

Kelvin-wave turbulence

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Collaborators

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Outline

- I. Introduction
 - · Classical vs. quantum turbulence

II. Kelvin Wave Turbulence Theory

Hamiltonian description, resonant wave interactions, kinetic equations, locality

III. Numerical Simulations

Biot-Savart and Gross-Pitaevskii equations

Nonlinearity, Complex Phenomena and Universality for Waves 15th-20th May 2022, Ile de Porquerolles, France

Turbulence at Large Scales

Polarized vortex bundles and K41

- Polarization of quantum vortex lines into bundles mimic vortex tubes of classical turbulence
- Polarization can be induced from large-scale mixing or from normal fluid at finite-temperature
- Similar Richardson cascade scenario of cascading vortex bundles
- Observation of Kolmogorov energy spectrum





Navier-Stokes



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Baggaley, JL, Barenghi, Phys. Rev. Lett. **109**, 205304, (2012)

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Navier-Stokes



Quantum vortex reconnections

- The classical-quantum vortex bundle analogy breaks down at scales near or below the inter-vortex scale ℓ
- Quantum vortex reconnections become important for the redistribution of energy





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Mechanisms of energy transport

- 1. Vortex ring emission
 - Rings emitted from reconnection region, directly transferring energy through tangle
- 2. Direct sound emission
 - Phonon emission at reconnection point
- 3. Generation of Kelvin waves
 - Energy and momentum transferred to helical Kelvin waves that propagate along individual quantized vortex lines





Vortex ring cascade at large angles

- A vortex reconnection of two (almost) anti-parallel vortices lead to a series of self-reconnections and the emission of multiple vortex rings
- Critical angle for ring generation in the Biot-Savart model is $\theta_c \simeq 0.942\pi$





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Reconnection angles in QT tangles

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Modulational instability and self-reconnection

 Strongly nonlinear Kelvin waves can lead to modulational instability and self reconnections

Salman, Phys. Rev. Lett. 111, 165301, (2013)













Isotropic homogeneous small-scale QT

- Polarization inhibits ring emission
- Vortex reconnections transfer large-scale energy to Kelvin waves at superfluid cross-over region
- Possible thermalisation at the intervortex scale
- Weakly nonlinear Kelvin wave interactions transfer energy to even smaller scales



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Wave turbulence description of Kelvin-wave cascade

• Theory for the non-equilibrium statistical description of the weakly nonlinear interaction of an ensemble of waves





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- Wave turbulence description of Kelvin-wave cascade
- Theory for the non-equilibrium statistical description of the weakly nonlinear interaction of an ensemble of waves
- Main theoretical results
 - 1. Nonlinear kinetic wave equation
 - 2. Steady-state power-law spectra for constant flux transfer of invariants
 - 3. But can easily study nonlinear evolution of higher-order moments and amplitude PDFs





Biot-Savart Hamiltonian description

$$\dot{\mathbf{s}} = \frac{\kappa}{4\pi} \oint_{\mathcal{L}} \frac{\mathbf{r} - \mathbf{s}}{\left|\mathbf{r} - \mathbf{s}\right|^3} \times \mathrm{d}\mathbf{r}$$



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$$a(z,t) = x(z,t) + iy(z,t)$$
 $i\kappa \frac{\partial a}{\partial t} = \frac{\delta \mathcal{H}}{\delta a^*}$

$$x = y = 0$$

$$\mathbf{s} = [x(z), y(z), z]$$

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$$\mathcal{H} = \frac{\kappa^2}{4\pi} \int \frac{1 + \operatorname{Re}\left[a^{\prime *}(z_1)a^{\prime}(z_2)\right]}{\sqrt{(z_1 - z_2)^2 + |a(z_1) - a(z_2)|^2}} \, \mathrm{d}z_1 \mathrm{d}z_2$$

Svistunov, Phys. Rev. B, **52**, 3647, (1995)



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- Regularization of integral by introducing cut-off $\xi < |z_2 - z_1|$

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Truncation and weak nonlinear expansion

- Regularization of integral by introducing cut-off $\xi < |z_2 z_1|$
- Expand Hamiltonian in powers of the canonical variable:

$$\epsilon = \frac{|a(z_1) - a(z_2)|}{|z_1 - z_2|} \ll 1$$

$$\mathcal{H} = \mathcal{H}_2 + \mathcal{H}_4 + \mathcal{H}_6 + \cdots$$

Hamiltonian-Fourier Representation



Wave action representation of the Hamiltonian

• Introduce wave action variables $a(z,t) = \kappa^{-1/2} \sum_{\mathbf{k}} a_{\mathbf{k}}(t) \exp(i \mathbf{k} z)$

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$$\mathcal{H} = \sum_{\mathbf{k}} \omega_{\mathbf{k}} a_{\mathbf{k}} a_{\mathbf{k}}^* + \frac{1}{4} \sum_{1,2,3,4} T_{3,4}^{1,2} a_1 a_2 a_3^* a_4^* \delta_{3,4}^{1,2} + \frac{1}{36} \sum_{1,2,3,4,5,6} W_{4,5,6}^{1,2,3} a_1 a_2 a_3 a_4^* a_5^* a_6^* \delta_{4,5,6}^{1,2,3} a_4 a_5^* a_6^* \delta_{4,5,6}^{1,2,3} a_5 a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_5^* a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_6^* a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_6^* a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_6^* a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_6^* a_6^* a_6^* a_6^* \delta_{4,5,6}^{1,2,3} a_6^* a_6^$$

$$a_1 = a_{\mathbf{k}_1}(t) \qquad T_{3,4}^{1,2} = T(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4) \qquad \delta_{3,4}^{1,2} = \delta(\mathbf{k}_1 + \mathbf{k}_2 - \mathbf{k}_3 - \mathbf{k}_4)$$

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Interaction coefficients

$$\omega_{\mathbf{k}} = \frac{\kappa\Lambda}{4\pi} \mathbf{k}^{2} - \frac{\kappa}{4\pi} \mathbf{k}^{2} \ln(\mathbf{k}\ell_{\text{eff}}), \qquad \Lambda = \ln\left(\ell_{\text{eff}}/\tilde{\xi}\right) \gg 1, \quad \tilde{\xi} = \xi e^{\gamma + \frac{3}{2}}$$
$$T_{3,4}^{1,2} = -\frac{\Lambda}{4\pi} \mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4} - \frac{1}{16\pi} \left[5\mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4} + \mathcal{F}_{3,4}^{1,2}\right]$$
$$W_{4,5,6}^{1,2,3} = \frac{9\Lambda}{8\pi\kappa} \mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4} \mathbf{k}_{5} \mathbf{k}_{6} + \frac{9}{32\pi\kappa} \left[7\mathbf{k}_{1} \mathbf{k}_{2} \mathbf{k}_{3} \mathbf{k}_{4} \mathbf{k}_{5} \mathbf{k}_{6} + \mathcal{G}_{4,5,6}^{1,2,3}\right]$$

• Separate logarithm divergent terms by introducing an effective length scale ℓ_{eff} • $\mathcal{F}_{3,4}^{1,2}$ and $\mathcal{G}_{4,5,6}^{1,2,3}$ are terms containing logarithmic contributions

JL *et al.* Phys. Rev. B, **81**, 104526, (2010)

Leading Order Integrability

Local Induction Approximation (LIA)

- If the cutoff is small then terms proportional to Λ give greatest contribution and diverge in the limit $\,\xi \to 0\,$
- Keeping only the leading divergent terms, then the Hamiltonian becomes

$$\mathcal{H} = \frac{\kappa^2 \Lambda}{2\pi} \int \sqrt{1 + \left|a'(z)\right|^2} \,\mathrm{d}z$$



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- Shown to be equivalent to the Local Induction Approximation (LIA)
- LIA implies only neighbouring vortex elements determine evolution and corresponds to integrable dynamics
- Subleading in Λ (non-LIA) terms are essential for turbulent Kelvin-wave interactions



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Double expansion in nonlinearity $\epsilon \ll 1$ and divergence $\Lambda^{-1} \ll 1$





Wave resonance

• Waves only transfer energy and momentum to each other when in resonance



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- In principle, only need to expand Hamiltonian up to first nonlinear term: $\mathcal{H} = \mathcal{H}_2 + \mathcal{H}_4$



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- Mode evolution equation

$$i\frac{\partial a_{\mathbf{k}}}{\partial t} = \frac{\delta \mathcal{H}}{\delta a_{\mathbf{k}}^*}$$

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$$i\frac{\partial a_{\mathbf{k}}}{\partial t} = \frac{\delta \mathcal{H}}{\delta a_{\mathbf{k}}^*} = \omega_k a_{\mathbf{k}} + \frac{1}{2} \sum_{1,2,3} T_{3,\mathbf{k}}^{1,2} a_1 a_2 a_3^* \delta_{3,\mathbf{k}}^{1,2}$$

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• Change variable into rotating coordinate frame $b_{\mathbf{k}} = a_{\mathbf{k}} \exp \left(i \, \omega_{\mathbf{k}} \, t \right)$



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Four-wave resonance condition Momentum con

Momentum conservation $\mathbf{k}_1 + \mathbf{k}_2 = \mathbf{k}_3 + \mathbf{k}$

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Energy conservation

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$$i\frac{\partial b_{\mathbf{k}}}{\partial t} = \frac{1}{2} \sum_{1,2,3} T_{3,\mathbf{k}}^{1,2} \ b_1 b_2 b_3^* \ \delta_{3,\mathbf{k}}^{1,2} \ \exp\left(-i \ \omega_{3,\mathbf{k}}^{1,2} \ t\right)$$

• Main nonlinear contribution when frequencies cancel: $\omega_{3,\mathbf{k}}^{1,2} \equiv \omega_1 + \omega_2 - \omega_3 - \omega_{\mathbf{k}} = 0$

Four-wave resonance condition

Momentum conservation

• This means that there are essentially two delta functions:

$$\mathbf{k}_1 + \mathbf{k}_2 = \mathbf{k}_3 + \mathbf{k}$$
$$\omega_1 + \omega_2 = \omega_3 + \omega_{\mathbf{k}}$$

 $\omega_1 + \omega_2 - \omega_3 + \omega_k$ Energy conservation T

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Only trivial resonances can solve resonance condition for Kelvin-wave frequency

 $\mathbf{k}_1 = \mathbf{k}_3, \quad \mathbf{k}_2 = \mathbf{k}, \quad \text{or} \quad \mathbf{k}_1 = \mathbf{k}, \quad \mathbf{k}_2 = \mathbf{k}_3$







Canonical transformation

• Trivial 4-wave resonances only lead to a nonlinear frequency shift of the linear dynamics



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Wave action density



Wave action density

• Of particular interest is the second order correlator function $\langle a_{\mathbf{k}}a_{\mathbf{k}_{1}}^{*}\rangle = n_{\mathbf{k}}\delta(\mathbf{k}-\mathbf{k}_{1})$



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Kolmogorov-Zakharov power-law solutions



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Locality

• With any KZ solutions, convergence of the collision integral must be ensured in order for the realizability of the stationary state



Nonlocal wave interactions

 Exact calculation of interaction coefficient enabled us to prove that six-wave collision integral diverges in the limit of two long Kelvin-waves

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JL *et al.* Phys. Rev. B, **81**, 104526, (2010) L'vov, Nazarenko, Low Temp. Phys. **36**, 785, (2010)



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Effective four-wave kinetic description

$$\frac{\partial n_{\mathbf{k}}}{\partial t} = \frac{\epsilon^8 \pi}{12} \int \left\{ |V_{\mathbf{k}}^{1,2,3}|^2 n_1 n_2 n_3 n_{\mathbf{k}} \left[\frac{1}{n_{\mathbf{k}}} - \frac{1}{n_1} - \frac{1}{n_2} - \frac{1}{n_3} \right] \delta_{1,2,3}^{\mathbf{k}} \delta \left(\omega_{1,2,3}^{\mathbf{k}} \right) \right. \\ \left. + 3 |V_1^{\mathbf{k},2,3}|^2 n_1 n_2 n_3 n_{\mathbf{k}} \left[\frac{1}{n_1} - \frac{1}{n_k} - \frac{1}{n_2} - \frac{1}{n_3} \right] \delta_{\mathbf{k},2,3}^1 \delta \left(\omega_{\mathbf{k},2,3}^1 \right) \right\} \, \mathrm{d}\mathbf{k}_1 \, \mathrm{d}\mathbf{k}_2 \, \mathrm{d}\mathbf{k}_3$$





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Boué et al. Phys. Rev. B, 84, 064516, (2011)



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Identification of Spectrum



History of simulations

- Many previous simulations (Vinen, Tsubota; Kozik, Svisuntov; Barenghi, Baggaley)
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Gross-Pitaevskii equation

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Conclusions and Perspectives



Energy dissipation in small-scale QT

- Evidence to say that Kelvin-waves are important for small-scale energy transfer for polarized vortex tangles in homogeneous and isotropic turbulence
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Perspectives

- Can we quantify the amount of energy transferred to Kelvin waves?
- Are Kelvin-waves weakly nonlinear in reality?
- Observation of Kelvin-wave cascade in velocity energy spectrum?
- UK Fluids Network SIG in Wave Turbulence focussed on strong nonlinearities in WT