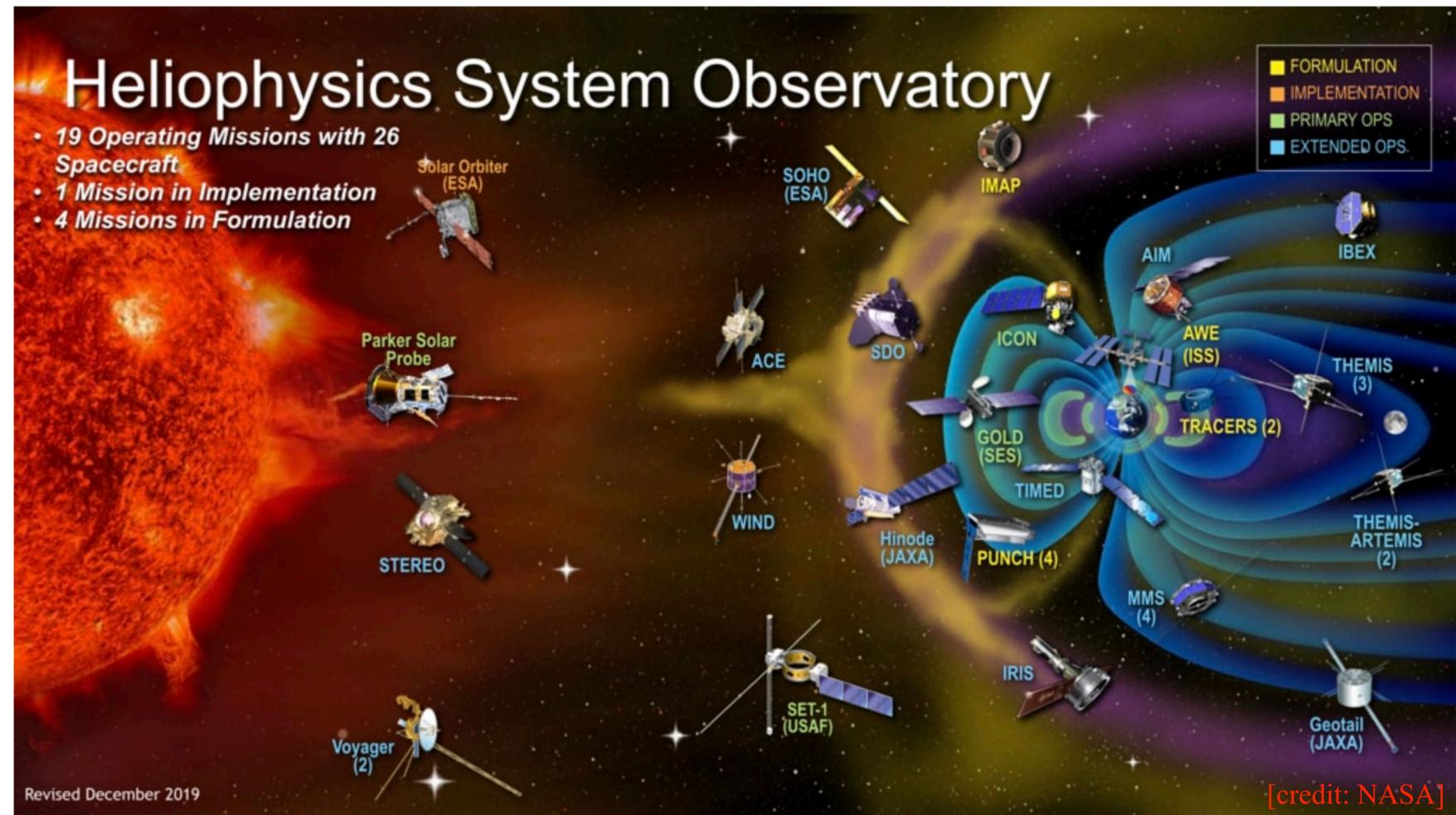
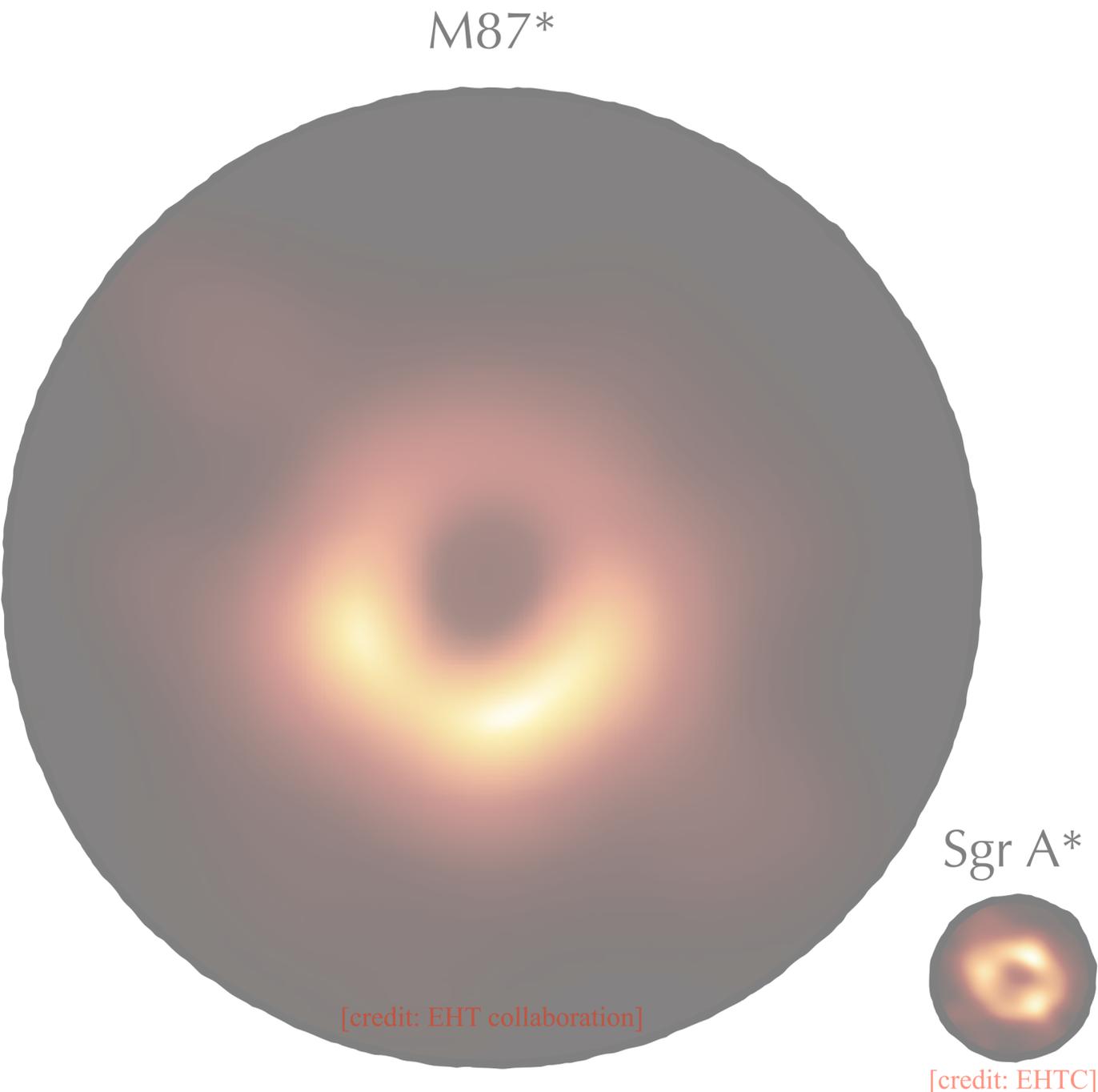


[dataset: Cerri et al., ApJL 2017]

Turbulence in space and astrophysical plasmas

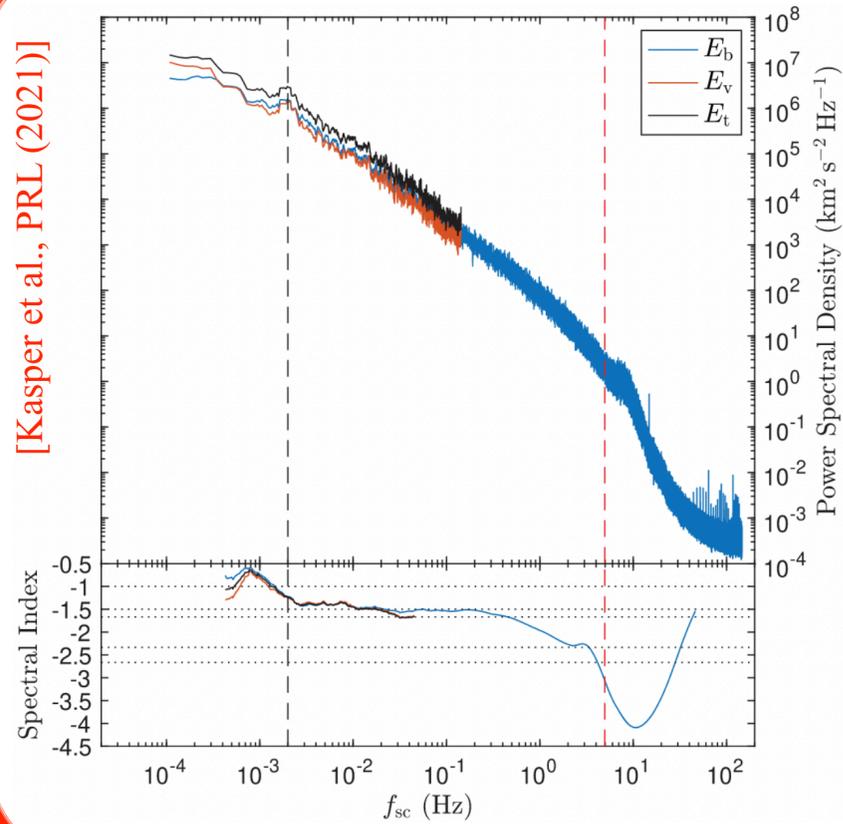


Turbulence in space and astrophysical plasmas

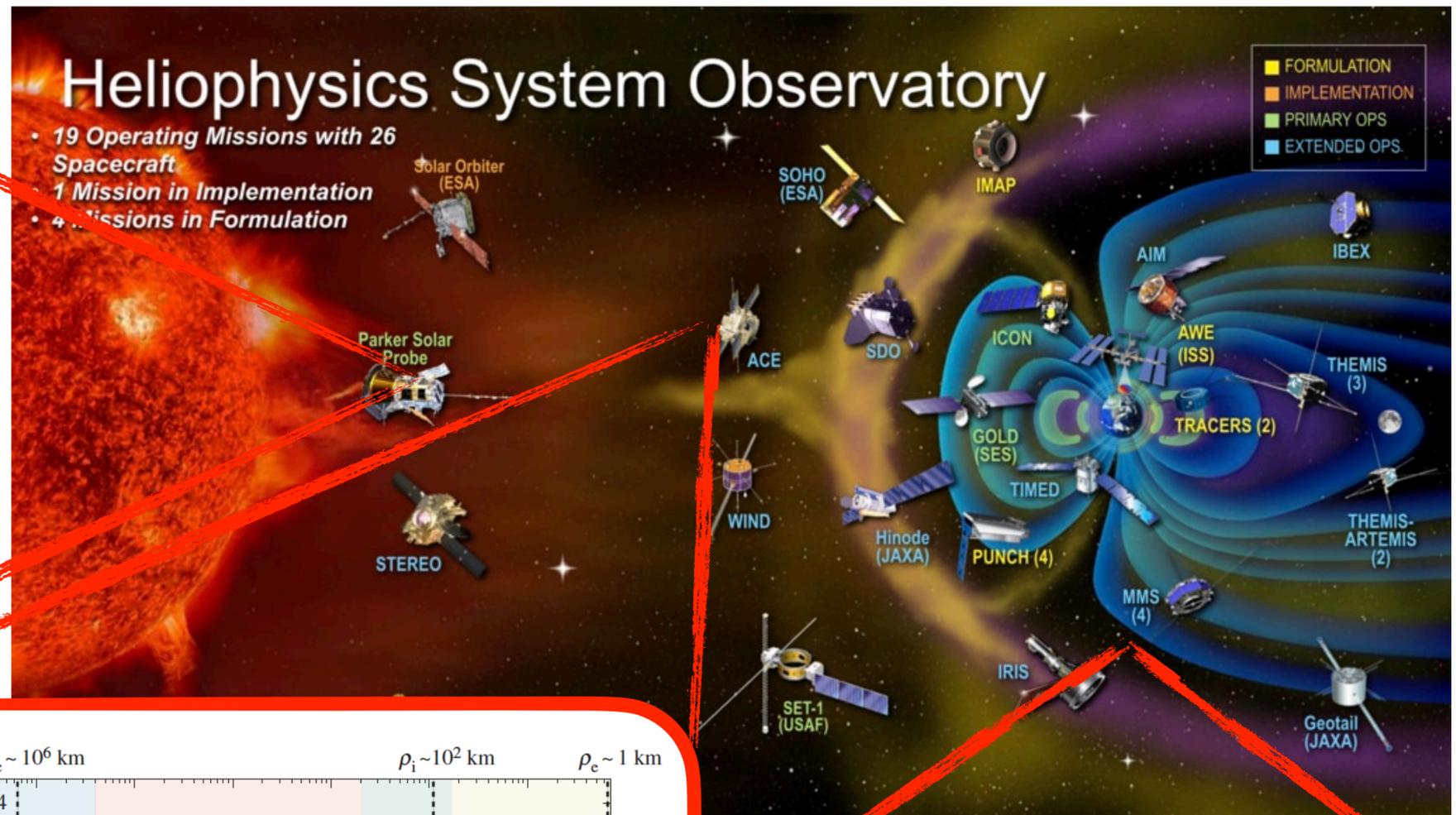


Turbulence in space and astrophysical plasmas

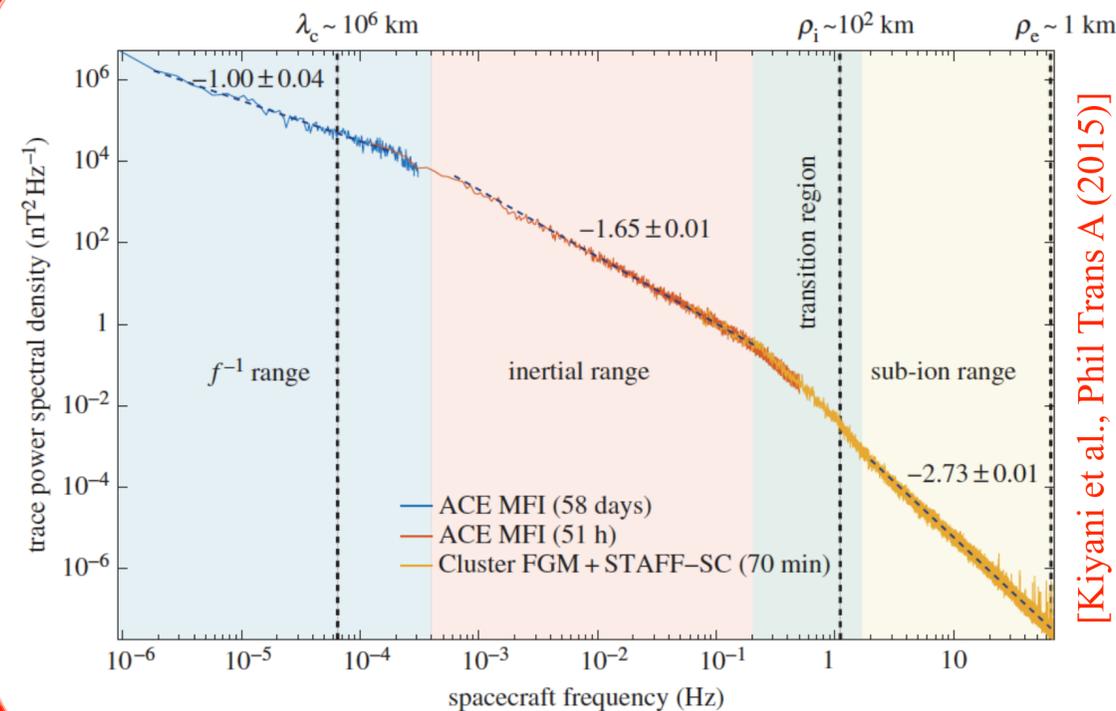
[Kasper et al., PRL (2021)]



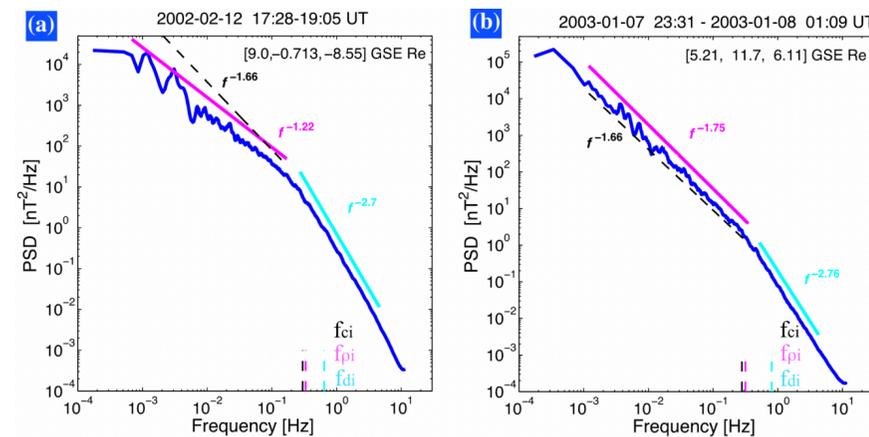
[credit: EHT collaboration]



[credit: NASA]



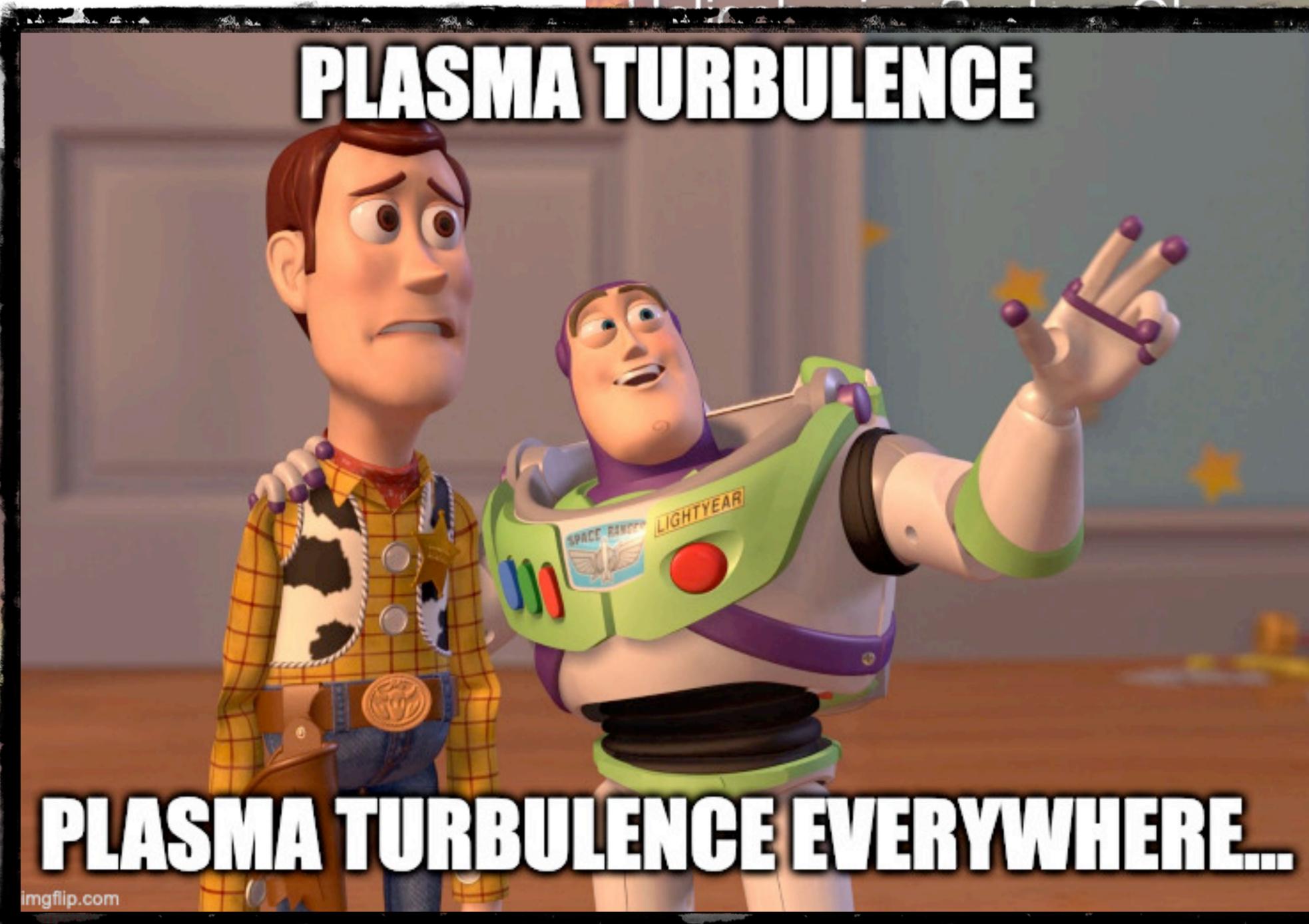
[Kiyani et al., Phil Trans A (2015)]



[Huang et al., ApJL 2017]

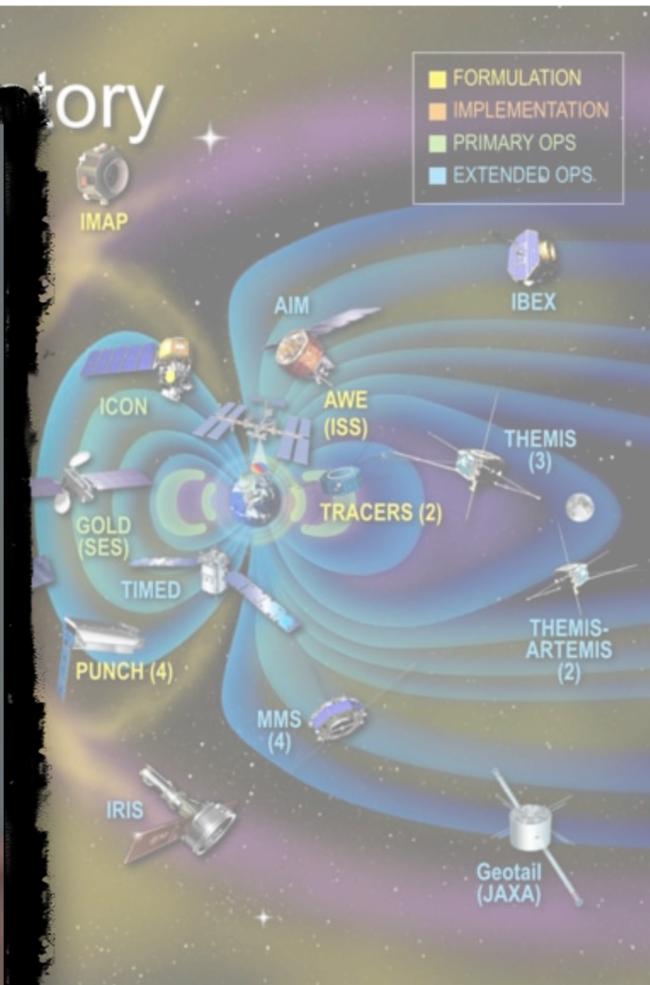
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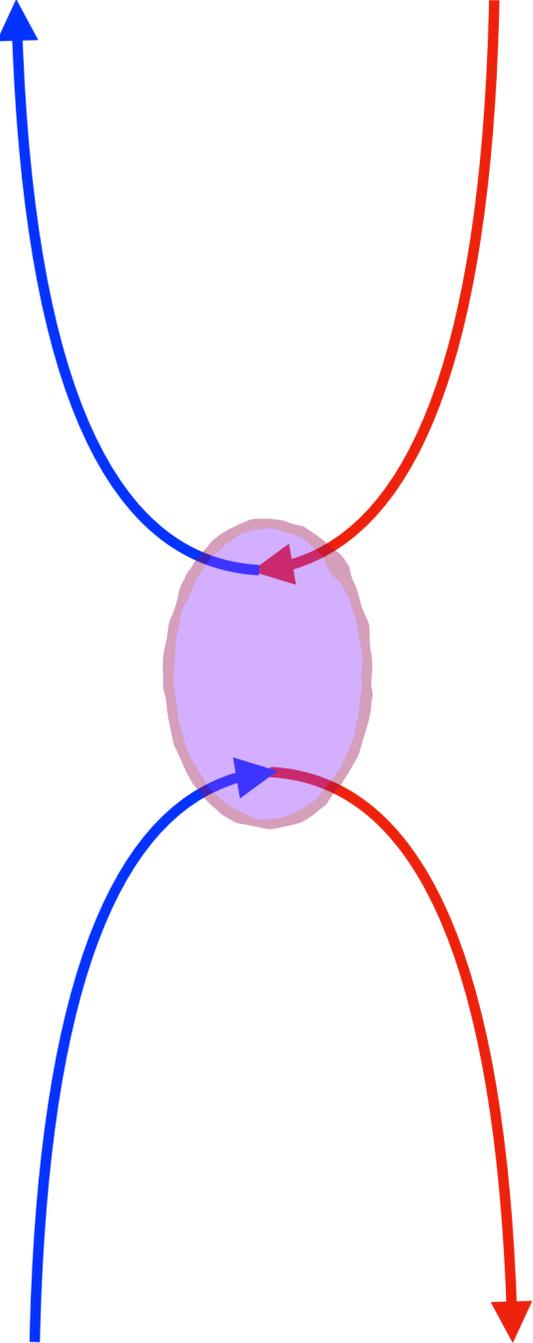
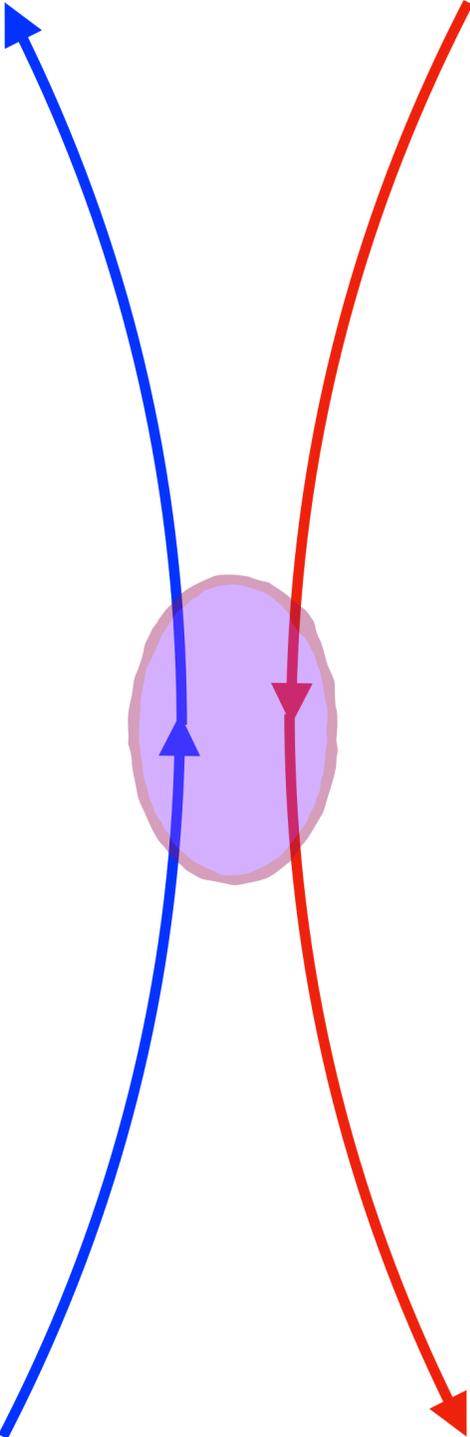
[credit: EHT col]

[credit: EHTC]



Introduction

Magnetic Reconnection



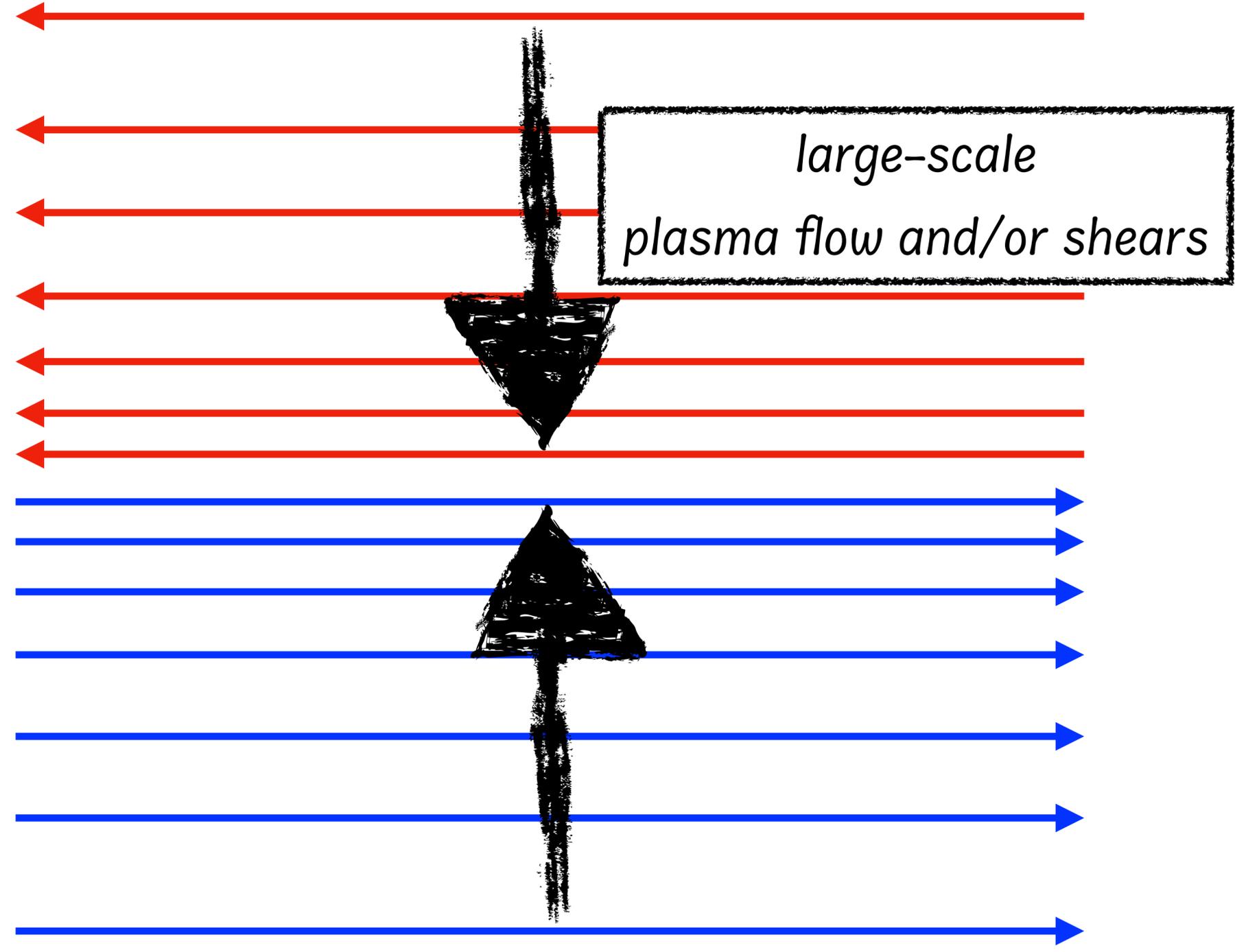
Magnetic reconnection: basic concepts

Magnetic-field reversal



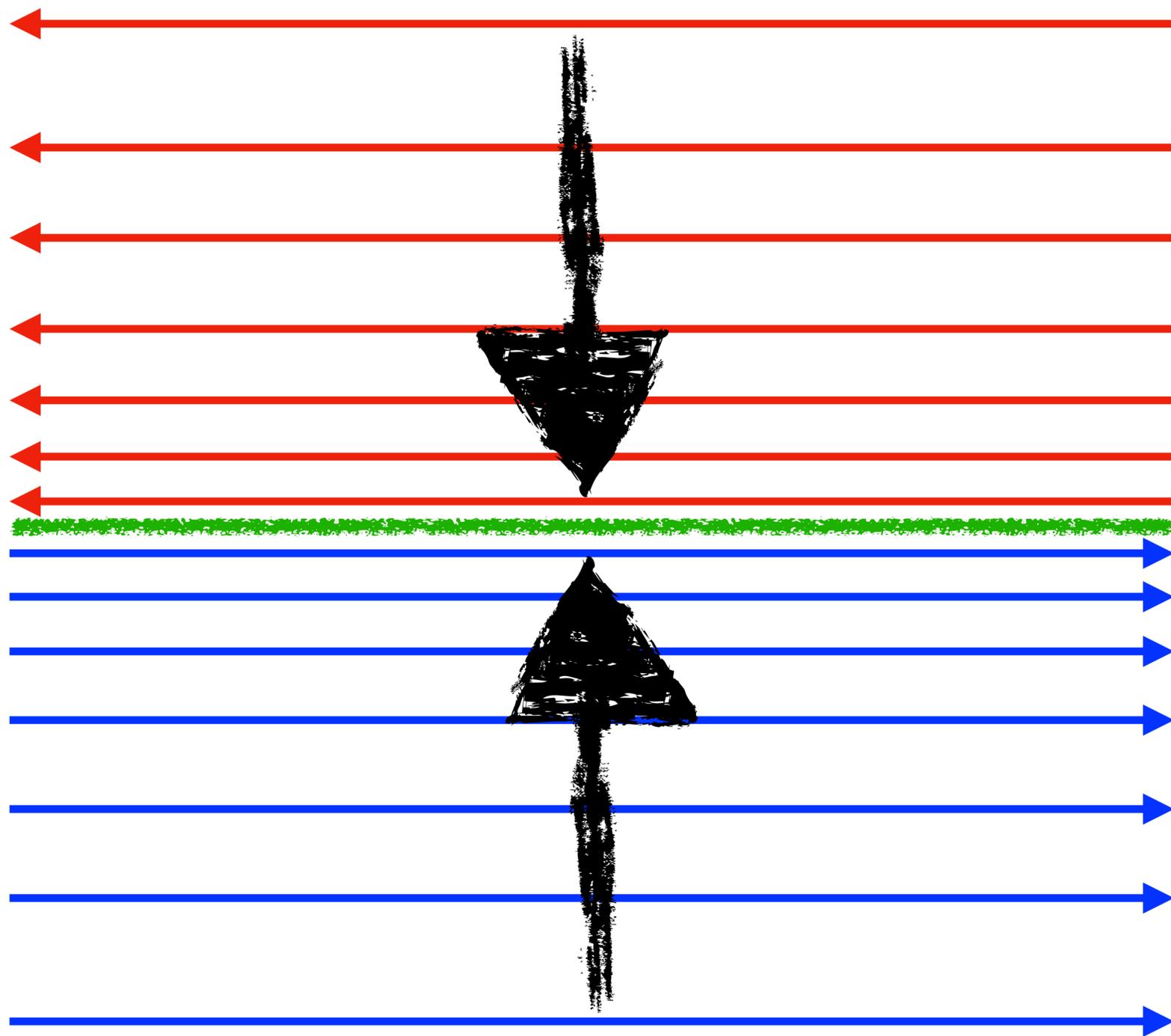
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Magnetic reconnection: basic concepts

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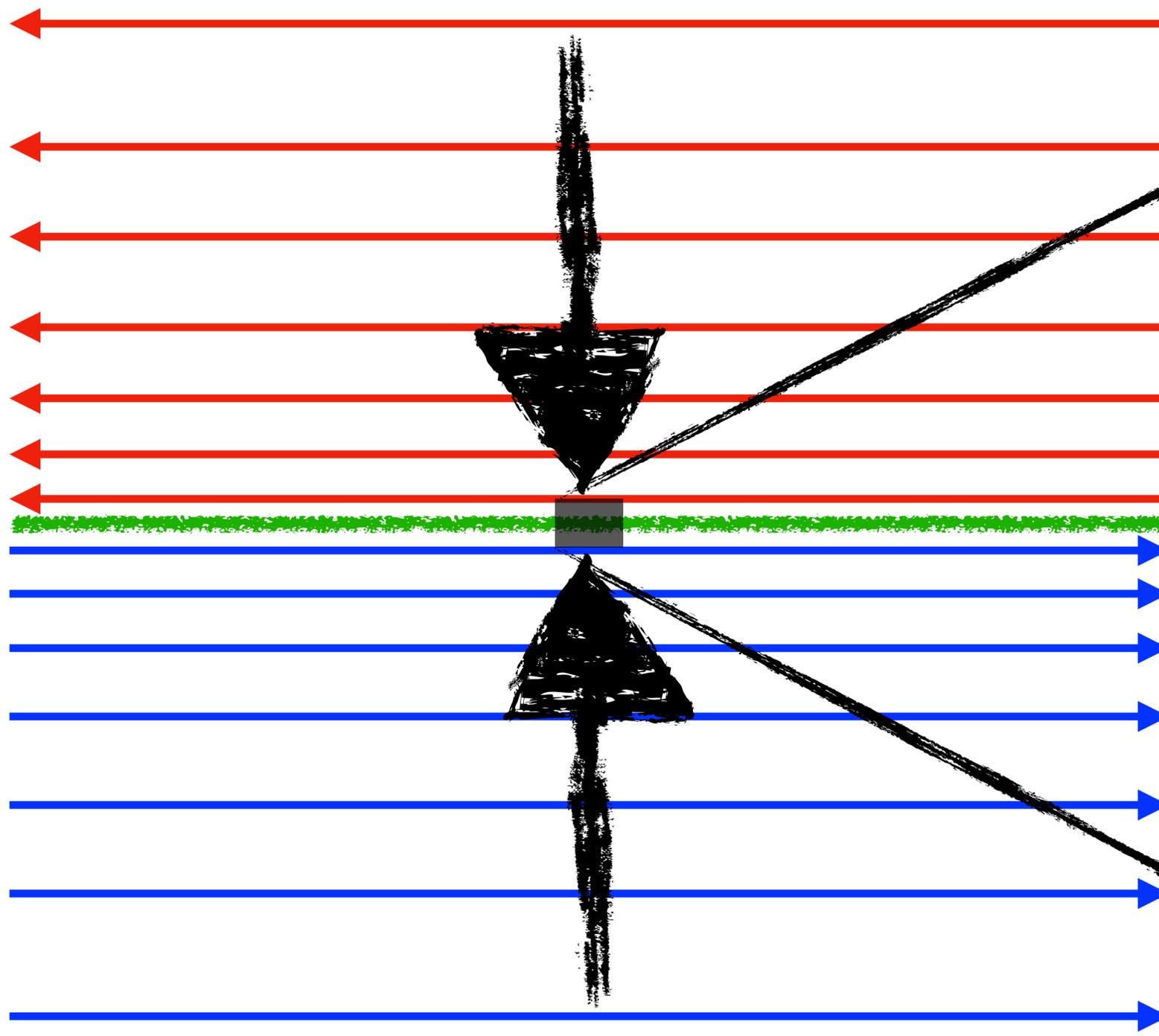


(current sheets = *regions with strong magnetic-field shear*)

current sheet

Magnetic reconnection: basic concepts

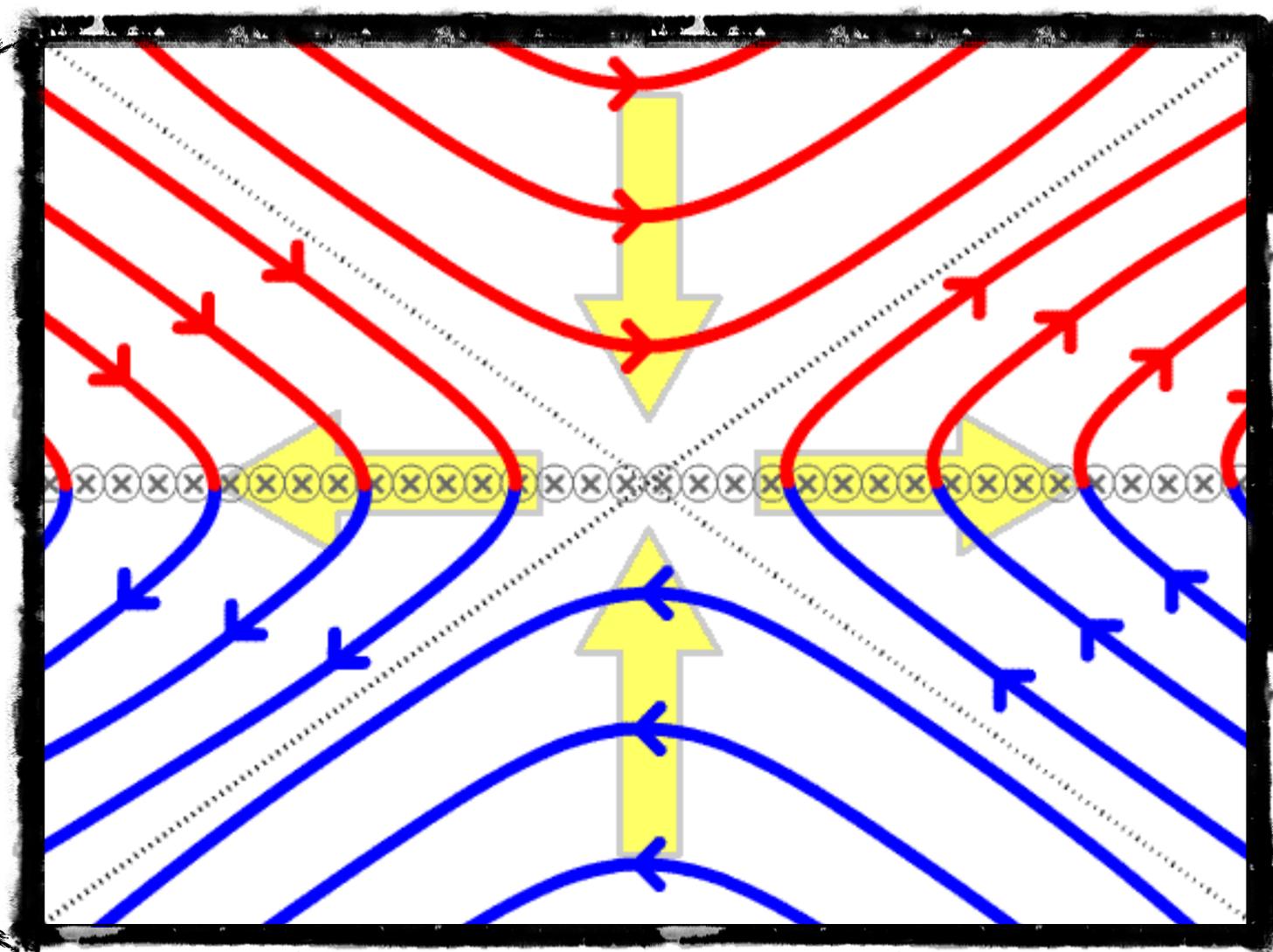
Magnetic-field reversal



if the magnetic shear is “*strong enough*(*)”

⇒ **RECONNECTION!**

(⇒ conversion of magnetic energy into plasma kinetic energy)



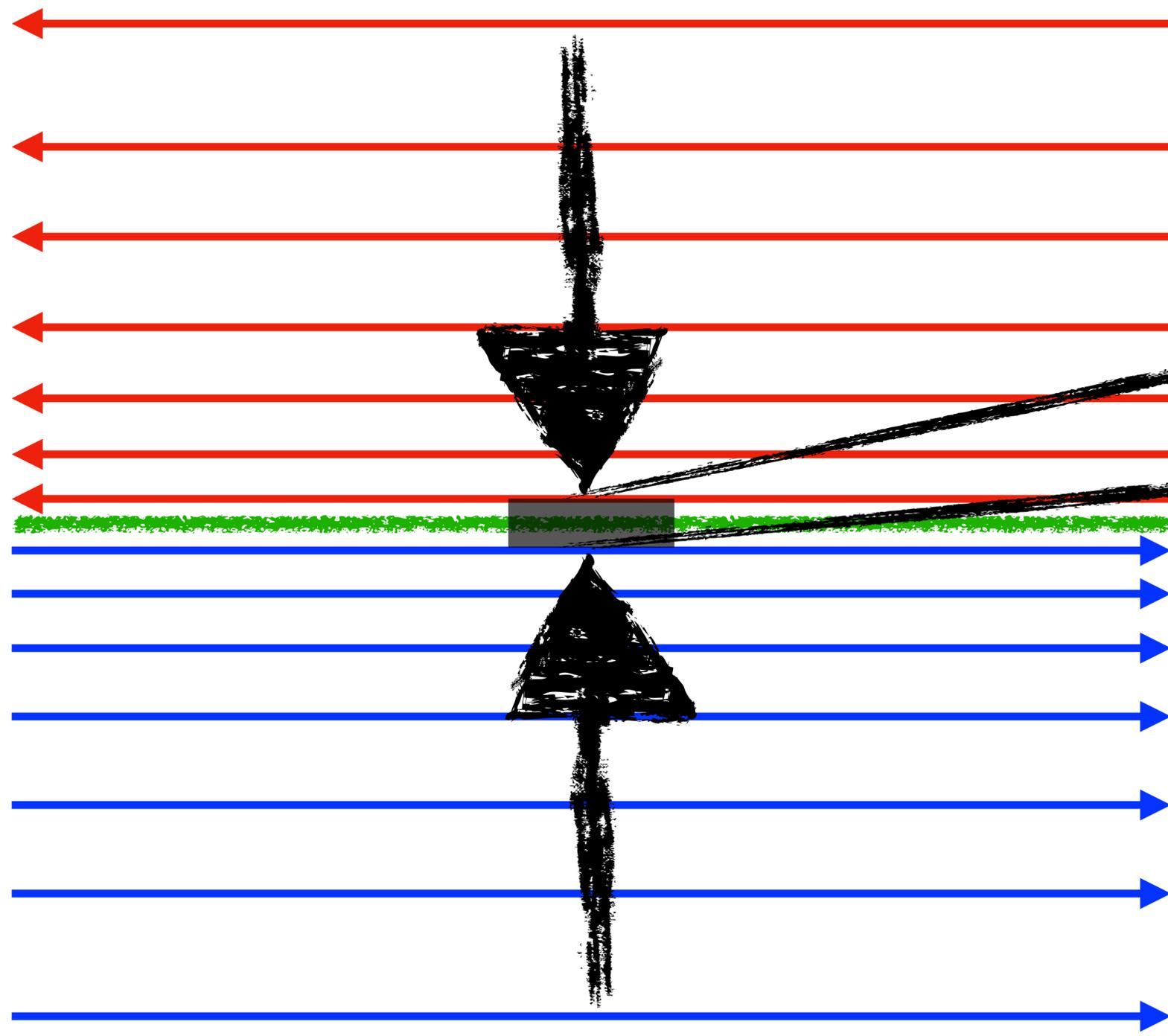
[gif from Wikipedia]

(*) there is a parameter called Δ' ... but this would require an entire lecture!

👉 ask Camille Granier and Emanuele Tassi, they know everything about it!

Magnetic reconnection: basic concepts

Magnetic-field reversal

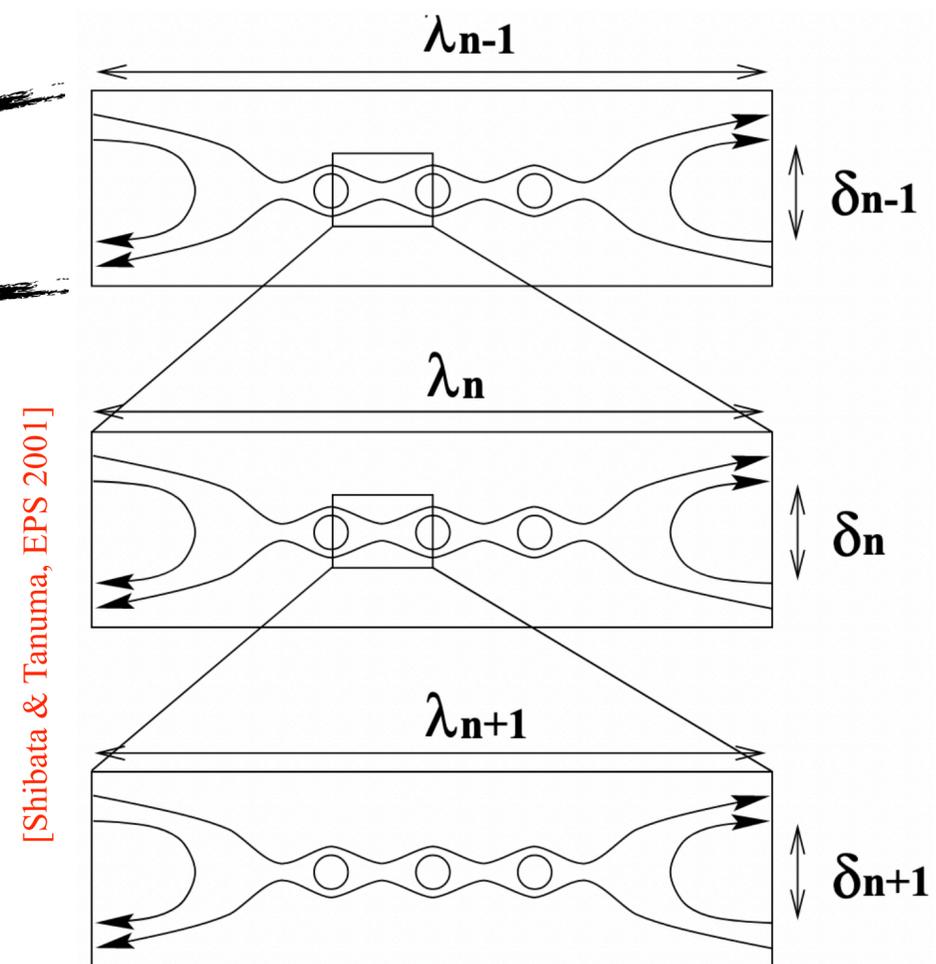


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“**PLASMOID**” regime with recursive/fractal reconnection



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Magnetic reconnection: basic concepts

Magnetic-field

magnetic shear is "strong enough(*)"

RECONNECTION!

plasma kinetic energy)

reconnection

☞ magnetic reconnection requires to **break magnetic-flux conservation**:
this is done *only by non-ideal MHD effects* (non-negligible at "small" scales!)

1. generalized Ohm's law:

$$\mathbf{E} + \frac{\mathbf{u}_i \times \mathbf{B}}{c} = \mathbf{R}$$

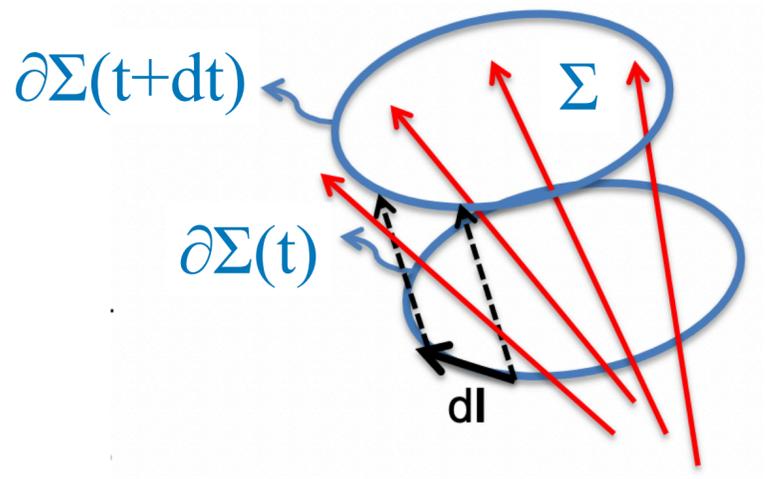
$\mathbf{R} = 0$ in ideal MHD

$$\mathbf{R} \equiv \underbrace{\eta \mathbf{J}}_{(I)} + \underbrace{\frac{\mathbf{J} \times \mathbf{B}}{nec}}_{(II)} - \underbrace{\frac{\nabla \cdot \bar{\mathbf{P}}_e}{ne}}_{(III)} + \underbrace{\frac{m_e}{ne^2} \left[\nabla \cdot \left(\mathbf{J} \mathbf{u}_i + \mathbf{u}_i \mathbf{J} - \frac{\mathbf{J} \mathbf{J}}{ne} \right) + \frac{\partial \mathbf{J}}{\partial t} \right]}_{(IV)}$$

2. Faraday's law:

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

$$\begin{aligned} \Rightarrow \frac{d\Phi_{\mathbf{B}}[\Sigma]}{dt} &\equiv \frac{d}{dt} \int_{\Sigma} \mathbf{B} \cdot d\mathbf{S} \equiv \int_{\Sigma} \partial_t \mathbf{B} \cdot d\mathbf{S} + \oint_{\partial\Sigma} \mathbf{B} \cdot (\mathbf{u} \times d\mathbf{l}) \\ &= -c \int_{\Sigma} (\nabla \times \mathbf{R}) \cdot d\mathbf{S} \end{aligned}$$



but this would require an entire lecture!

... or Emanuele Tassi, they know everything about it!

Magnetic reconnection: basic concepts

Magnetic-field

plasma turbulence:

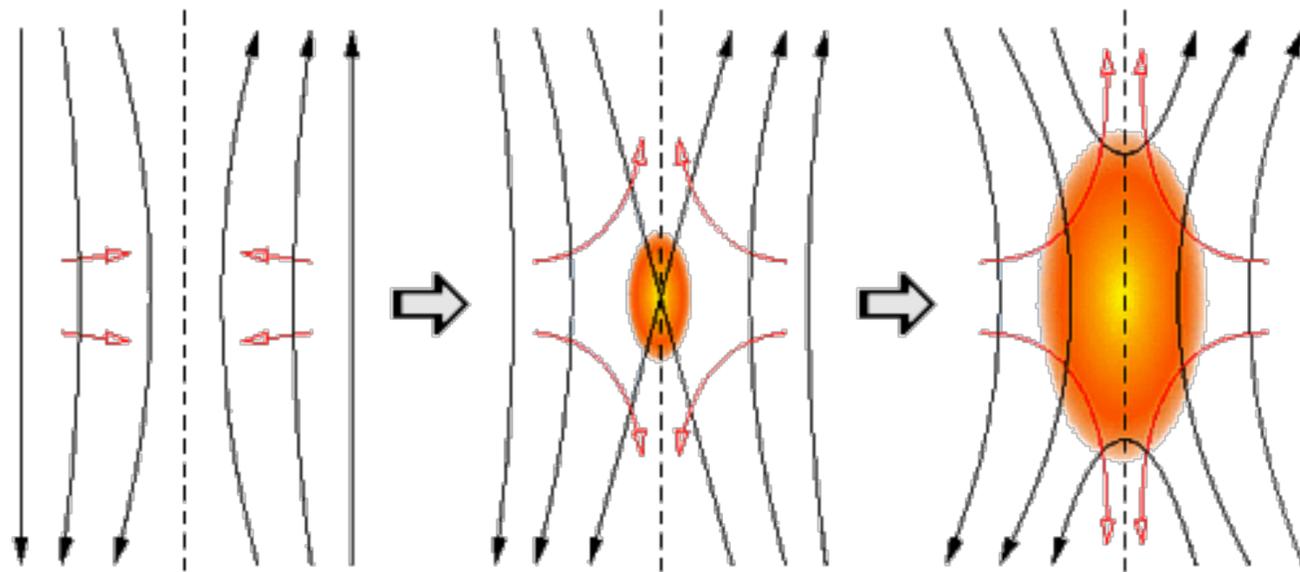
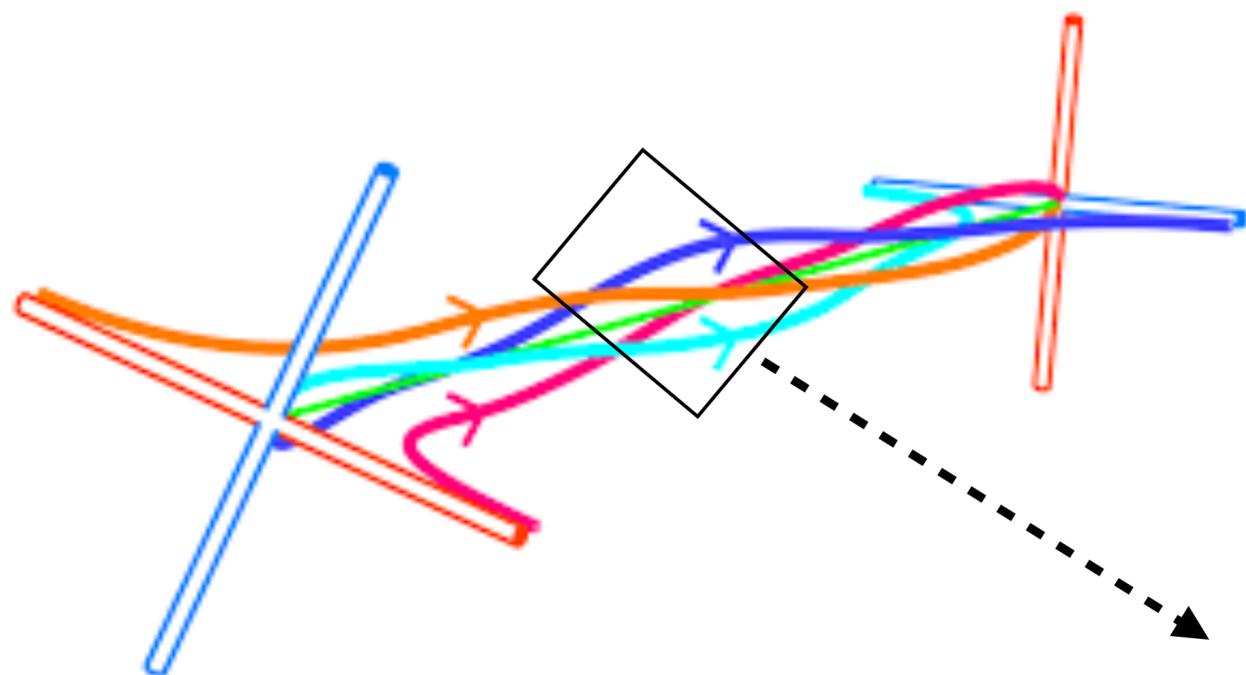
1. *naturally develops current sheets*
2. *naturally generates small scales* (eventually "activating" non-ideal MHD effects)

magnetic shear is "strong enough(*)"

RECONNECTION!

plasma kinetic energy)

reconnection



n+1

but this would require an entire lecture!

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Introduction

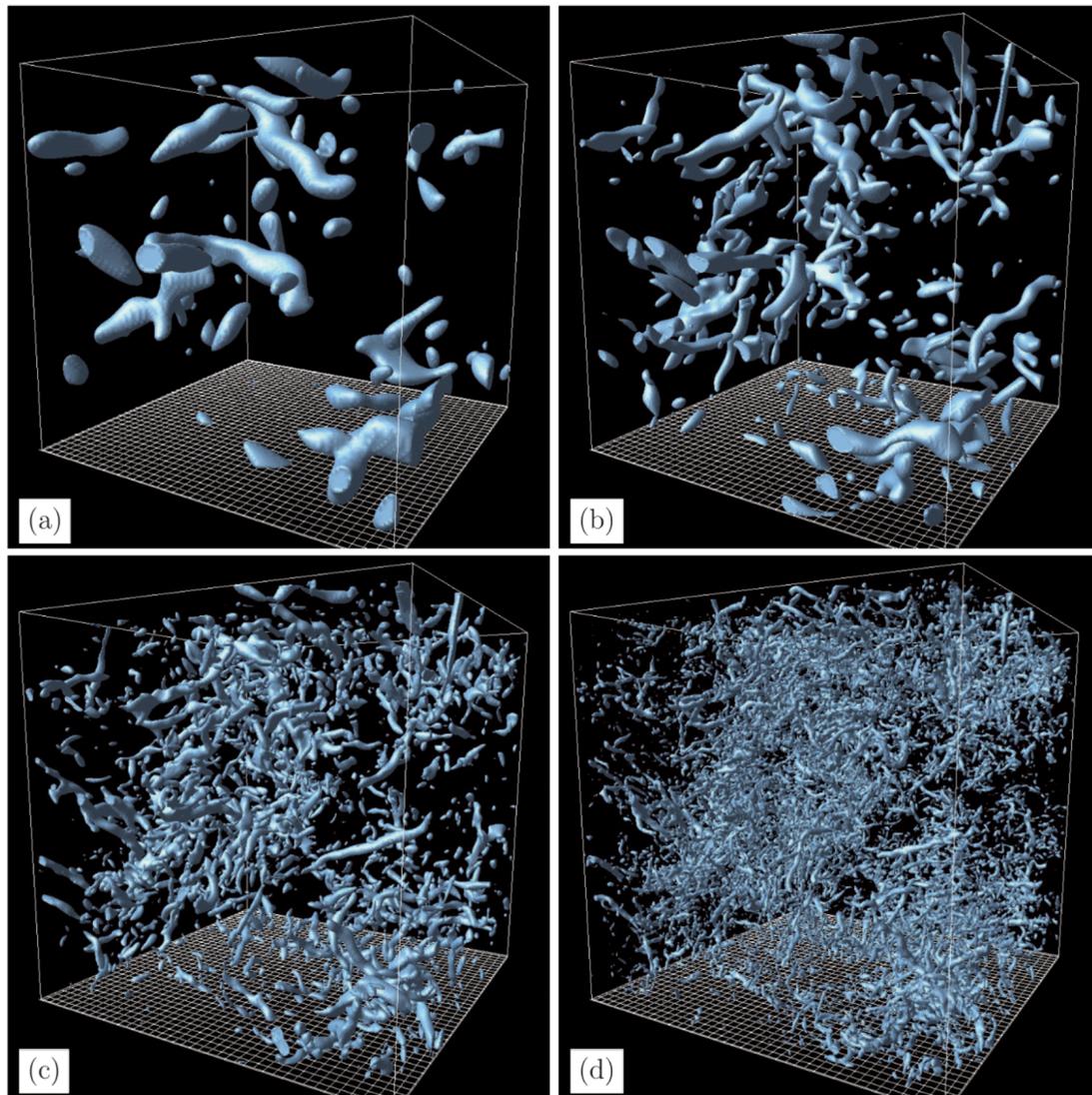


Alfvénic Turbulence

Alfvénic Turbulence

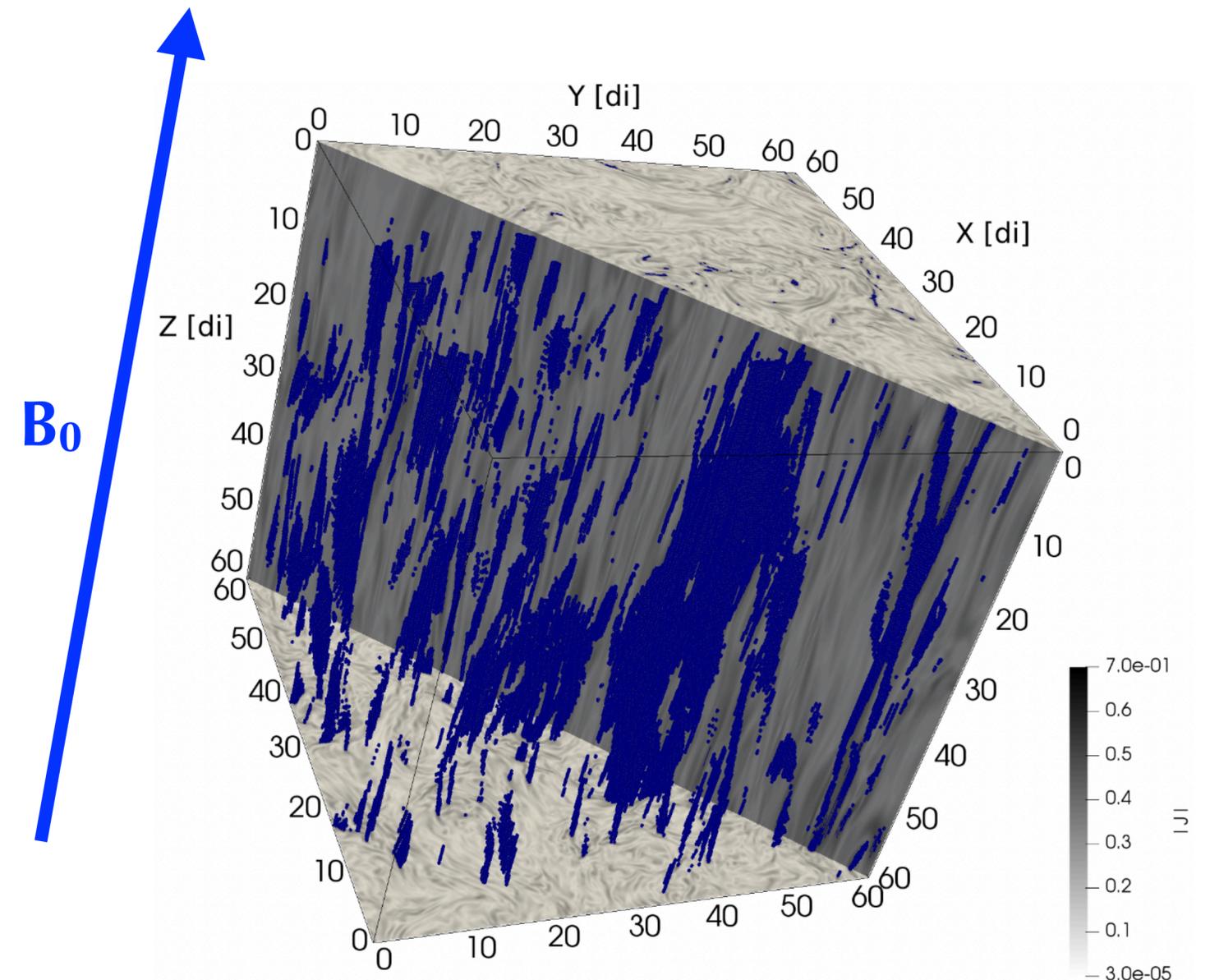
👉 Immediate visual difference: *anisotropy of structures*

hydrodynamic turbulence ($\mathbf{B} = 0$)



[Goto, PTPS 2012]

MHD turbulence ($\mathbf{B} \neq 0$)



[Sisti et al., A&A 2021]

Alfvénic Turbulence

The MHD equations in the Elsässer formulation

$$\frac{\partial \mathbf{z}^+}{\partial t} + \mathbf{z}^- \cdot \nabla \mathbf{z}^+ = -\nabla P;$$

$$\frac{\partial \mathbf{z}^-}{\partial t} + \mathbf{z}^+ \cdot \nabla \mathbf{z}^- = -\nabla P;$$

$$\mathbf{z}^\pm = \mathbf{u} \pm \frac{\mathbf{B}}{\sqrt{4\pi\rho}}$$

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Alfvén waves traveling “up” or “down” the magnetic field \mathbf{B}

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non-linear interaction only between counter-propagating Alfvén waves

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Alfvén waves traveling “up” or “down” the magnetic field \mathbf{B}

Alfvénic turbulence ~ interaction of counter-propagating AWs

Alfvénic Turbulence

Split background and fluctuations:

$$\mathbf{B} = \mathbf{B}_0 + \delta\mathbf{b}$$

$$\mathbf{u} = \mathbf{u}_0 + \delta\mathbf{u}$$

$$v_A = \frac{B_0}{\sqrt{4\pi\rho_0}}$$

$$\delta\mathbf{z}^\pm = \delta\mathbf{u} \pm \frac{\delta\mathbf{b}}{4\pi\rho_0}$$

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$$\Rightarrow \left(\frac{\partial}{\partial t} \mp \mathbf{v}_A \cdot \nabla \right) \delta\mathbf{z}^\pm + (\delta\mathbf{z}^\mp \cdot \nabla) \delta\mathbf{z}^\pm = \dots (!) \text{ turbulence needs finite dissipation!}$$

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non-linear frequency: $\omega_{nl} = k_{\perp} \delta z^{\mp}$

\Rightarrow non-linearity parameter: $\chi \doteq \frac{\omega_{nl}}{\omega_A} = \frac{k_{\perp} \delta z^{\mp}}{k_{\parallel} v_A}$

$\ll 1$ ("WEAK")

~ 1 ("STRONG")

Phenomenology of Alfvénic Turbulence

weak Alfvénic turbulence: a quick phenomenological derivation of the spectrum

☞ *for a formal derivation, see, e.g.,*

[Ng & Bhattacharjee, PoP 1996]

[Galtier, Nazarenko, Newell, Pouquet, JPP 2000]

[Schekochihin, arXiv:2010.00699]

(because, yes, it's the last talk on Friday and we all want to go to lunch!)

Phenomenology of Alfvénic Turbulence

weak Alfvénic turbulence: a quick phenomenological derivation of the spectrum

$$\begin{aligned} \mathbf{k}_1 + \mathbf{k}_2 &= \mathbf{k}_3 \\ \omega(k_{\parallel,1}) + \omega(k_{\parallel,2}) &= \omega(k_{\parallel,3}) \end{aligned} \Rightarrow \text{no parallel cascade } (k_{\parallel} = \text{cst.}) \text{ only a cascade in } k_{\perp}!$$

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(i.e., $\Delta(\delta z)/\delta z \sim 1$)

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\Rightarrow assume changes accumulates
as a random walk:

$$N_{\text{inter.}} \sim \left(\frac{\delta z}{\Delta(\delta z)} \right)^2 \sim \frac{1}{\chi^2} \Rightarrow \tau_{\text{casc}} \sim N \tau_A \sim \frac{\tau_{\text{nl}}^2}{\tau_A} = \frac{\tau_{\text{nl}}}{\chi} \quad \text{CASCADE TIME}$$

Phenomenology of Alfvénic Turbulence

weak Alfvénic turbulence: a quick phenomenological derivation of the spectrum

👉 fluctuations' scaling and energy spectrum
from constant energy flux through scales:

$$\frac{\delta z^2}{\tau_{\text{casc}}} \sim \varepsilon = \text{const.}$$

\Rightarrow

$$\delta z \propto k_{\perp}^{-1/2}$$

\Rightarrow

$$\mathcal{E}_{\delta z} \propto k_{\perp}^{-2}$$

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⚠ A very important consequence of these scalings is that **an initially weak Alfvénic cascade will not remain weak!**

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non-linear frequency increases with decreasing scales,
while linear frequency is constant because there is no parallel cascade:

$$\omega_{\text{nl}} = k_{\perp} \delta z \sim k_{\perp}^{1/2}$$

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\Rightarrow

$$\chi \sim k_{\perp}^{1/2}$$

\Rightarrow

$$\frac{\lambda_{\perp}^{\text{CB}}}{\ell_{\parallel,0}} \sim \left(\frac{\varepsilon \ell_{\parallel,0}}{v_{\text{A}}^3} \right)^{1/2} \sim \left(\frac{\delta z_0}{v_{\text{A}}} \right)^{3/2} \approx \chi_0^{3/2} \quad (\ll 1)$$

transition to critical balance ($\chi \sim 1$)

Phenomenology of Alfvénic Turbulence

critically balanced (strong) Alfvénic turbulence: a quick phenomenological derivation

☞ *for further details, see, e.g.,*

[Goldreich & Sridhar, ApJ 1995]

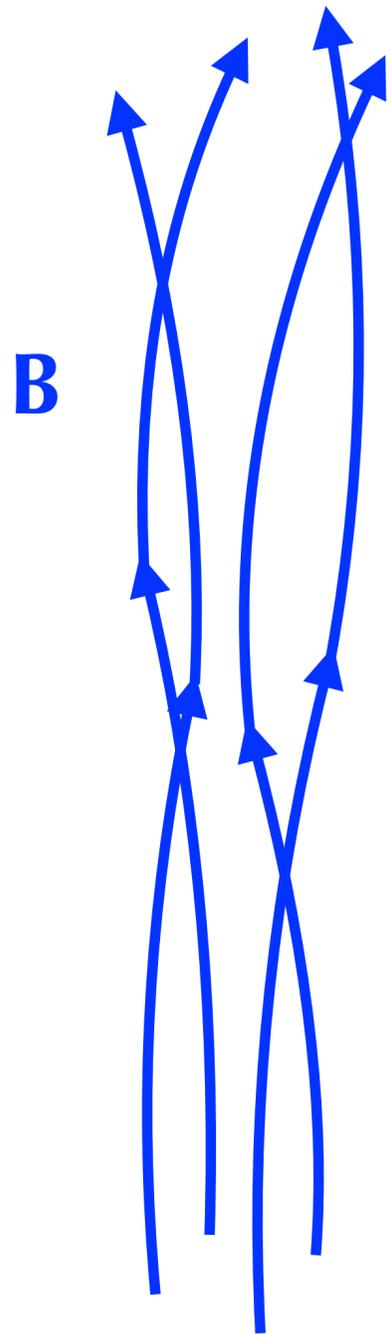
[Oughton & Matthaeus, ApJ 2020]

[Schekochihin, arXiv:2010.00699]

(same reason as before)

Phenomenology of Alfvénic Turbulence

critically balanced (strong) Alfvénic turbulence: a quick phenomenological derivation



☞ At this point, linear, non-linear, and cascade timescales match each other:

$$\tau_{nl} \sim \tau_A \quad \Rightarrow \quad \tau_{casc} \sim \tau_{nl}$$

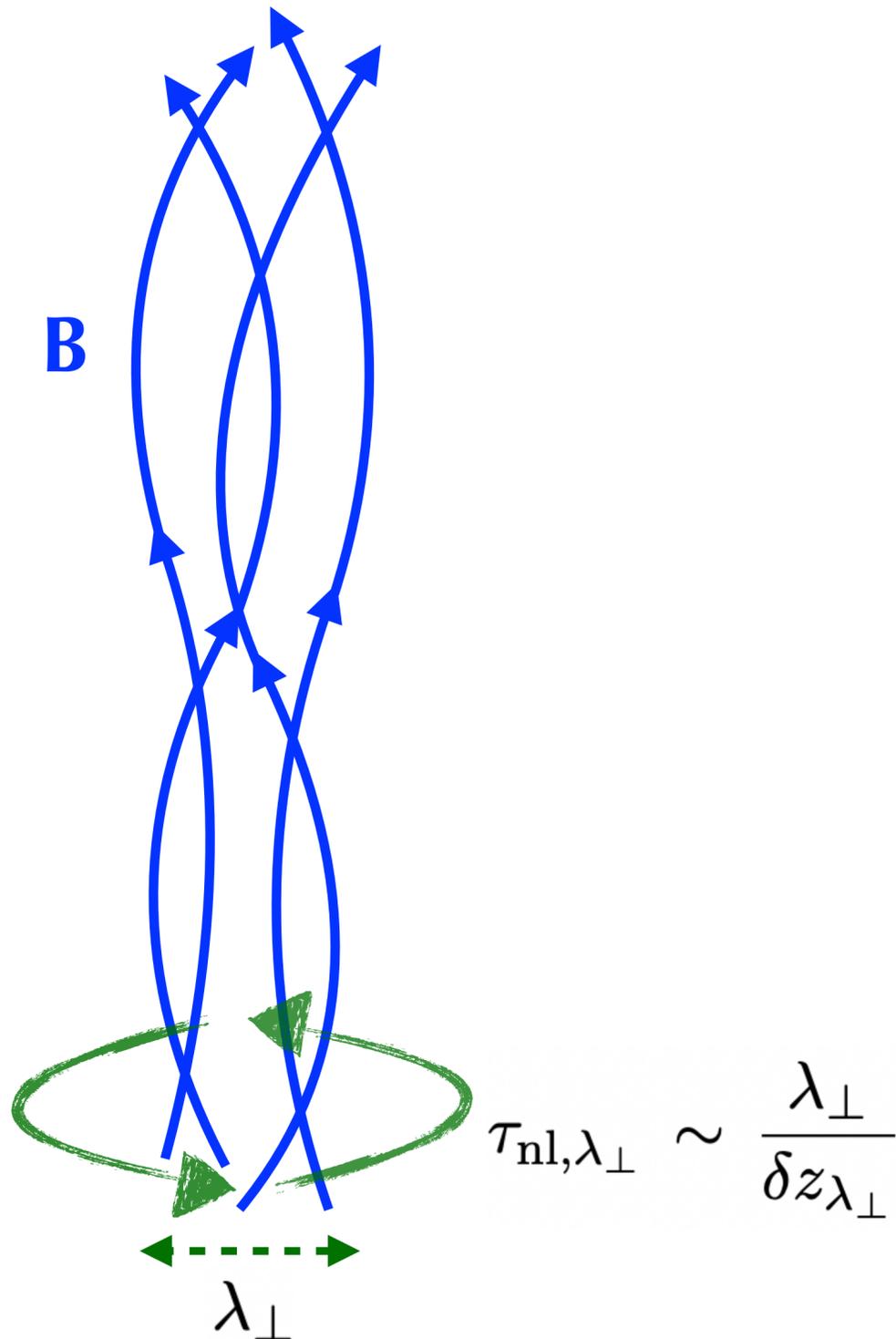
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you can see the *“critical-balance condition”* as the result of causality:



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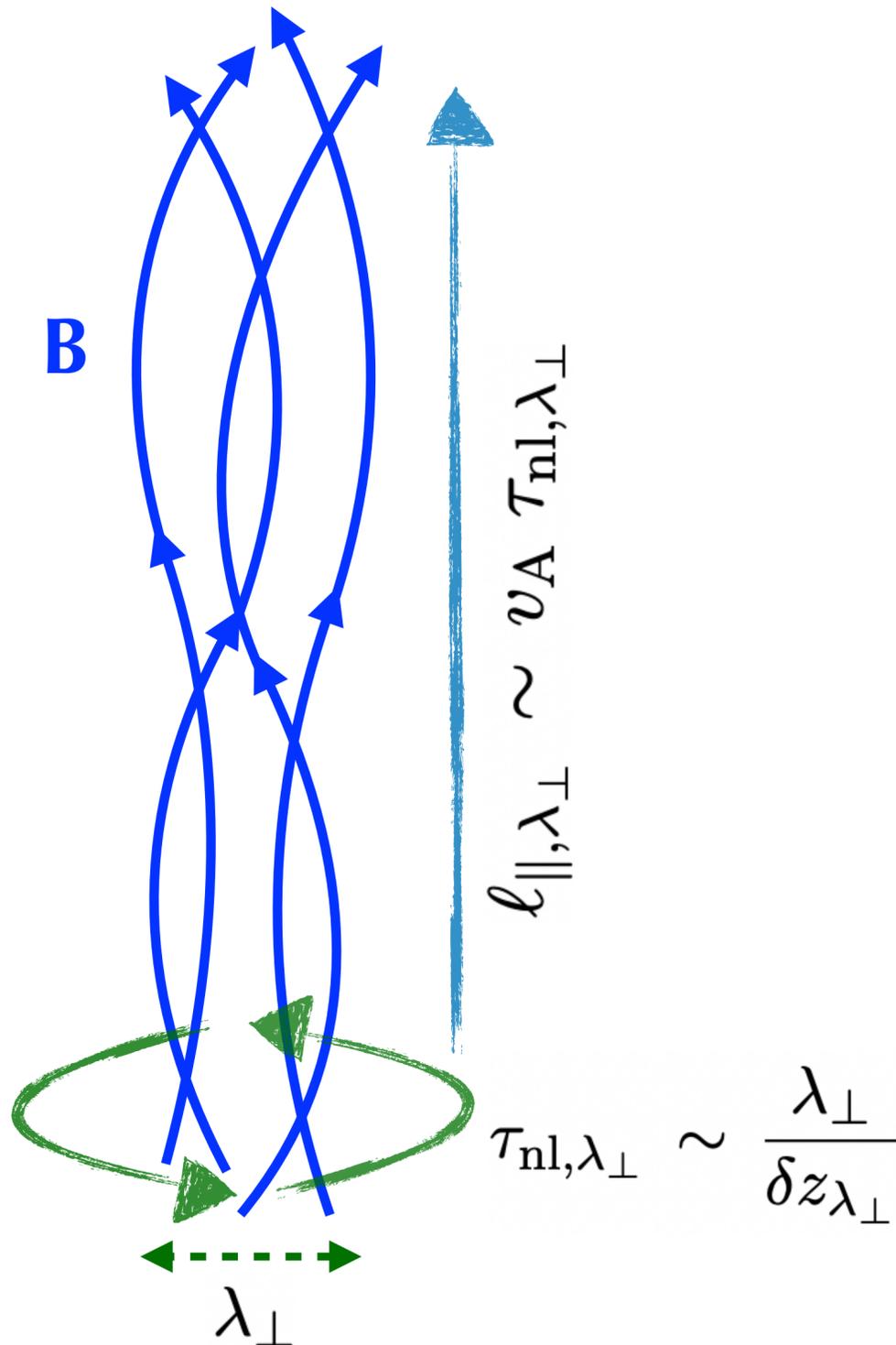
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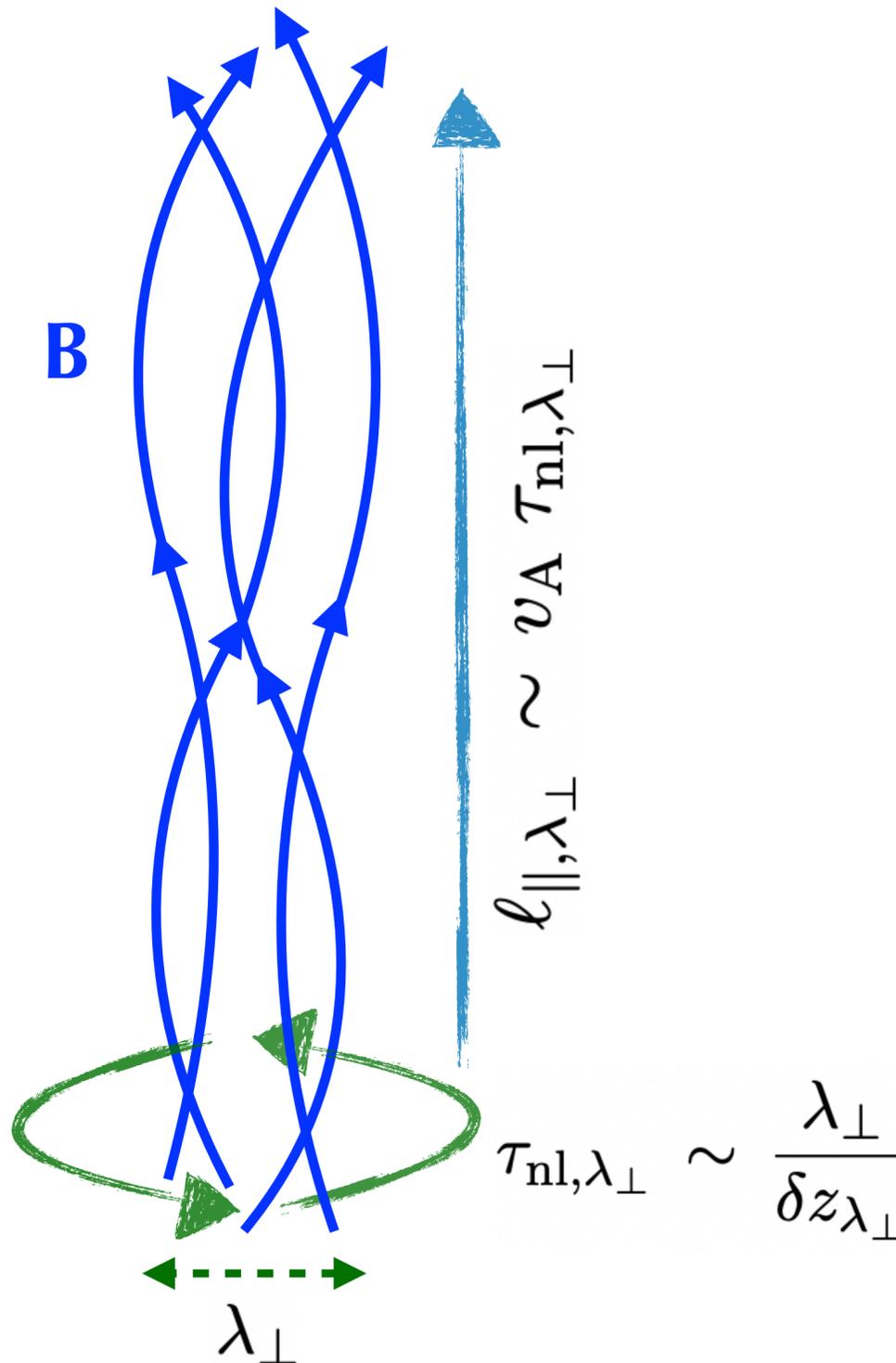
the information about Alfvénic fluctuations decorrelating in the perpendicular plane over an eddy turn-over time τ_{nl} can only propagate along the field for a length $\ell_{||}$ at maximum speed v_A .

“So... CB is essentially AWs trying to keep up with the turbulent eddies...”



Phenomenology of Alfvénic Turbulence

critically balanced (strong) Alfvénic turbulence: a quick phenomenological derivation



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$$\tau_{nl} \sim \tau_A \Rightarrow \tau_{casc} \sim \tau_{nl}$$

you can see the *“critical-balance condition” as the result of causality:*

the information about Alfvénic fluctuations decorrelating in the perpendicular plane over an eddy turn-over time τ_{nl} can only propagate along the field for a length $l_{||}$ at maximum speed v_A .

“So... CB is essentially AWs trying to keep up with the turbulent eddies...”

Therefore, once $\tau_{nl} \sim \tau_A$ is reached, the balance is maintained.

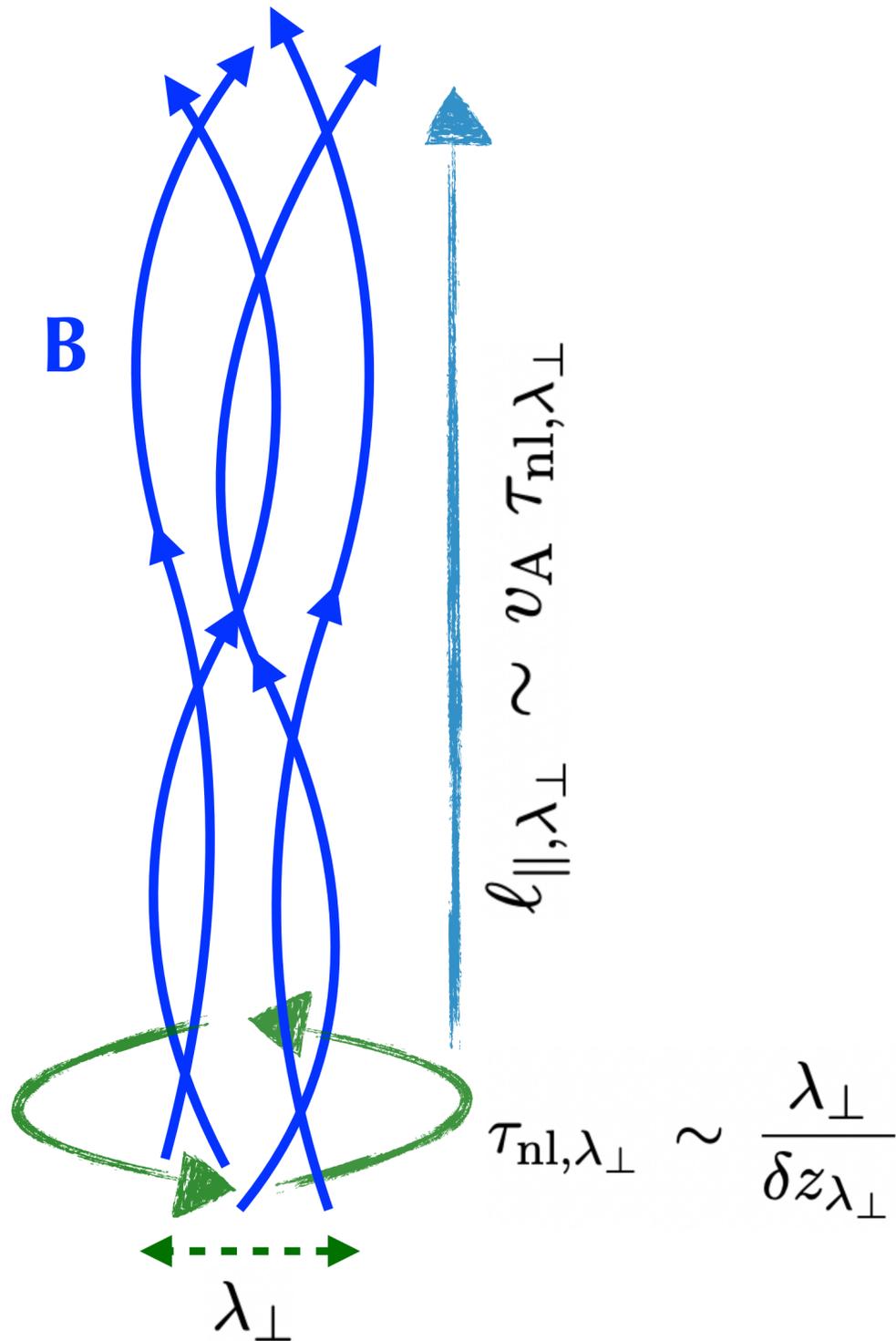
(In principle, this could be done by continuing the cascade with $\tau_{nl} = \text{const.}$, or by generating smaller $l_{||}$ such that $\tau_A \sim l_{||}/v_A \sim \tau_{nl}$ keeps holding... it is the latter)

Phenomenology of Alfvénic Turbulence

critically balanced (strong) Alfvénic turbulence: a quick phenomenological derivation

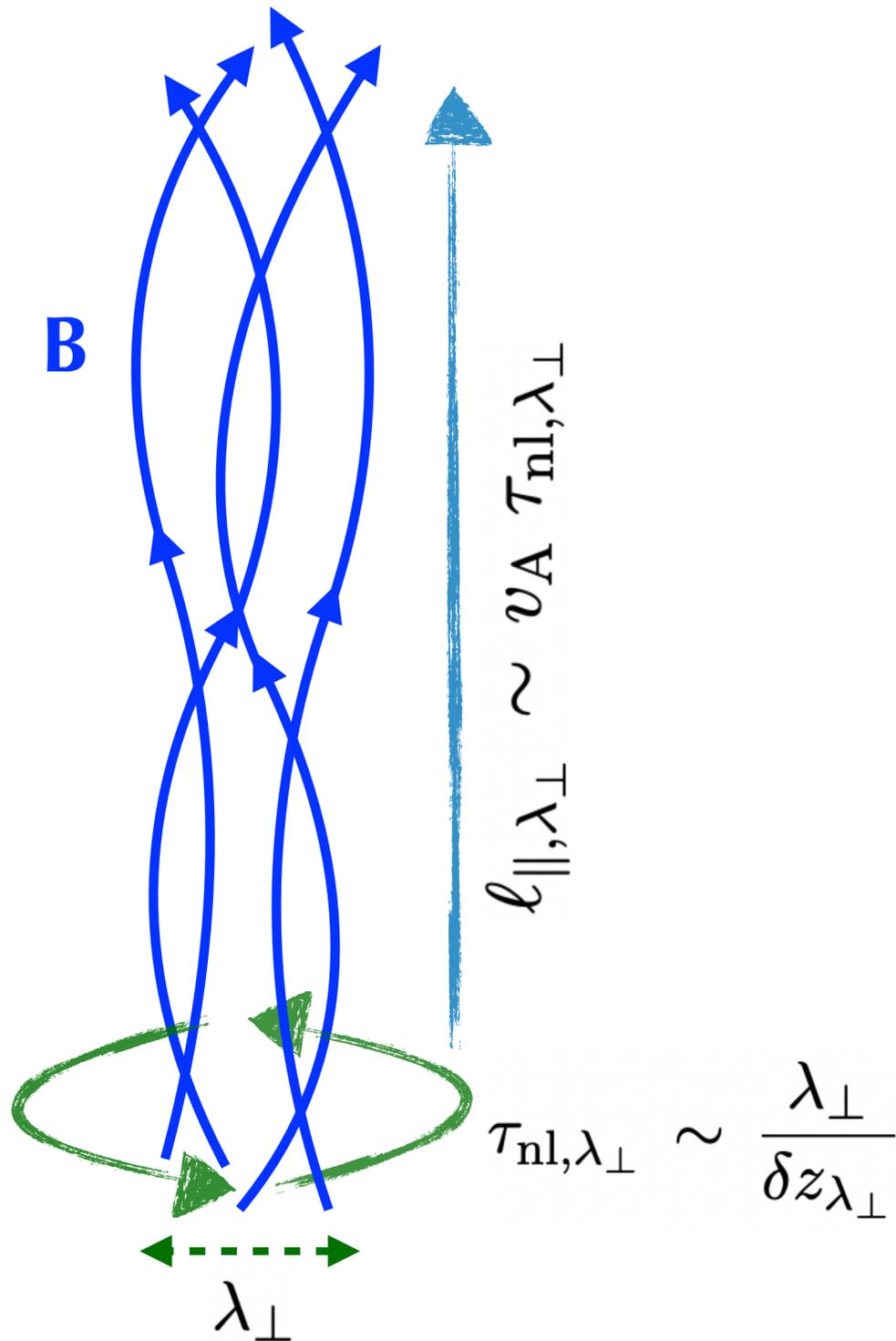
☞ At this point, linear, non-linear, and cascade timescales match each other:

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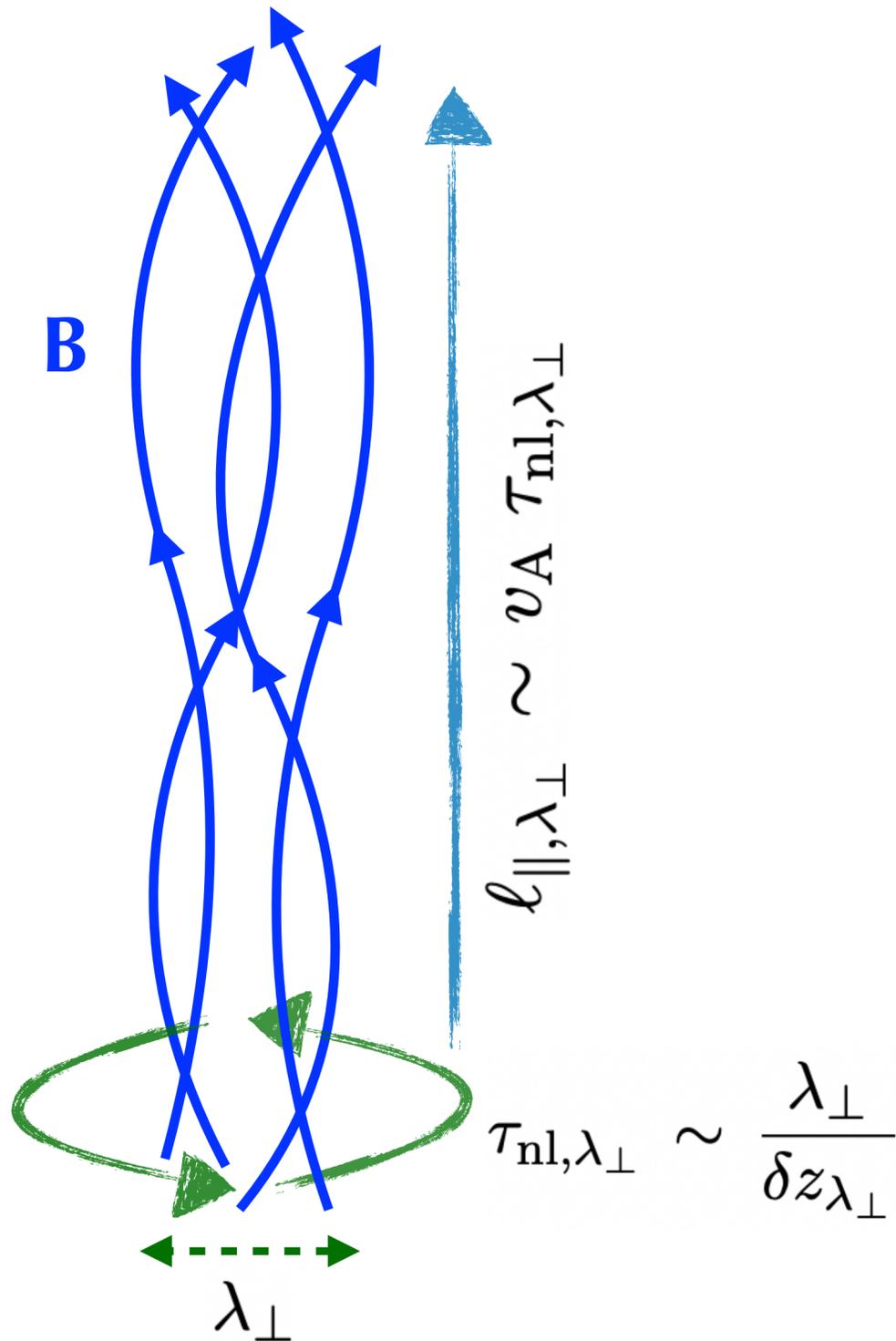
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☞ fluctuations' scaling + spectrum from $\varepsilon = \text{const.}$ (*you know the drill*):

$$\frac{\delta z_{k_{\perp}}^2}{\tau_{nl, k_{\perp}}} \sim \varepsilon = \text{const.} \Rightarrow \delta z_{k_{\perp}} \propto k_{\perp}^{-1/3} \Rightarrow \mathcal{E}_{\delta z}(k_{\perp}) \propto k_{\perp}^{-5/3}$$

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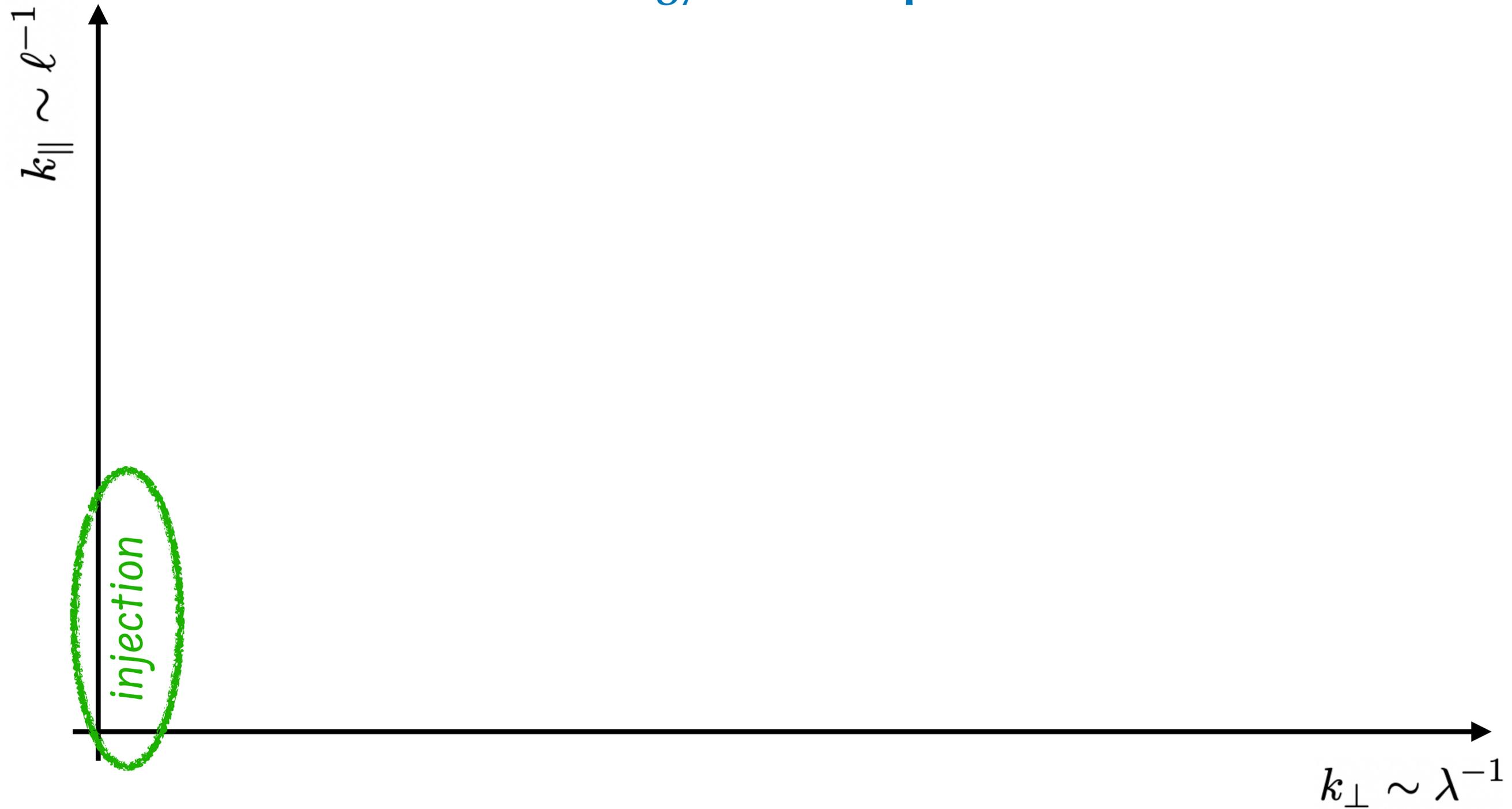
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☞ now, you can also compute the fluctuations' wavenumber anisotropy:

$$k_{\perp} \delta z_{k_{\perp}} \sim k_{\parallel} v_A \Rightarrow k_{\parallel} \propto k_{\perp}^{2/3} \left(\Rightarrow \mathcal{E}_{\delta z}(k_{\parallel}) \propto k_{\parallel}^{-2} \right)$$

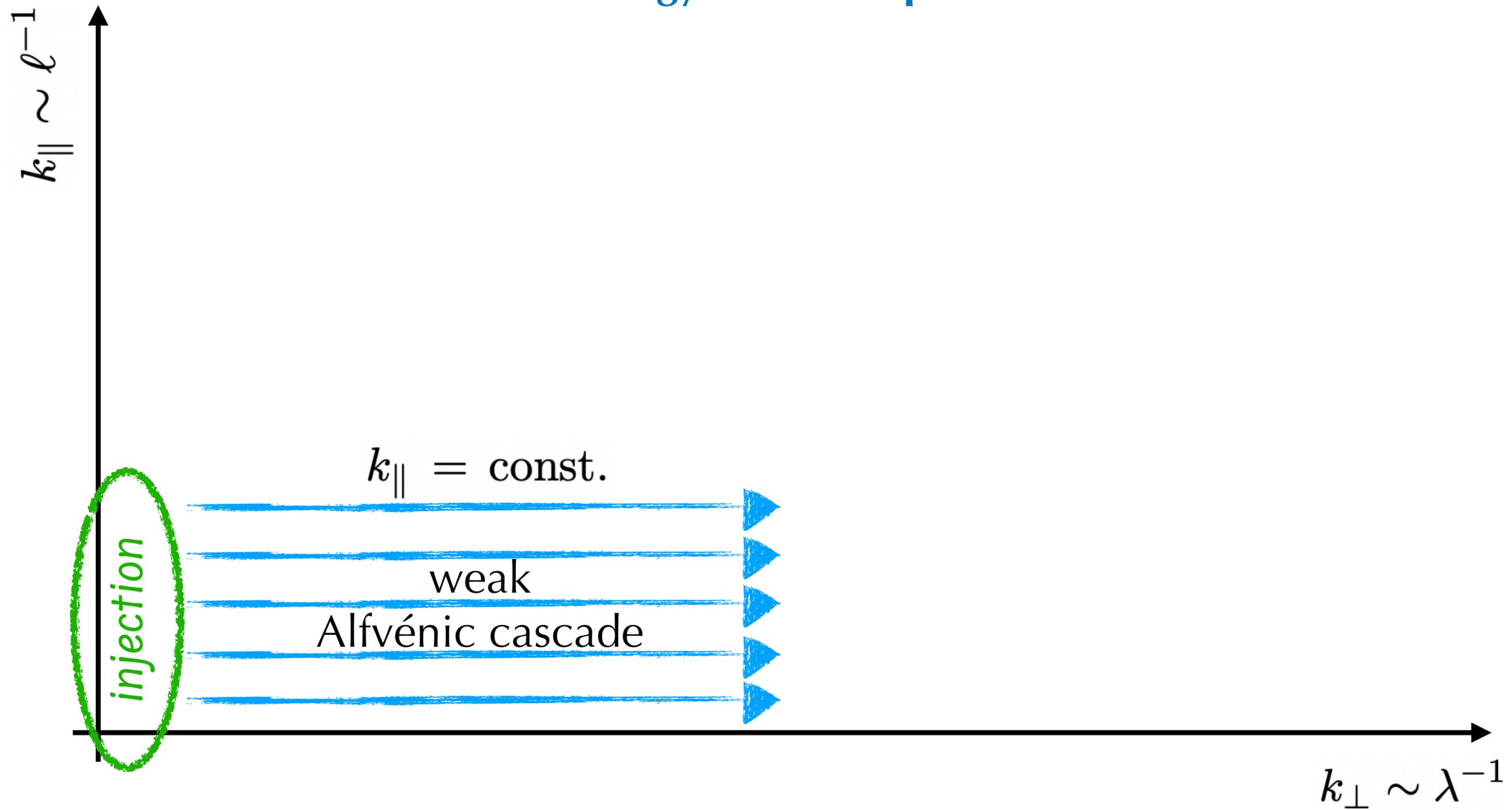
Phenomenology of Alfvénic Turbulence

Energy flux in k space



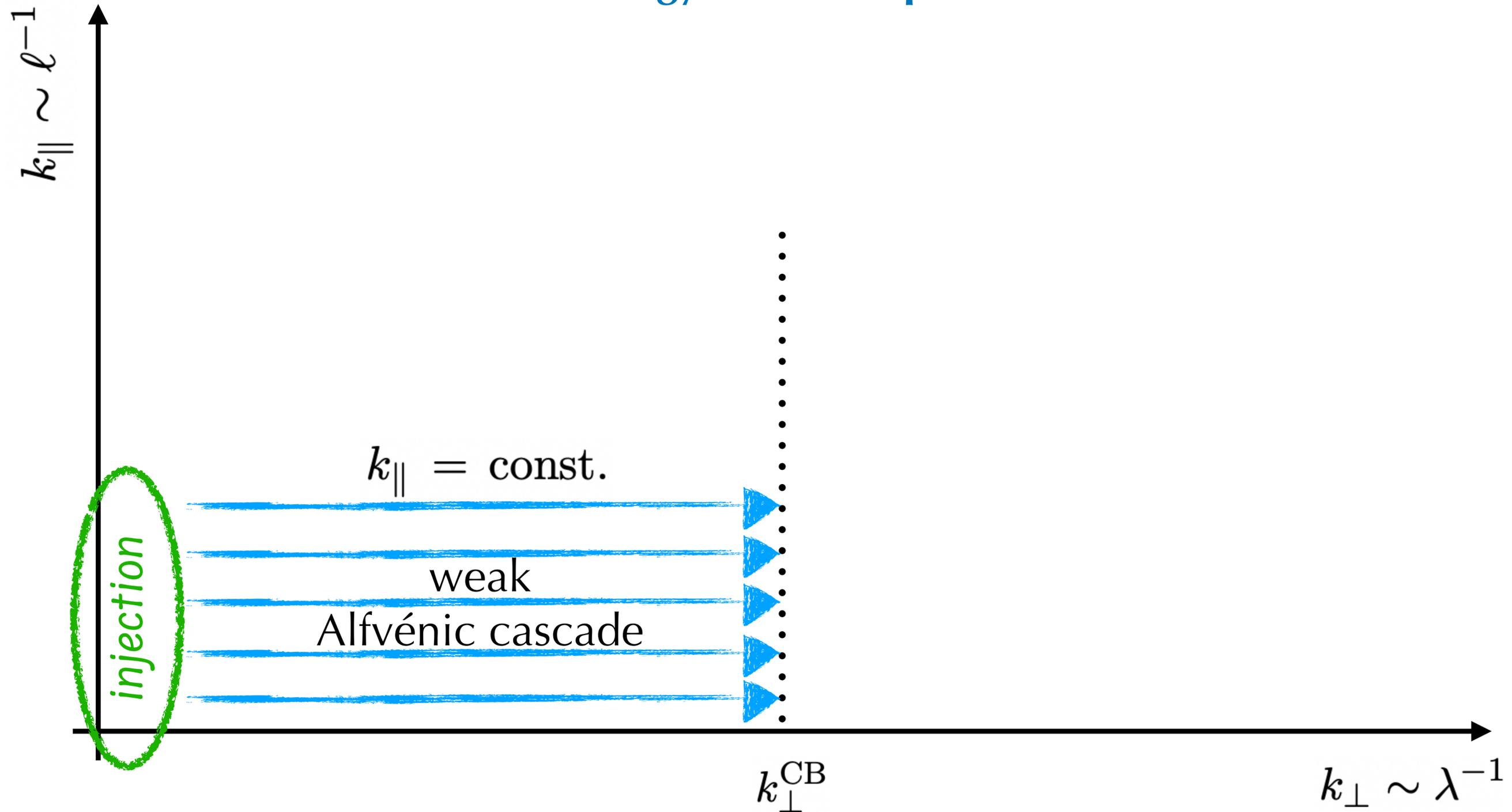
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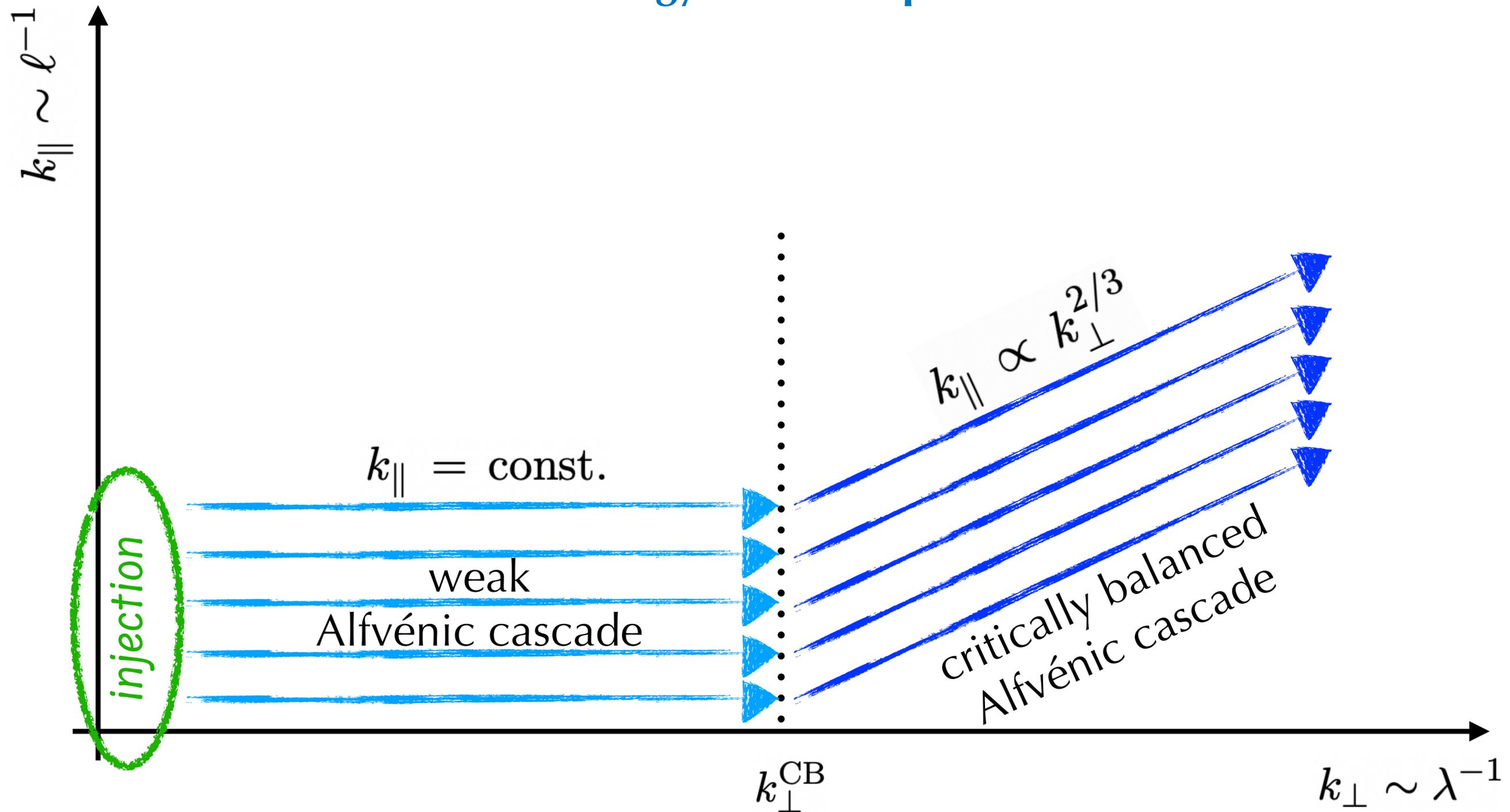
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Phenomenology of Alfvénic Turbulence

Energy flux in k space



Further Developments

Developments in Theoretical Models



Further Developments in Theoretical Models

dynamic alignment in Alfvénic turbulence: three-dimensional anisotropy

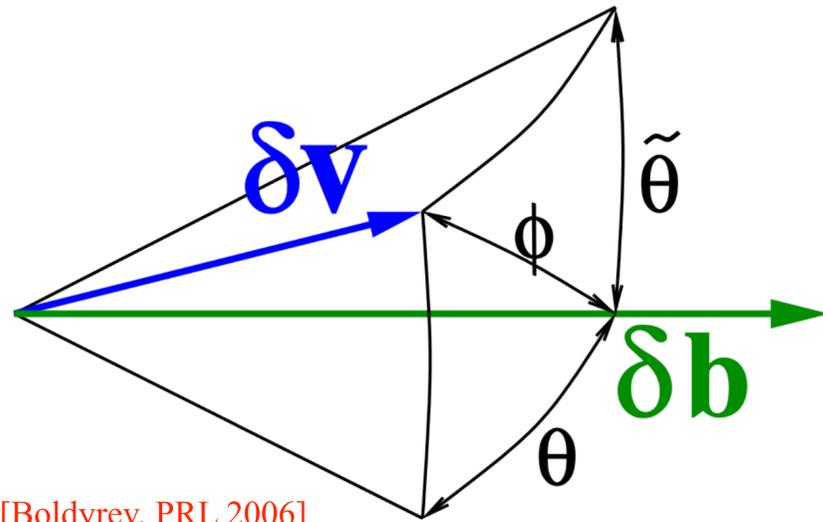
☞ *for further details, see, e.g.,*

[Boldyrev, PRL 2006]

[Schekochihin, arXiv:2010.00699]

Further Developments in Theoretical Models

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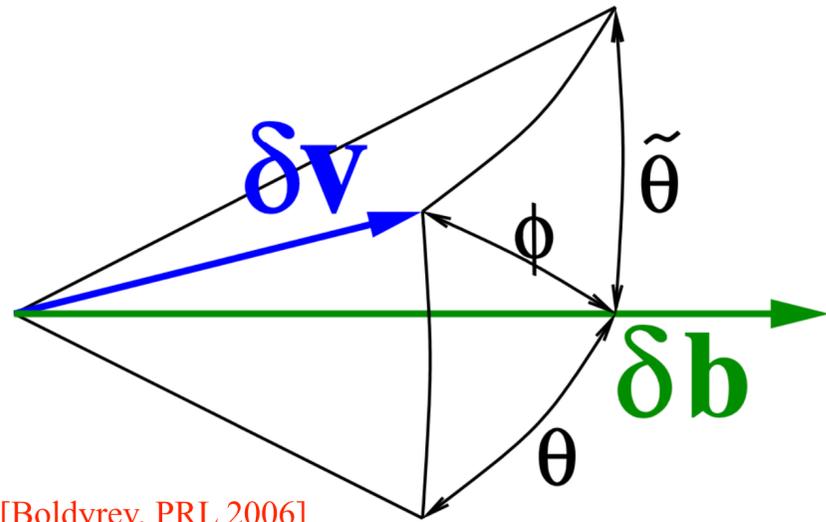


[Boldyrev, PRL 2006]

- Observations and simulations show that $\delta \mathbf{v}_\lambda$ and $\delta \mathbf{b}_\lambda$ have a spontaneous tendency to align in the plane perpendicular to the local mean field $\langle \mathbf{B} \rangle_\lambda$, within an angle θ_λ
(e.g., Podesta et al., JGR 2009; Hnat et al., PRE 2011; Mason et al., ApJ 2011; Wicks et al., PRL 2013; Mallet et al., MNRAS 2016; ...)

Further Developments in Theoretical Models

dynamic alignment in Alfvénic turbulence: three-dimensional anisotropy



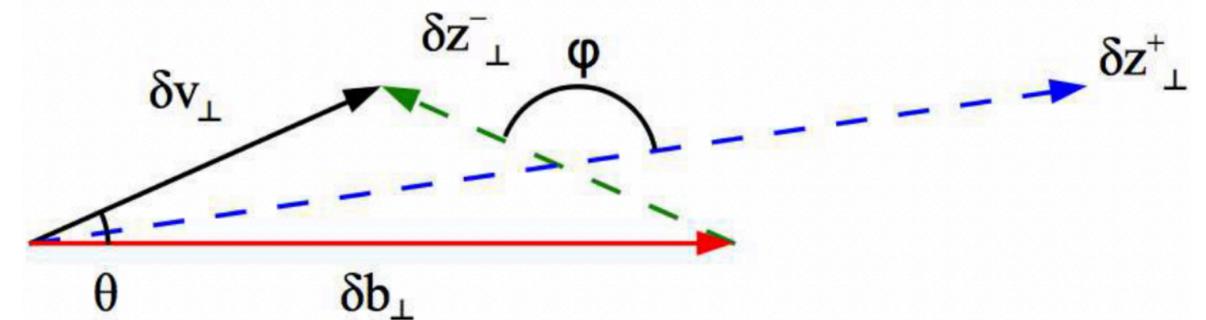
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! the alignment between $\delta \mathbf{v}_\lambda$ and $\delta \mathbf{b}_\lambda$ is *not the same* as the alignment between $\delta \mathbf{z}^+_\lambda$ and $\delta \mathbf{z}^-_\lambda$!

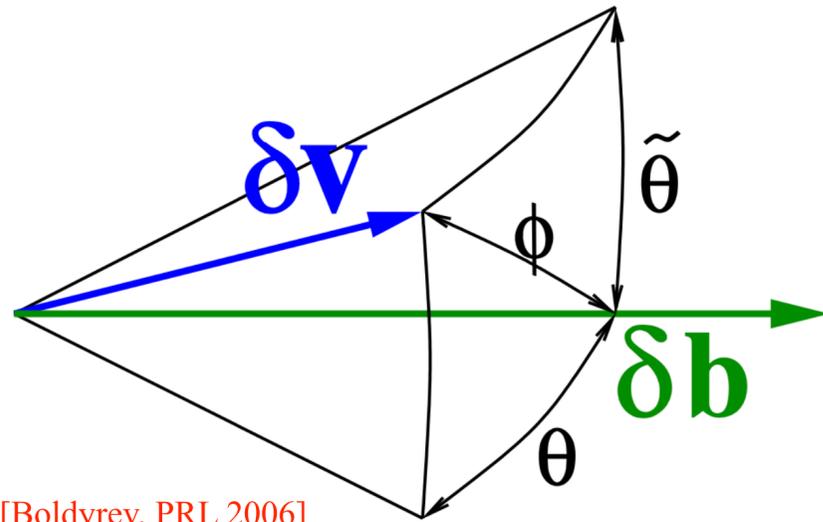
(but they are related: see Schekochihin arXiv:2010.00699)



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Further Developments in Theoretical Models

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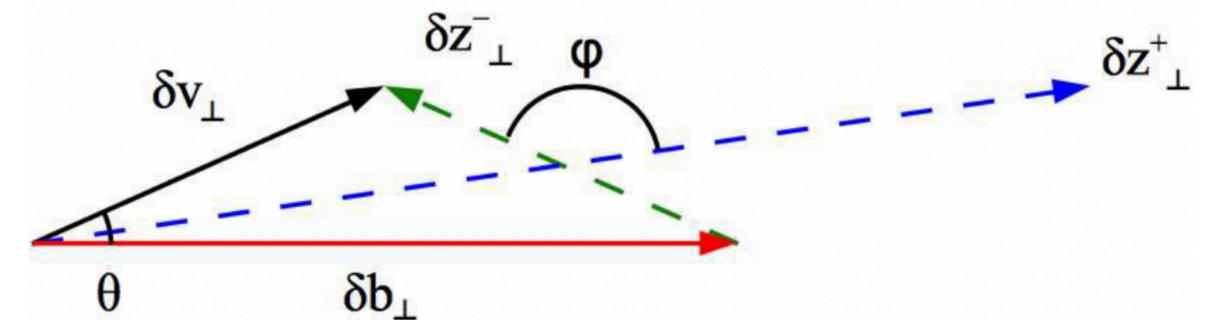
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Observations and simulations show that $\delta\mathbf{v}_\lambda$ and $\delta\mathbf{b}_\lambda$ have a spontaneous tendency to align in the plane perpendicular to the local mean field $\langle\mathbf{B}\rangle_\lambda$, within an angle θ_λ

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[Wicks et al., PRL 2013]

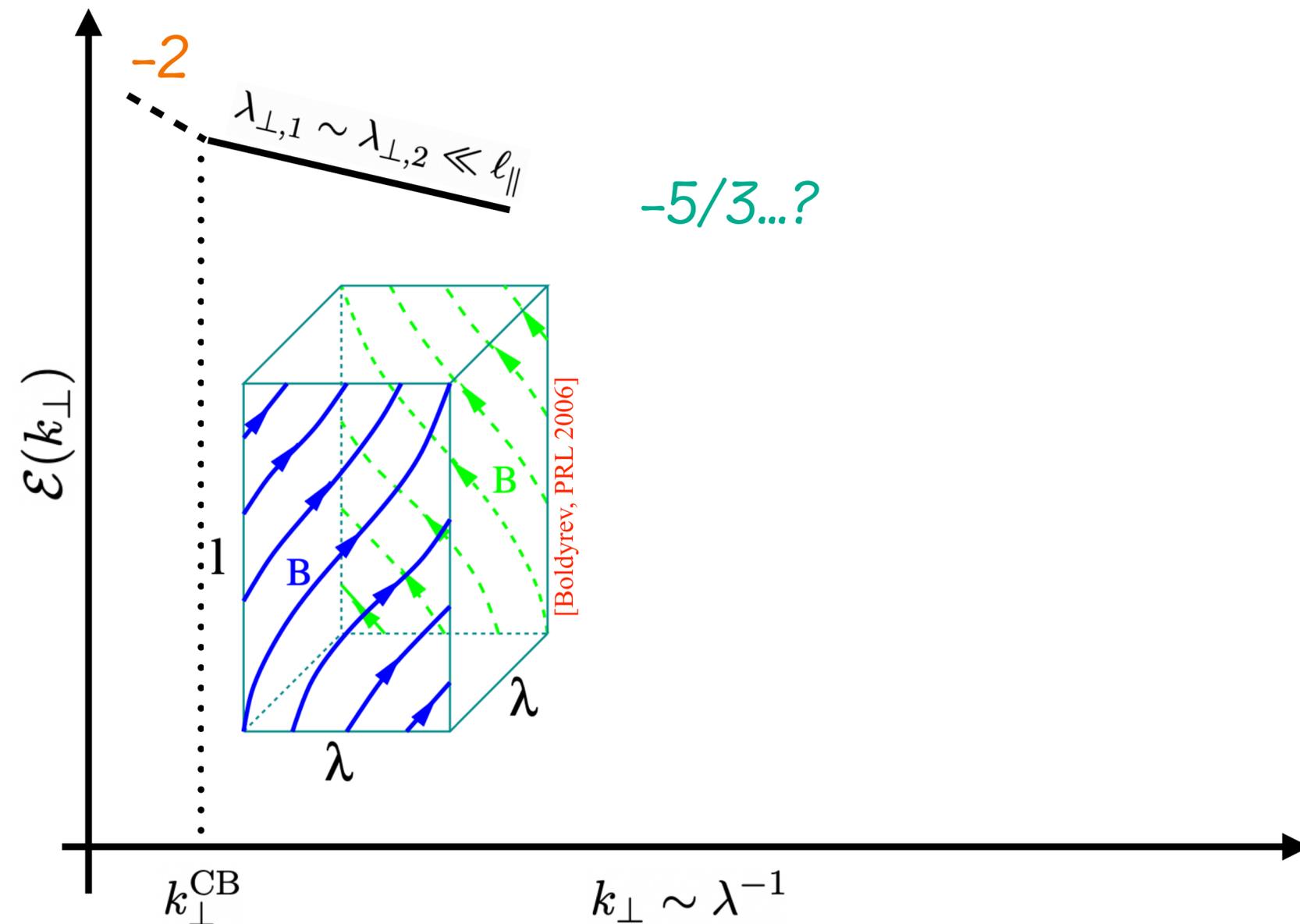
alignment \Rightarrow depletion of non-linearities: $\delta\mathbf{z}^\mp \cdot \nabla \delta\mathbf{z}^\pm \sim \sin \varphi_\lambda \frac{\delta z_\lambda^2}{\lambda} \approx \varphi_\lambda \frac{\delta z_\lambda^2}{\lambda} \longleftrightarrow \theta_\lambda \frac{\delta v_\lambda^2}{\lambda}$

⚠ but remember that *fluctuations cannot be perfectly aligned* ($\theta_\lambda = 0$) in order to have a non-linear cascade

Further Developments in Theoretical Models

dynamic alignment in Alfvénic turbulence: three-dimensional anisotropy

The effect of alignment is not only to make the non-linear interactions weaker, but also to induce anisotropy in the plane perpendicular to the magnetic field \mathbf{B}

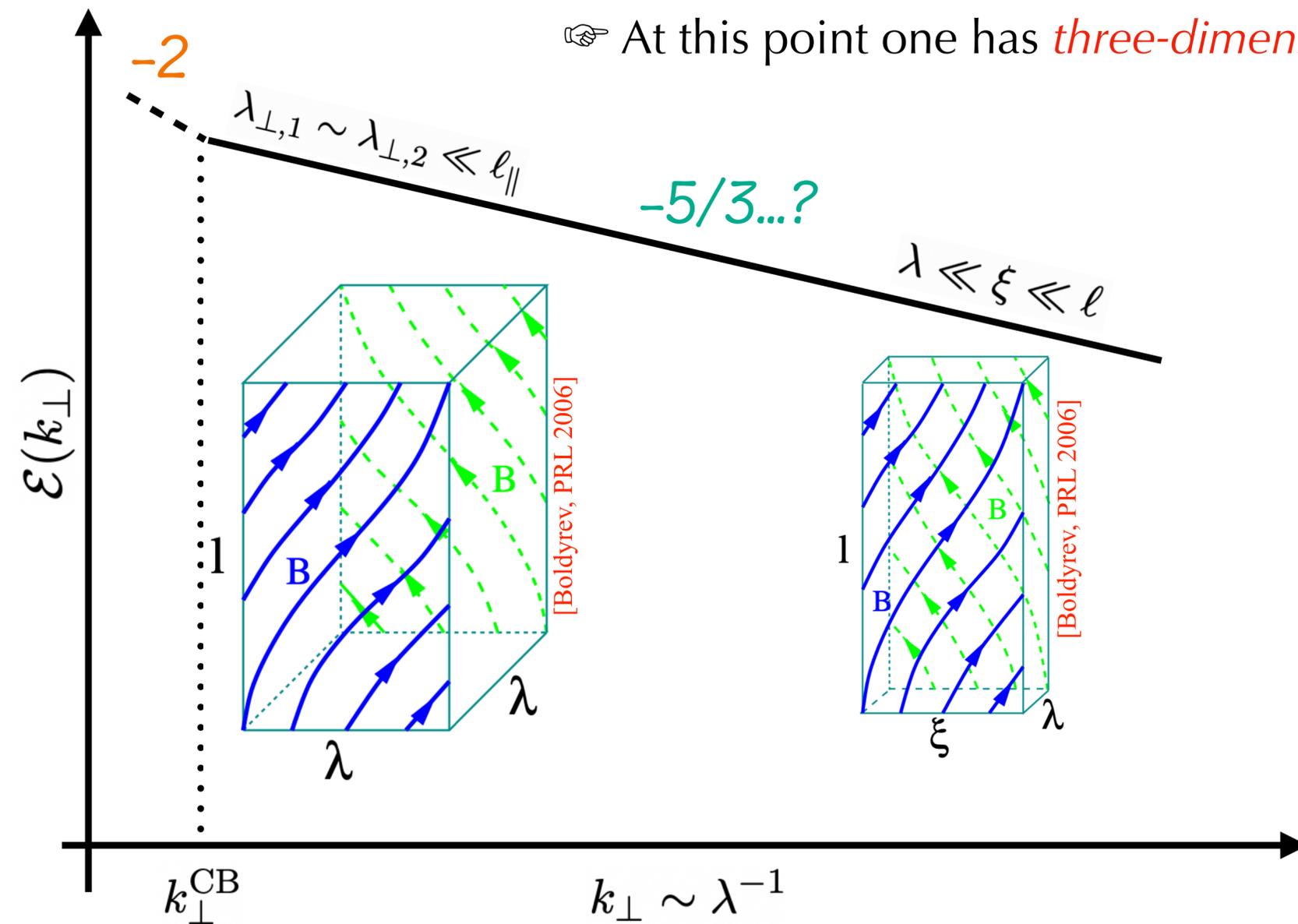


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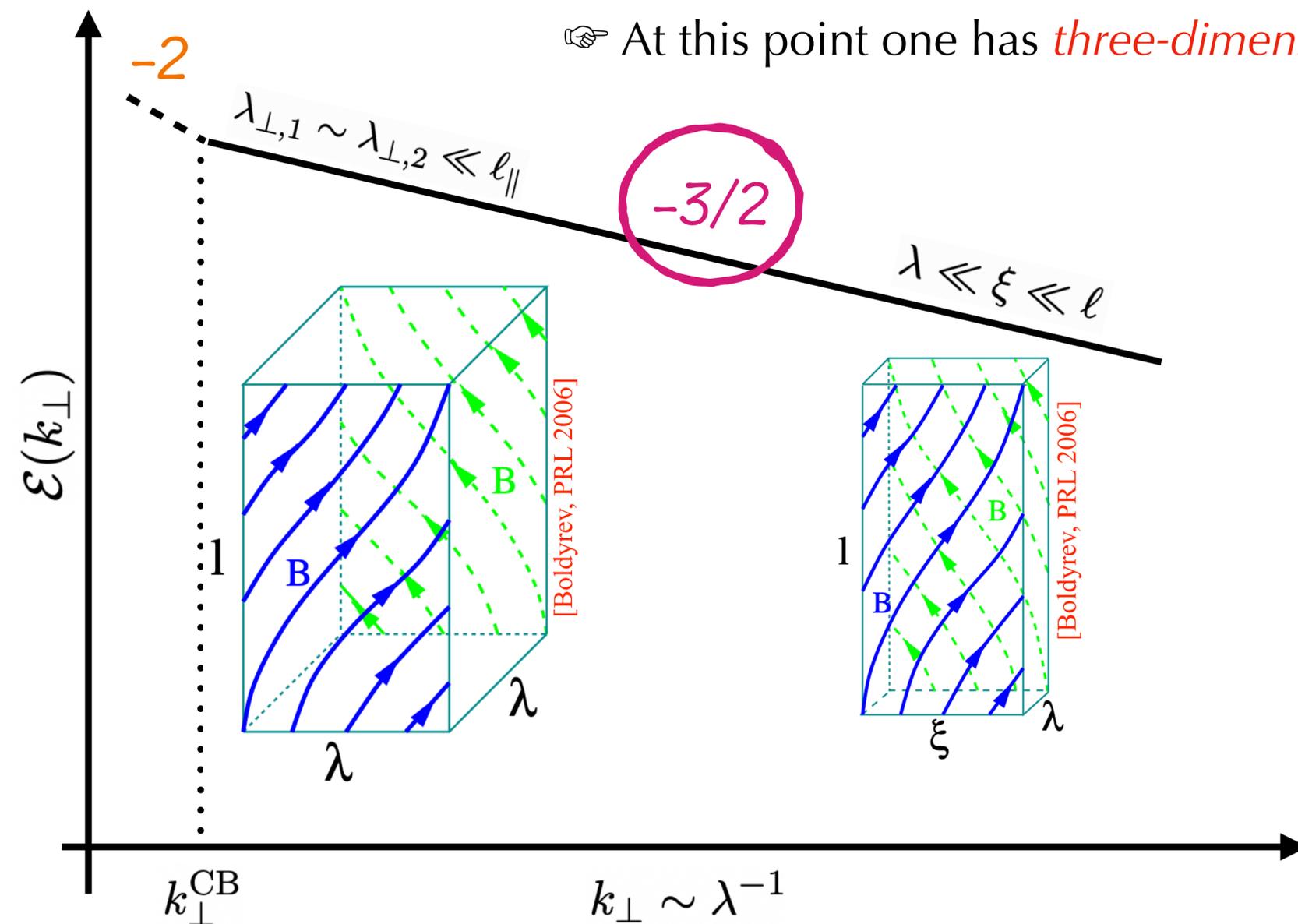


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long story short:

(see Boldyrev, PRL 2006 for the derivation)

$$\theta_{k_{\perp}} \propto k_{\perp}^{-1/4} \Rightarrow \delta v_{k_{\perp}} \propto k_{\perp}^{-1/4} \Rightarrow \mathcal{E}(k_{\perp}) \propto k_{\perp}^{-3/2}$$

(also, now $k_{\parallel} \propto k_{\perp}^{1/2}$)

Recent Developments in Theoretical Models

reconnection-mediated regime in Alfvénic turbulence

☞ *for further details, see, e.g.,*

[Boldyrev & Loureiro, ApJ 2017]

[Mallet, Schekochihin, Chandran, MNRAS 2017]

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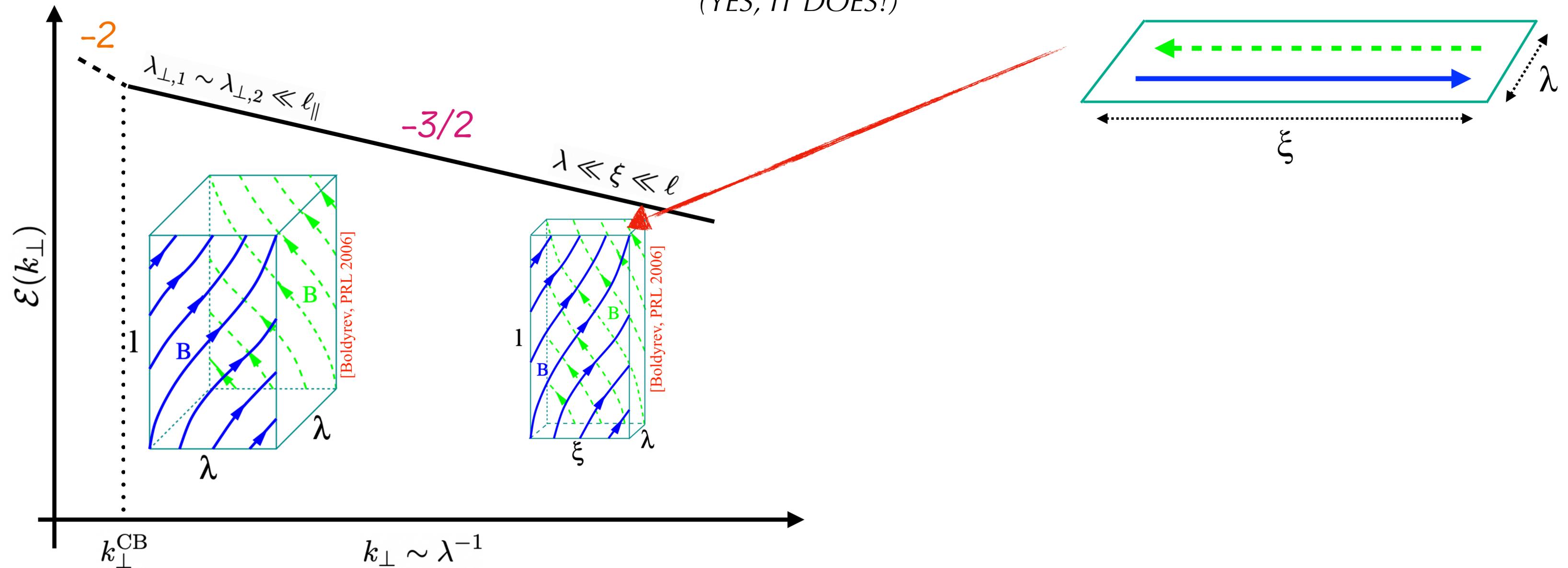
Recent Developments in Theoretical Models

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So, we had *three-dimensional anisotropy*, right? ... wait a minute!

doesn't 3D anisotropy of the turbulent eddies *look like a current sheet in the plane perpendicular to B ?*

(YES, IT DOES!)



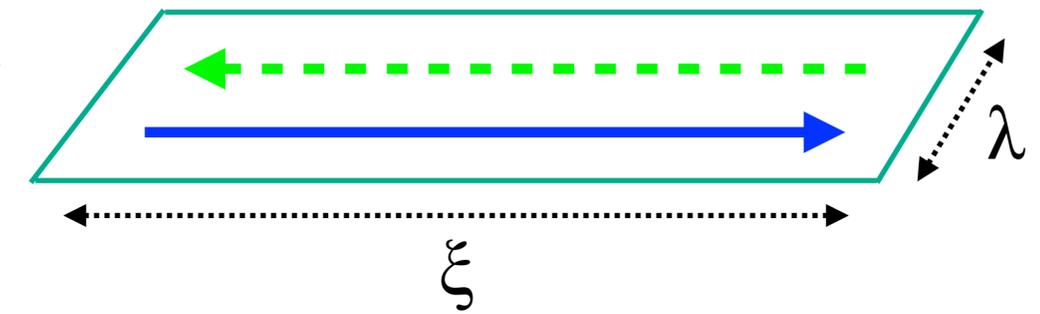
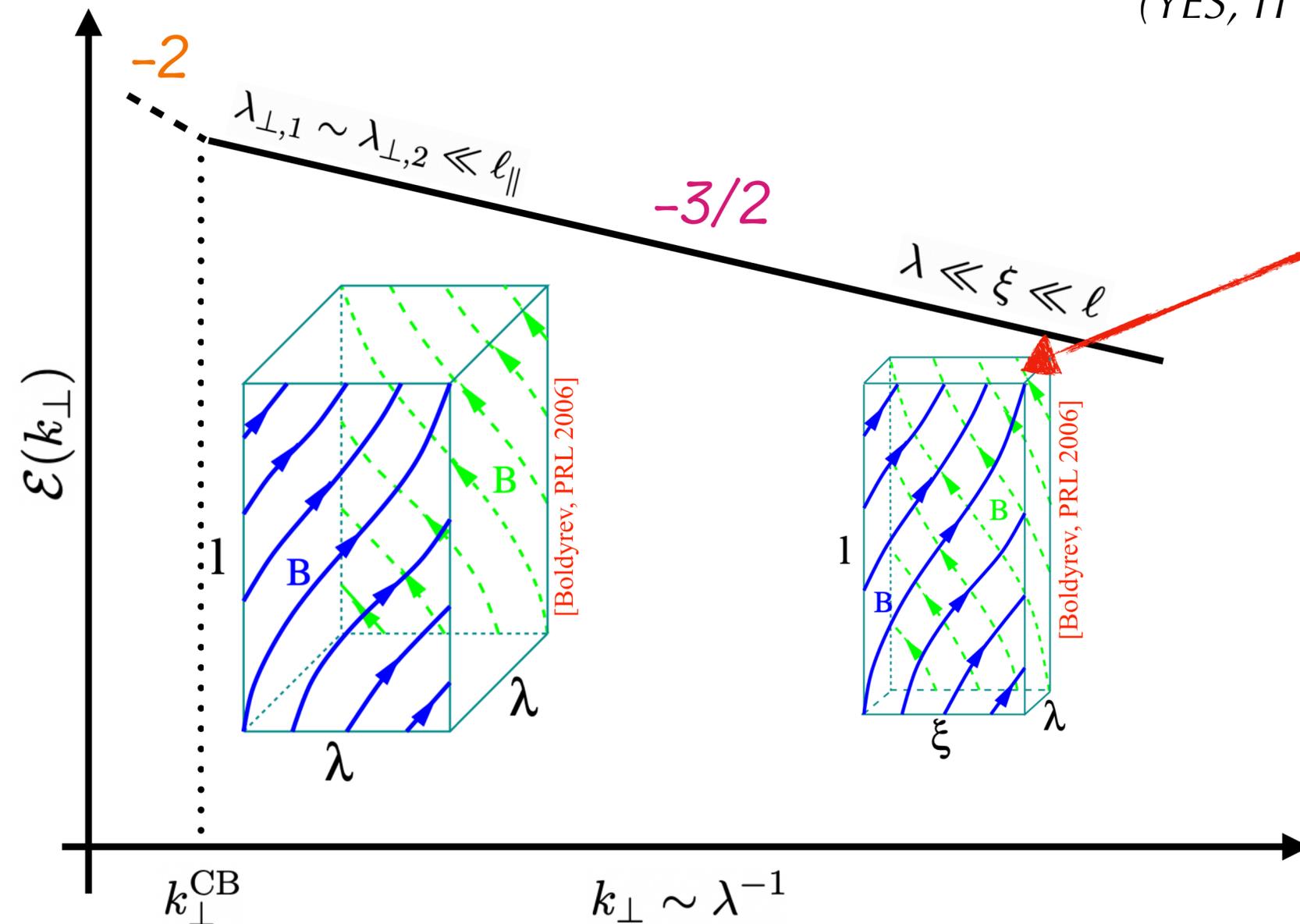
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☞ if the eddies at a scale live “long enough” for the tearing instability (i.e., reconnection) to grow, then we can imagine that this process will be responsible for the production of small-scale magnetic fluctuations

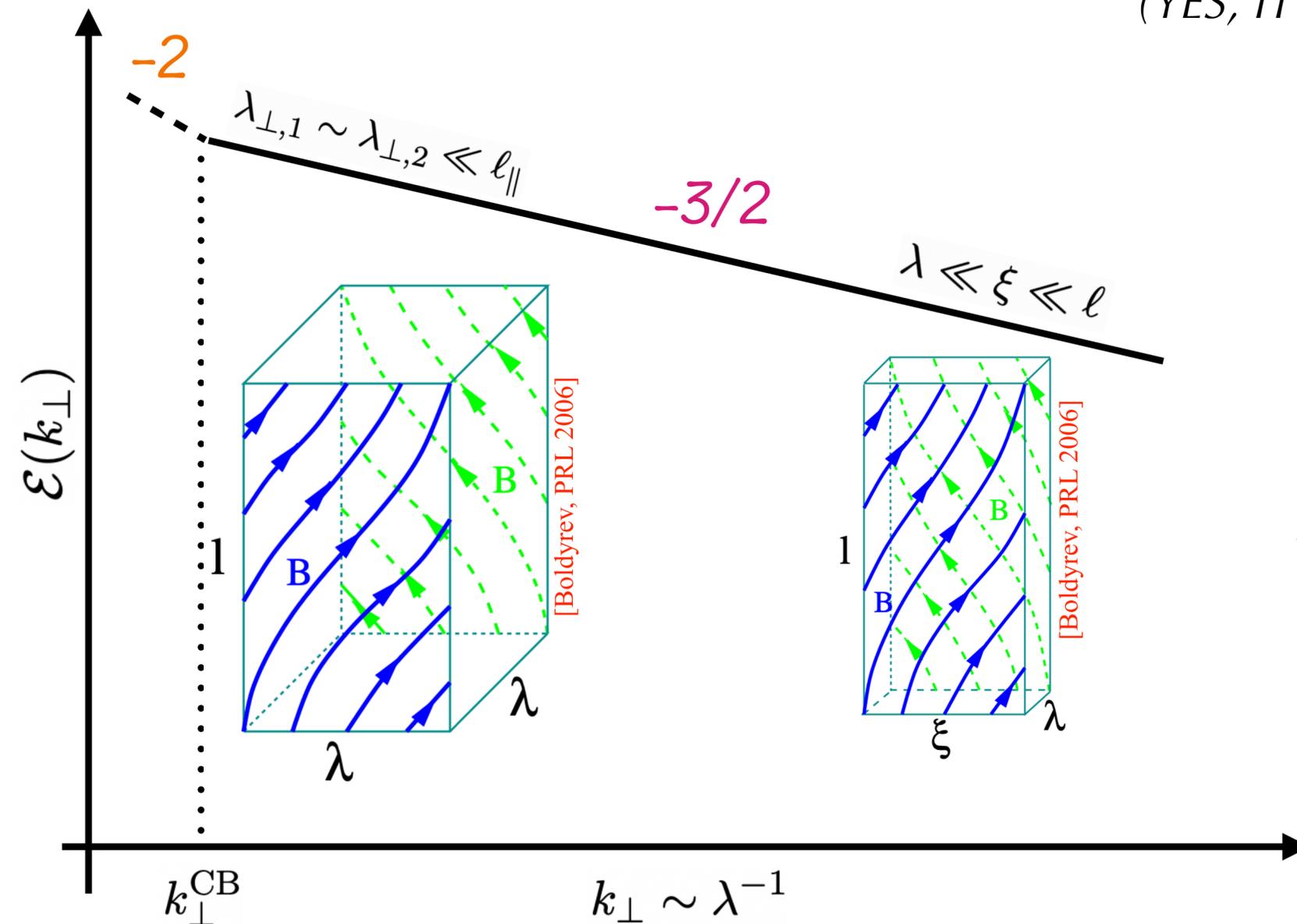
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eddy lifetime:

$$\tau_{\text{nl},k_{\perp}} \sim (\theta_{k_{\perp}} k_{\perp} \delta v_{k_{\perp}})^{-1} \propto k_{\perp}^{-1/2}$$

tearing growth rate:

$$\gamma_{k_{\perp}}^{\text{rec}} \sim k_{\perp} \delta v_{k_{\perp}} \left(\frac{\delta v_{k_{\perp}}}{k_{\perp} \eta} \right)^{-1/2} \propto S_0^{-1/2} k_{\perp}^{11/8}$$

$$\left(\eta : \text{resistivity}, \quad S \doteq \frac{v_A \ell_0}{\eta} : \text{Lundquist number} \right)$$

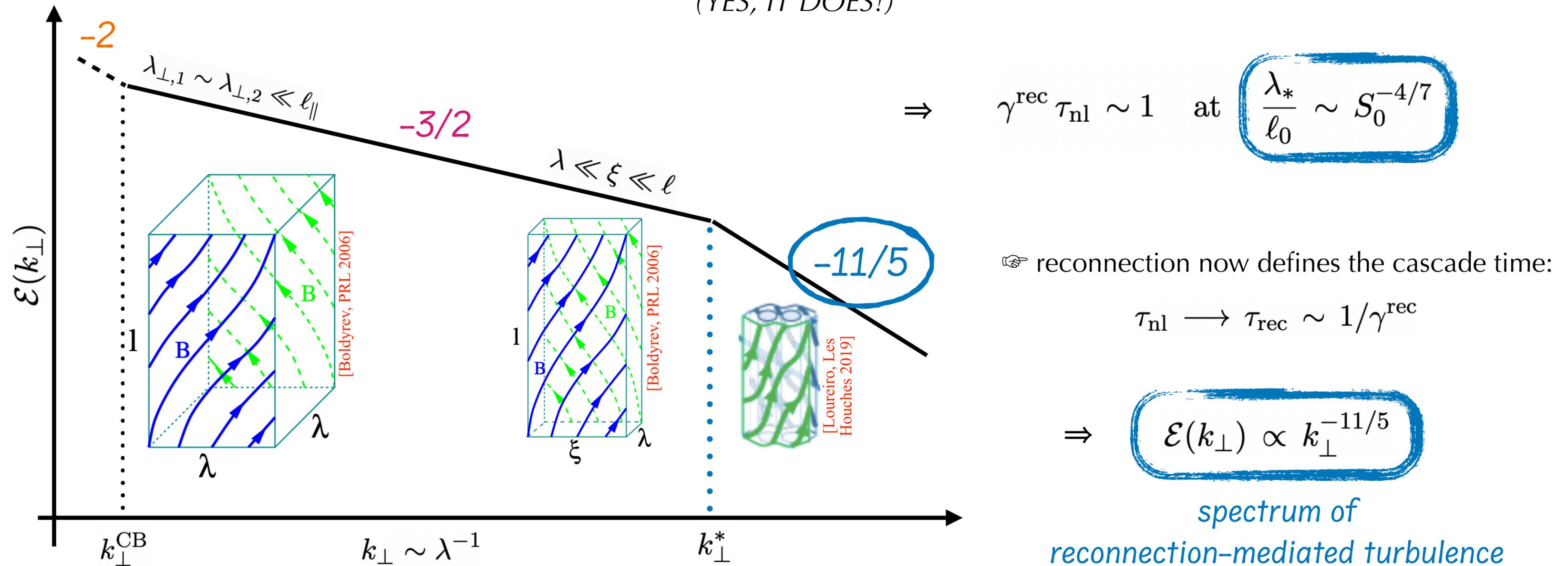
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Recent Developments in Theoretical Models

Is this the end of the story?!

Of course not! ... we have totally neglected **KINETIC EFFECTS...**

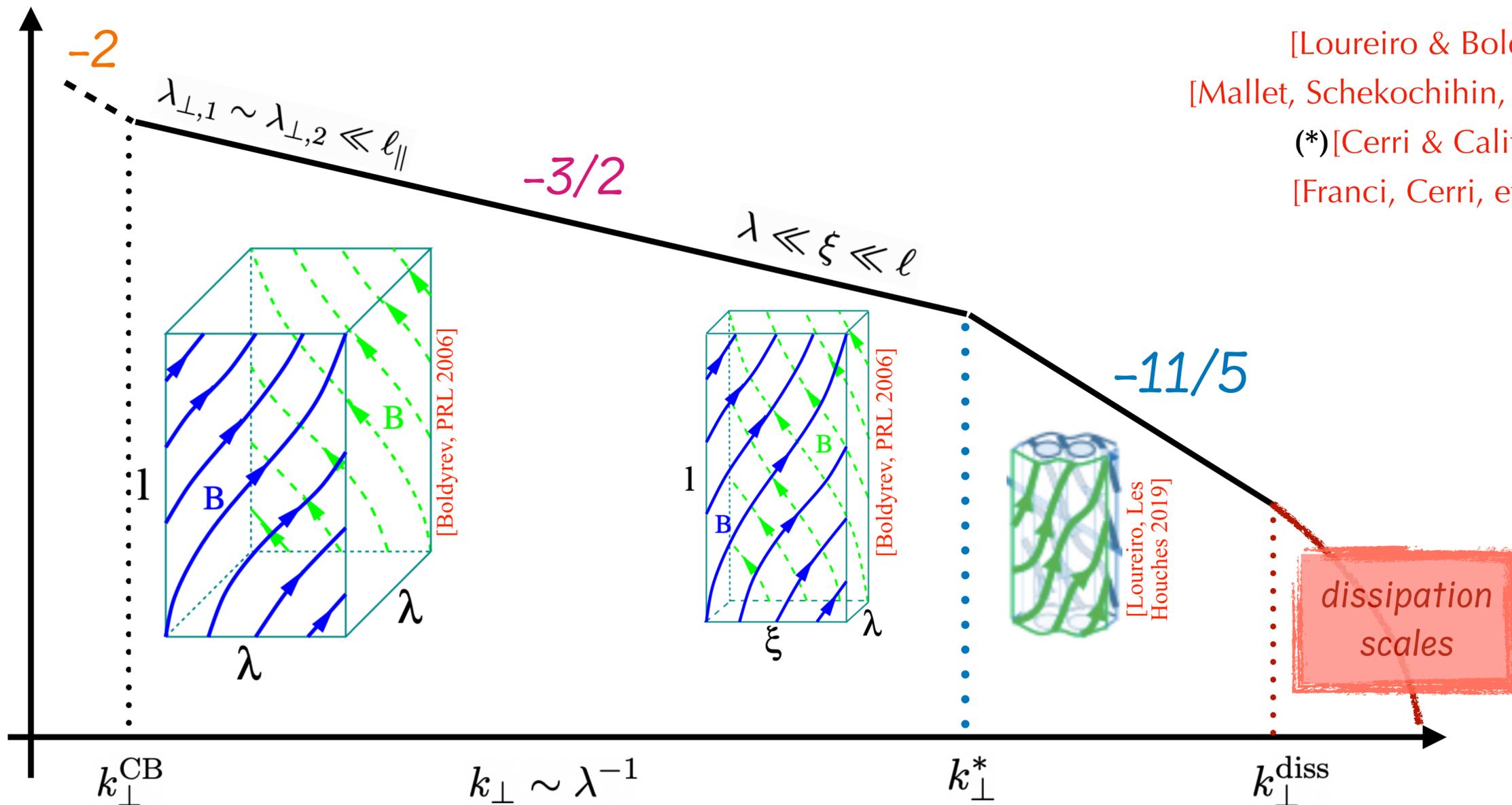
☞ reconnection-mediated turbulence with kinetic effects (theory & simulations):

[Loureiro & Boldyrev, ApJ 2017]

[Mallet, Schekochihin, Chandran, JPP 2017]

(*) [Cerri & Califano, NJP 2017]

[Franci, Cerri, et al., ApJL 2017]



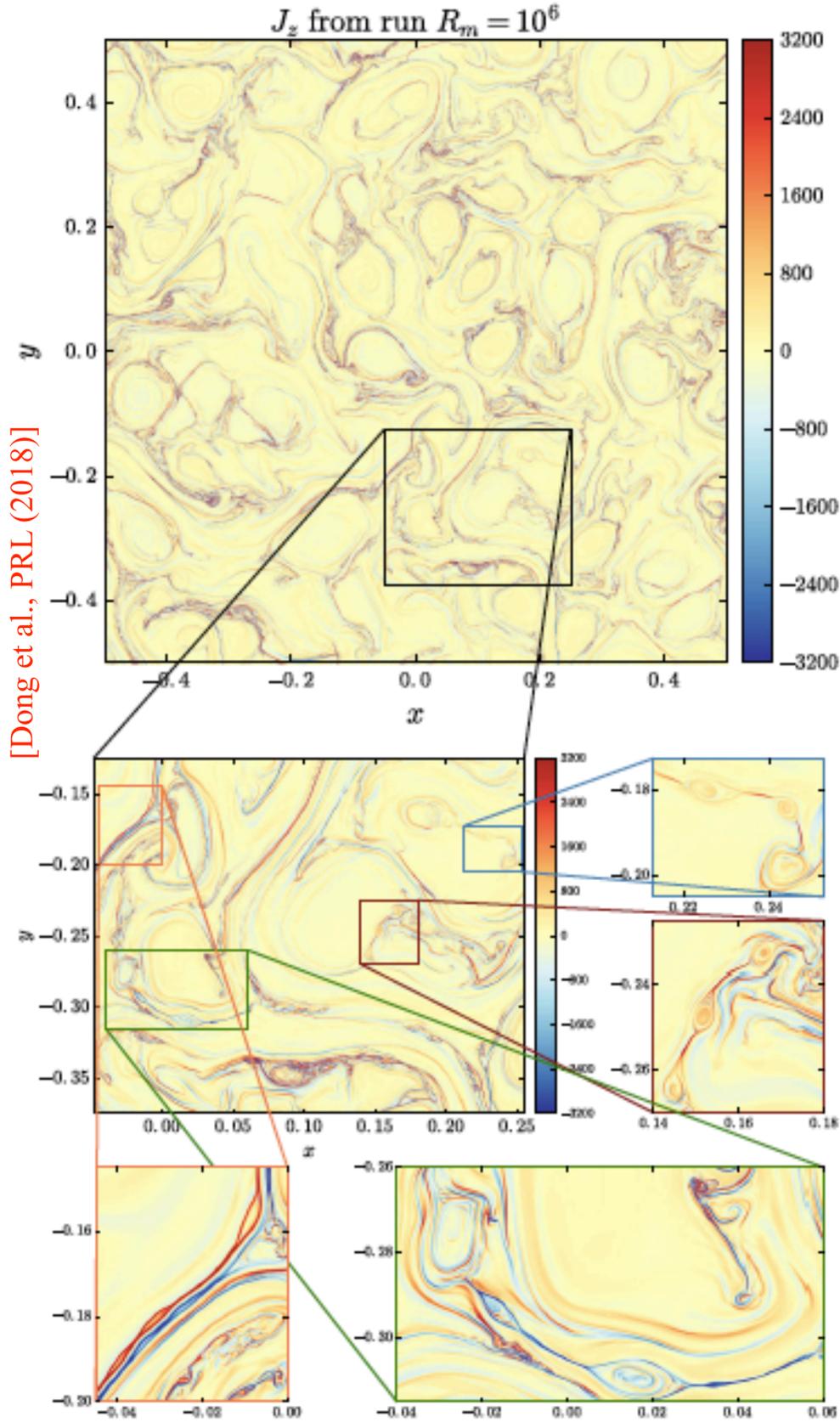
(*) Actually this was the first suggestion of the existence of a reconnection-mediated regime: from a kinetic simulation, before any theory existed!

Progress via Numerical Simulations



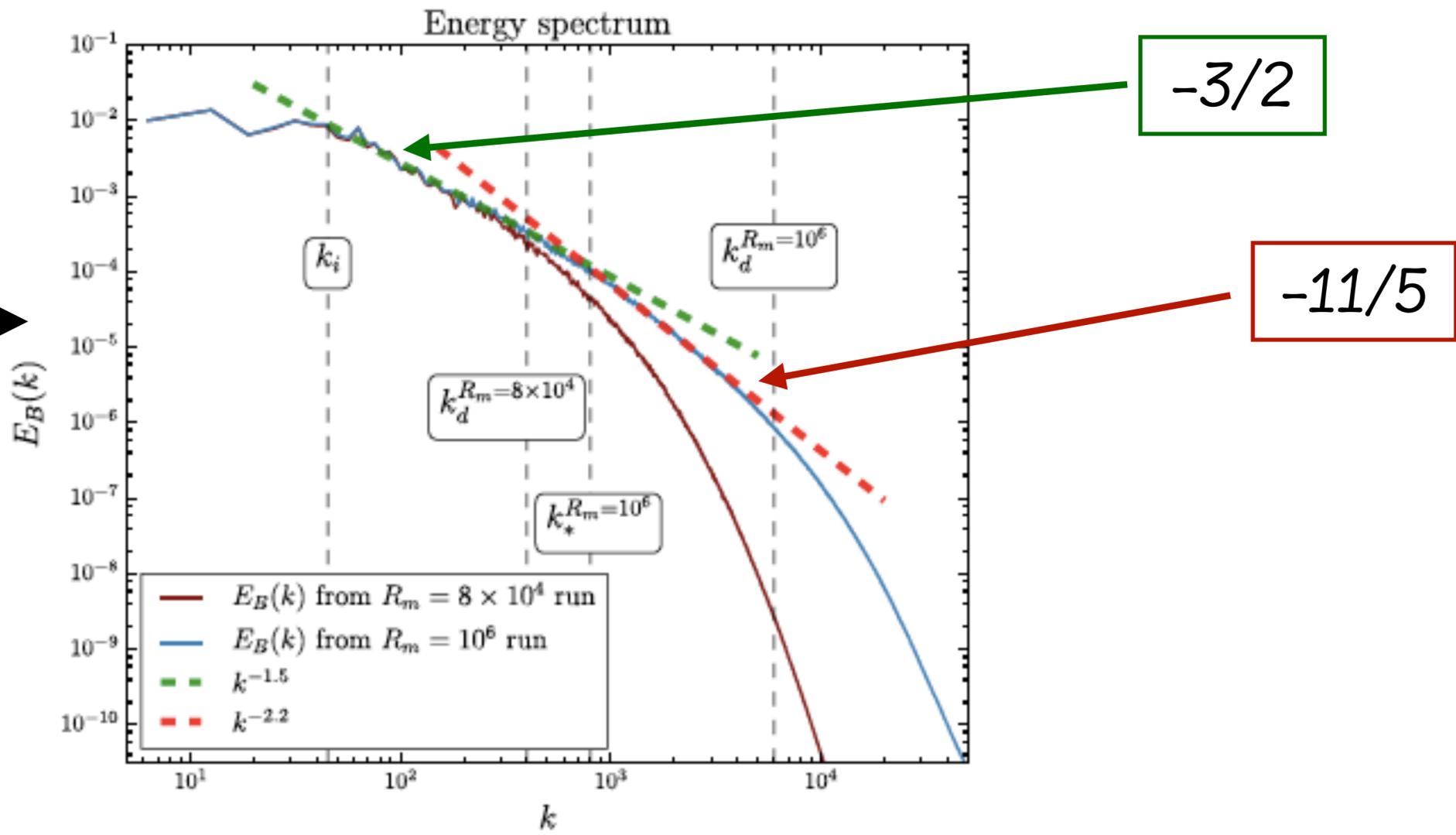
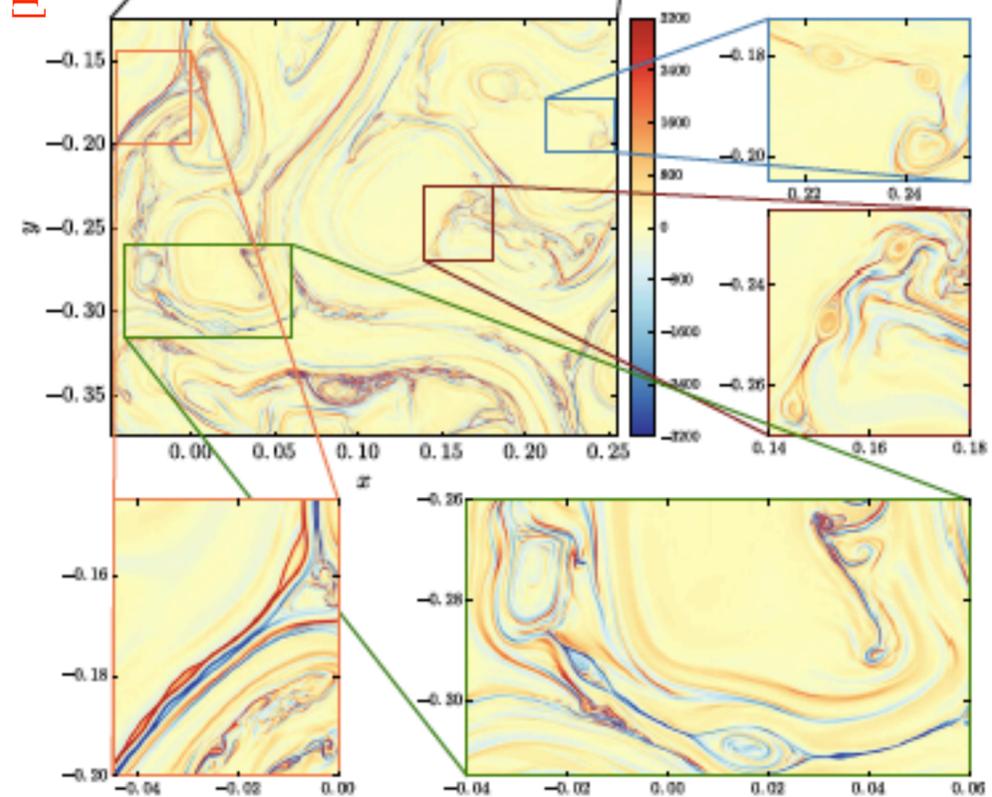
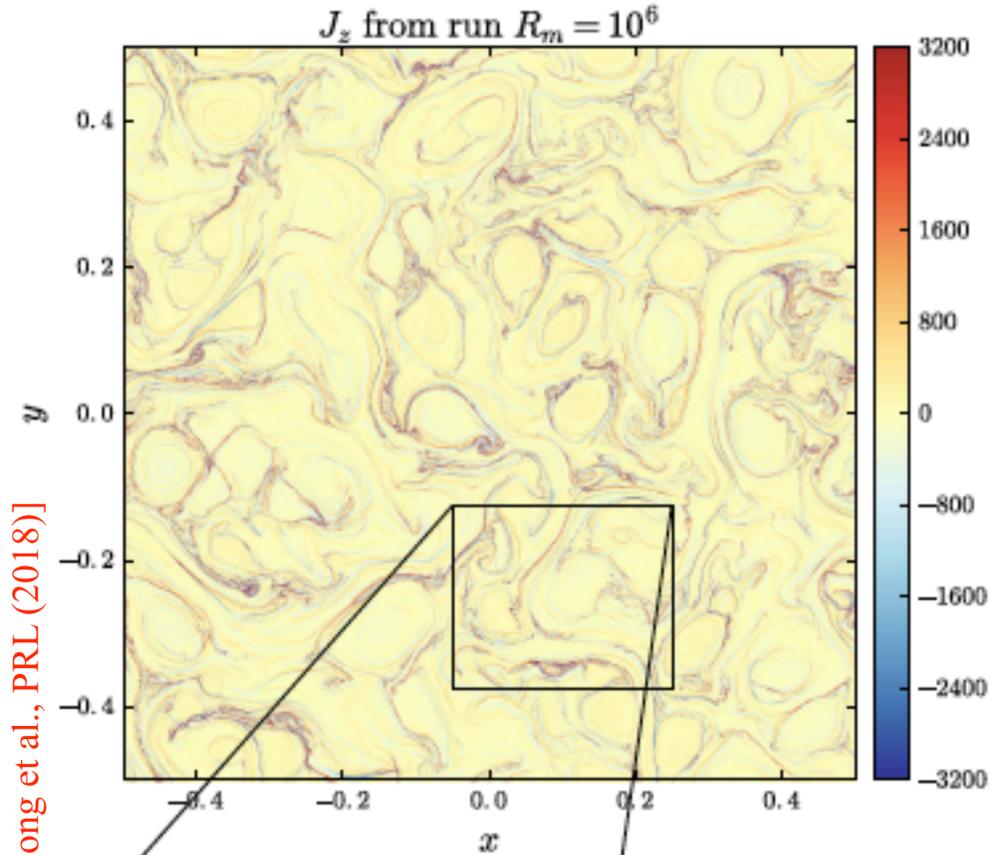
Recent Developments via Numerical Simulations

[Dong et al., PRL (2018)]



Recent Developments via Numerical Simulations

[Dong et al., PRL (2018)]



so far, the only evidence of reconnection-mediated turbulence in MHD

- only in **2D** geometry
- requires *extremely large Lundquist numbers* (grid: 64000^2 !!!)

Recent Developments via Numerical Simulations

Can we do better, and can it be done in 3D?

Recent Developments via Numerical Simulations

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just go back to a **basic 3D setup**: start from the ***building blocks of the Alfvénic cascade!***

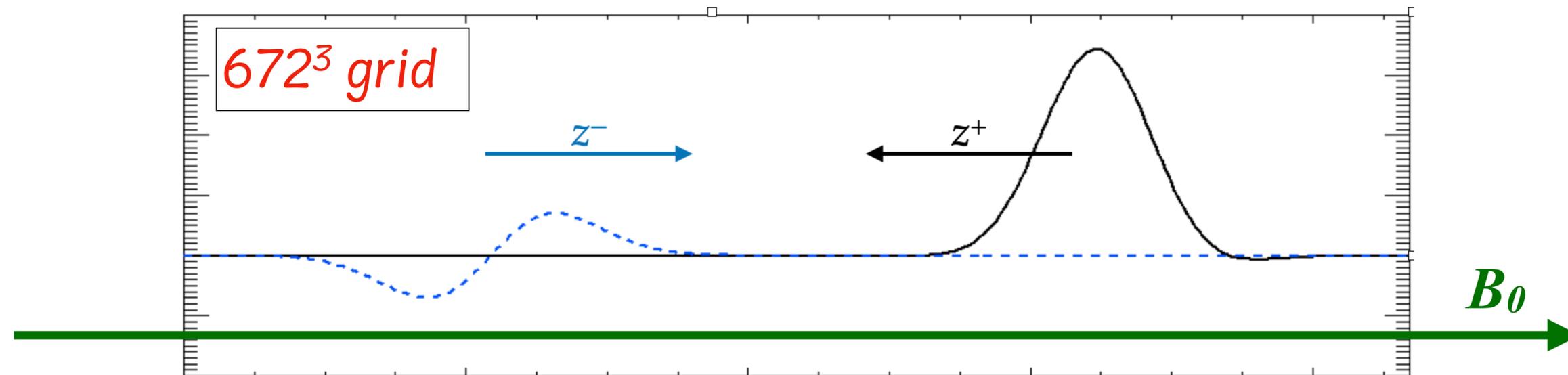
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Simulations performed with the *Hamiltonian 2-fields gyro-fluid* model/code by Passot, Tassi, Sulem, and Laveder

👉 model retains *only Alfvén & kinetic-Alfvén modes*, assumes *strong anisotropy* ($k_{\parallel} \ll k_{\perp}$), ...

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But we do it “WISELY”, i.e., with a “trick”:

$$\gamma^{\text{rec}} \tau_{\text{nl}} \sim 1$$

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Usually, one tries to increase γ^{rec} by achieving large S : requires extreme resolution!

Recent Developments via Numerical Simulations

[Cerri, Passot, Laveder, Sulem, Kunz, ApJL (to be submitted)]

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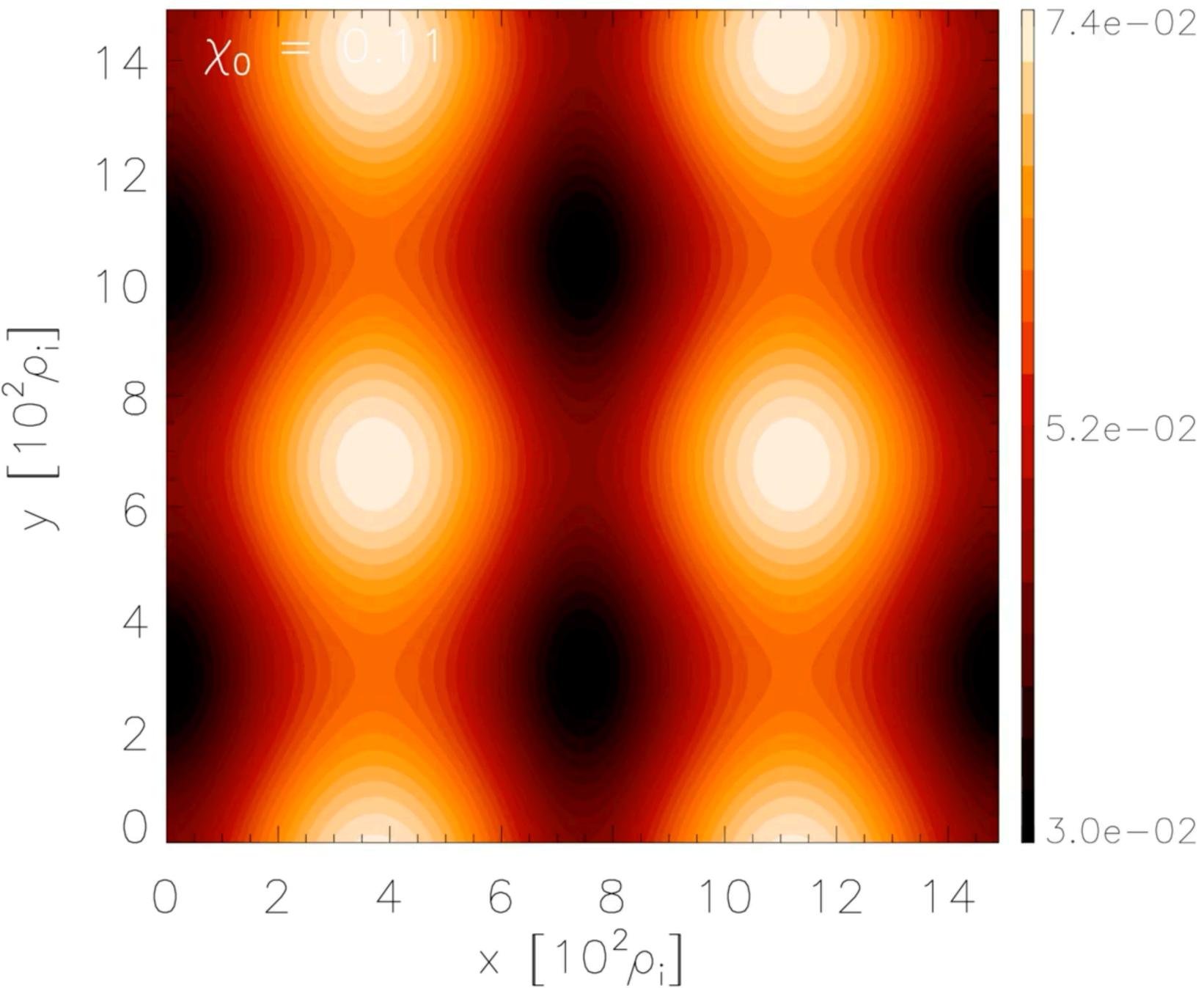
$$\gamma^{\text{rec}} \tau_{\text{nl}} \approx 1$$

Let's increase the non-linear time instead! (by considering a smaller non-linear parameter, $\chi < 1$)

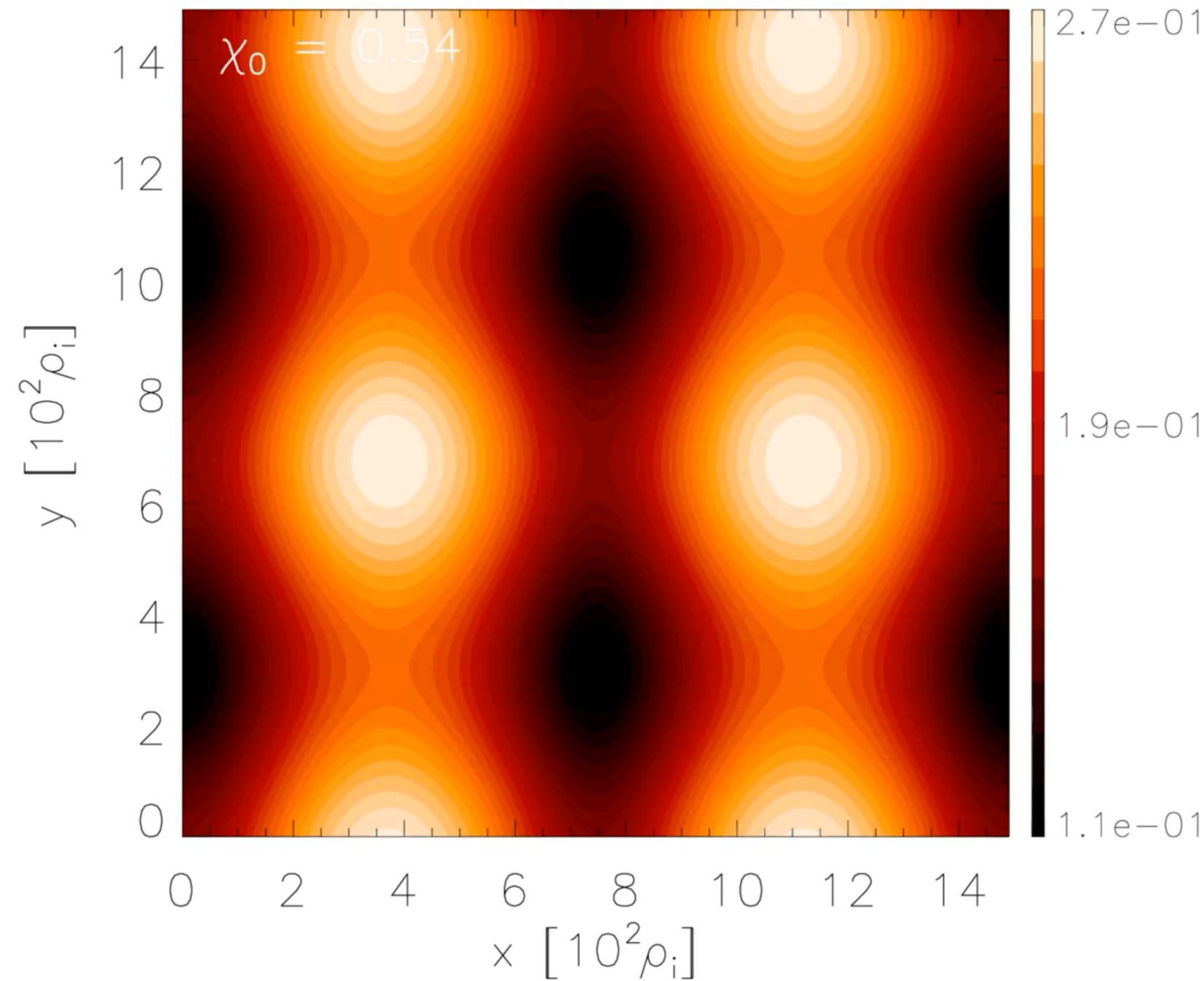
Recent Developments via Numerical Simulations

[Cerri, Passot, Laveder, Sulem, Kunz, ApJL (to be submitted)]

$$\langle \delta \mathbf{b}_\perp \rangle_z / \mathbf{B}_0 \quad (\chi_0 \sim 0.1)$$



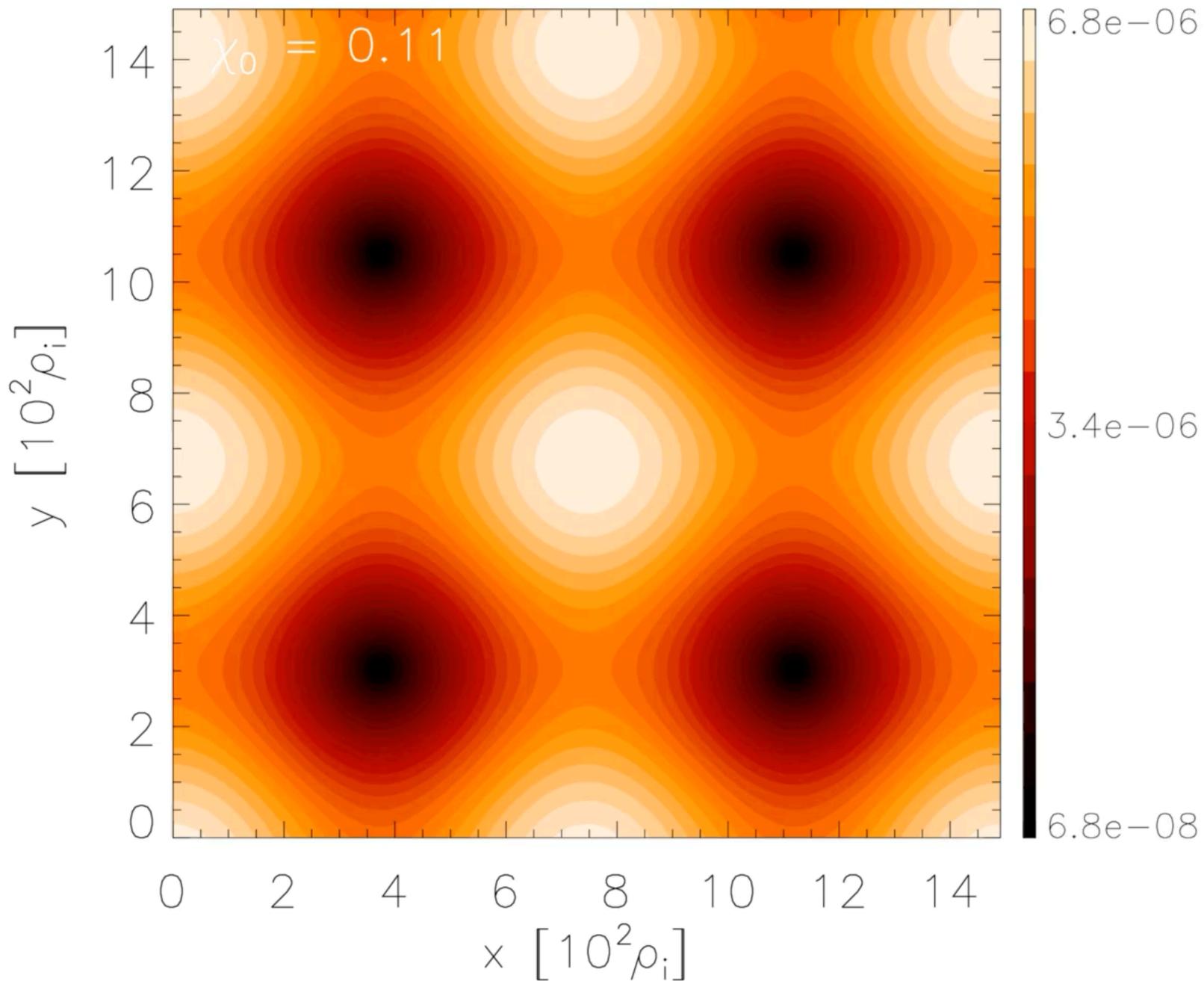
$$\langle \delta \mathbf{b}_\perp \rangle_z / \mathbf{B}_0 \quad (\chi_0 \sim 0.5)$$



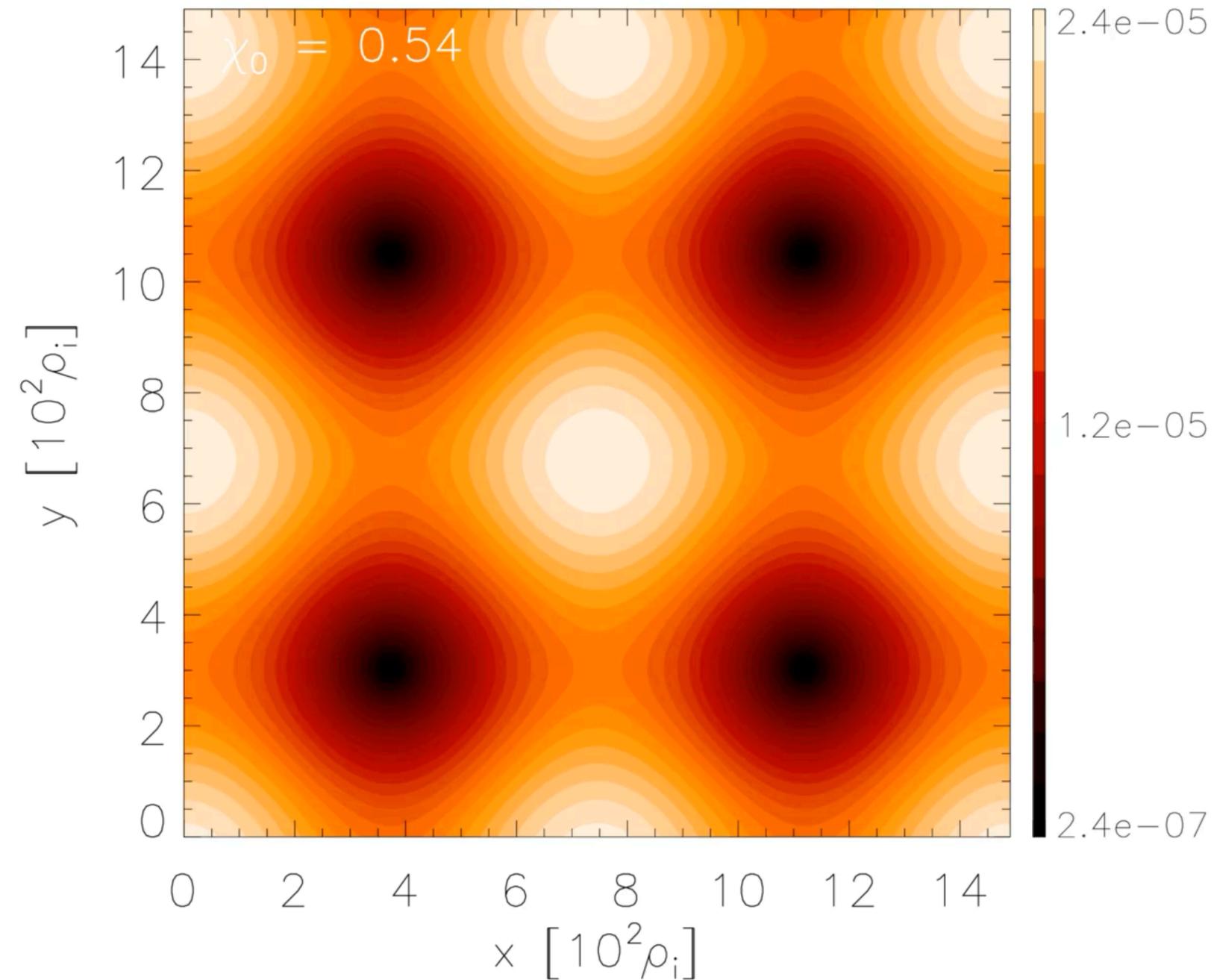
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$\delta \mathbf{b}_\perp |_{z=L/2} / \mathbf{B}_0$ ($\chi_0 \sim 0.1$)



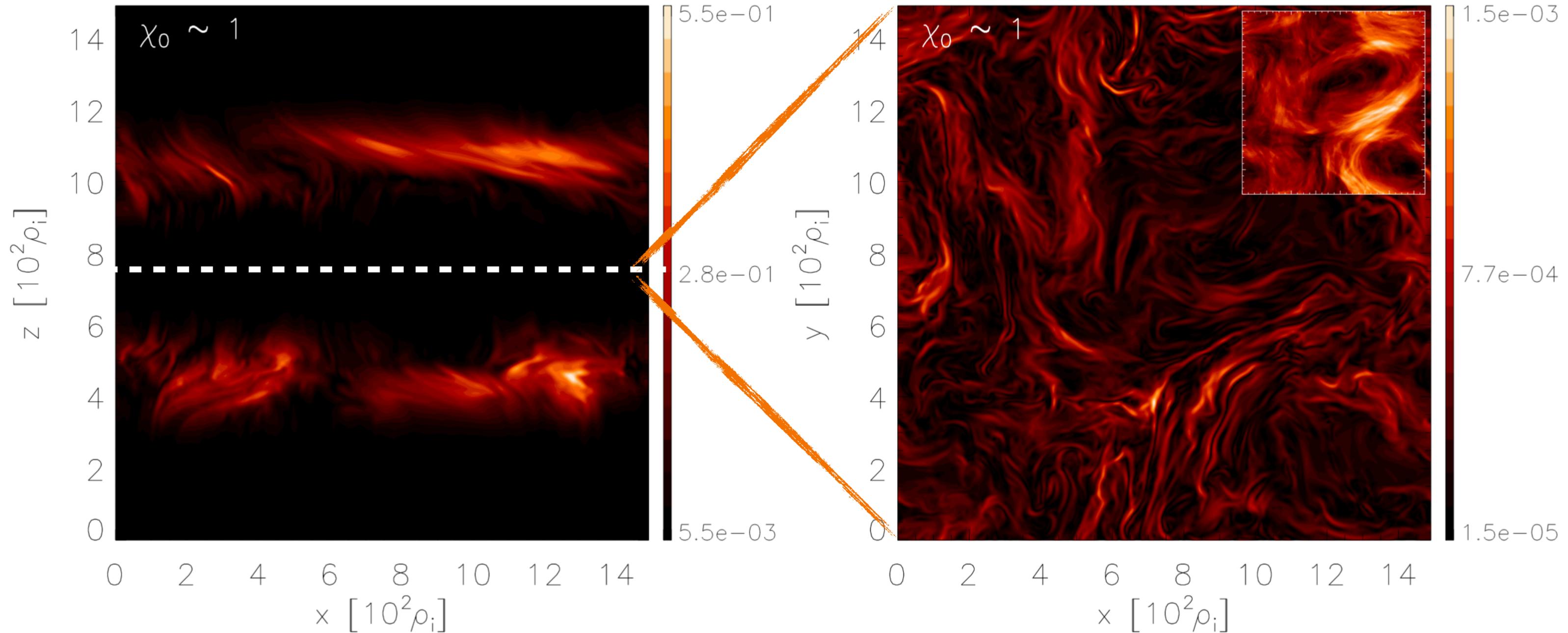
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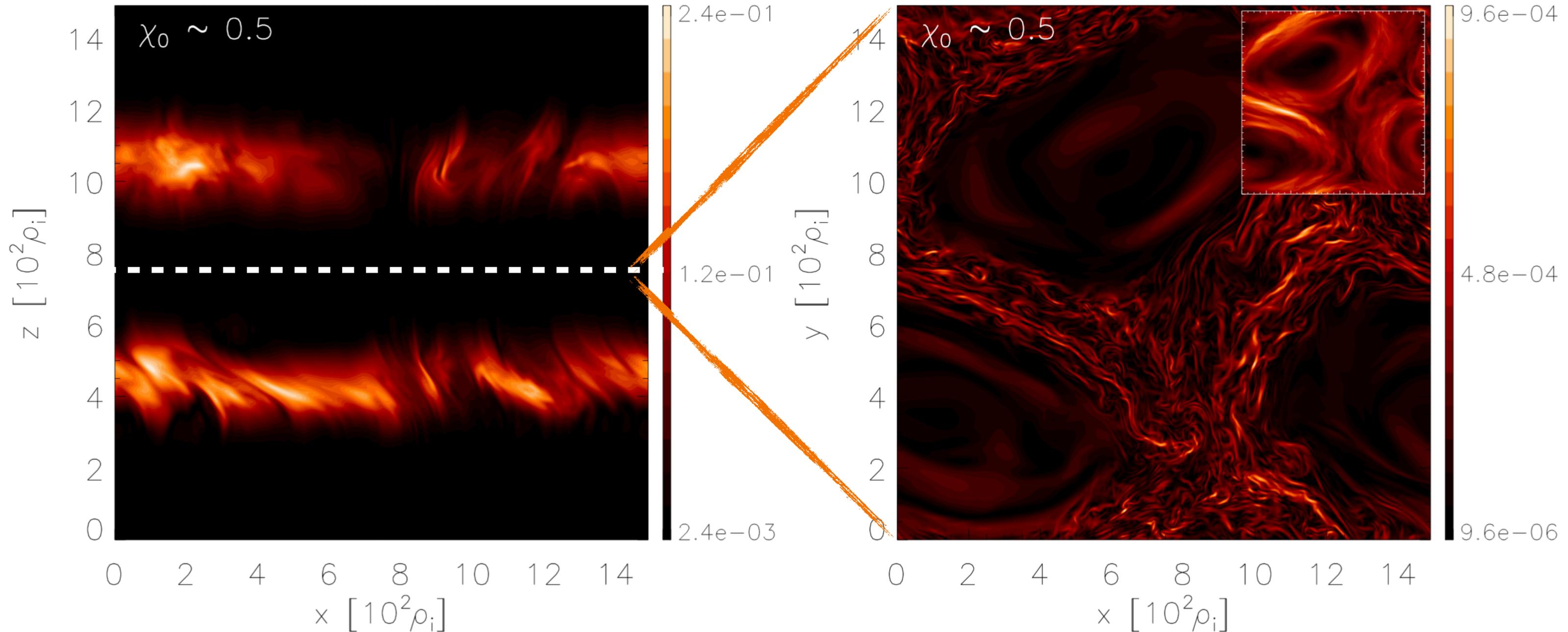
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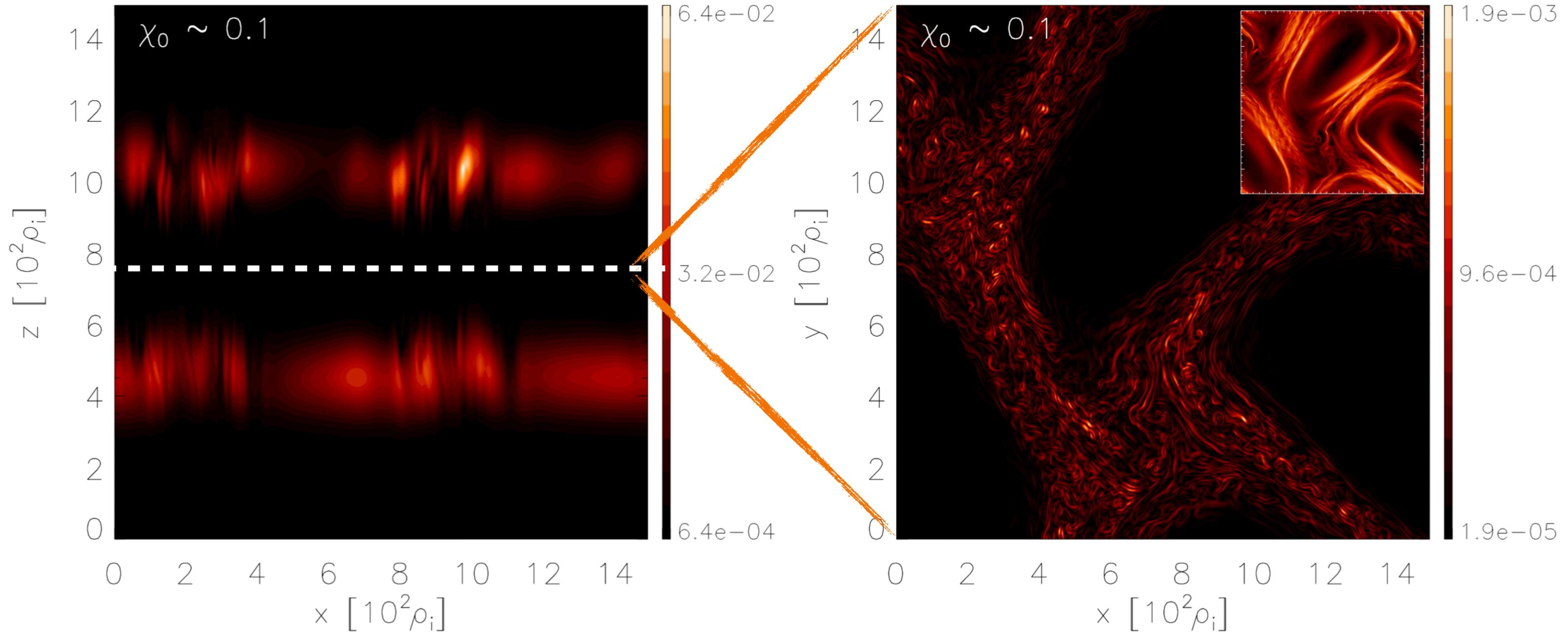
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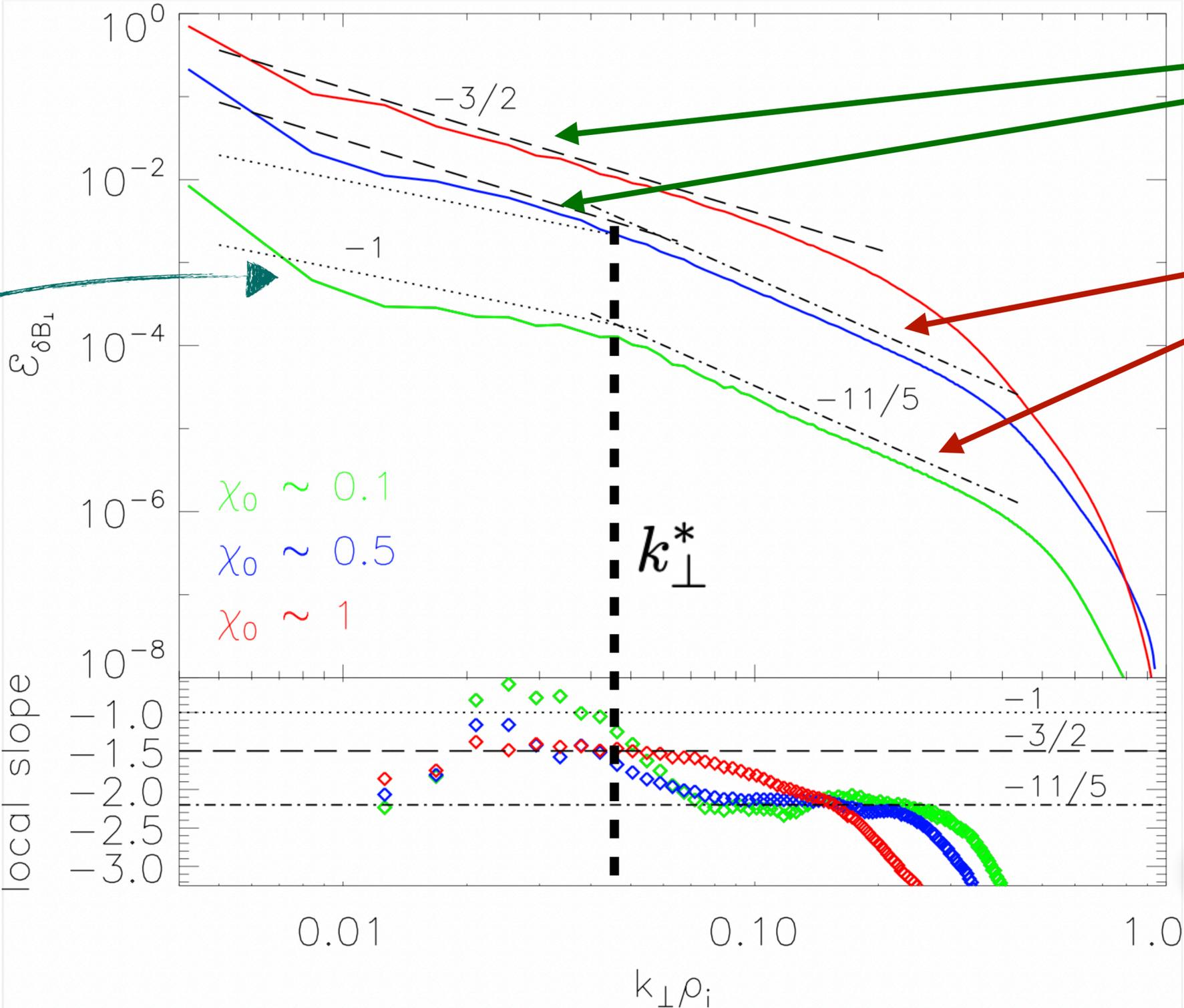
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[Cerri, Passot, Laveder, Sulem, Kunz, ApJL (to be submitted)]



reconnection-mediated regime obtained even before a large-scale turbulence state is generated (just from reconnection of current sheet generated by AWs shearing each other)

$-3/2$

$-11/5$

First proof in 3D!!!
(from reduced-MHD)

Take home message(s)

☞ *The fate of weak MHD turbulence is to become strong...*

...but which type of strong MHD turbulence?

☞ *A new turbulence regime exists, mediated by magnetic reconnection*

...emergence of reconnection-mediated turbulence depends both on the Lundquist number and/or on the strength of the nonlinearities

☞ We have now provided a *proof via 3D simulations (from a first-principle setup)*

☞ **BUT: there are still a lot of open questions... we need smart(er than me) people!**

Thank you for your attention!